

10 Mpc

$z=1.2$



BH growth and the impact of AGN feedback on galaxy evolution

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Collaborators:

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1 Mpc

Horizon-AGN simulation
<http://horizon-AGN.projet-horizon.fr>

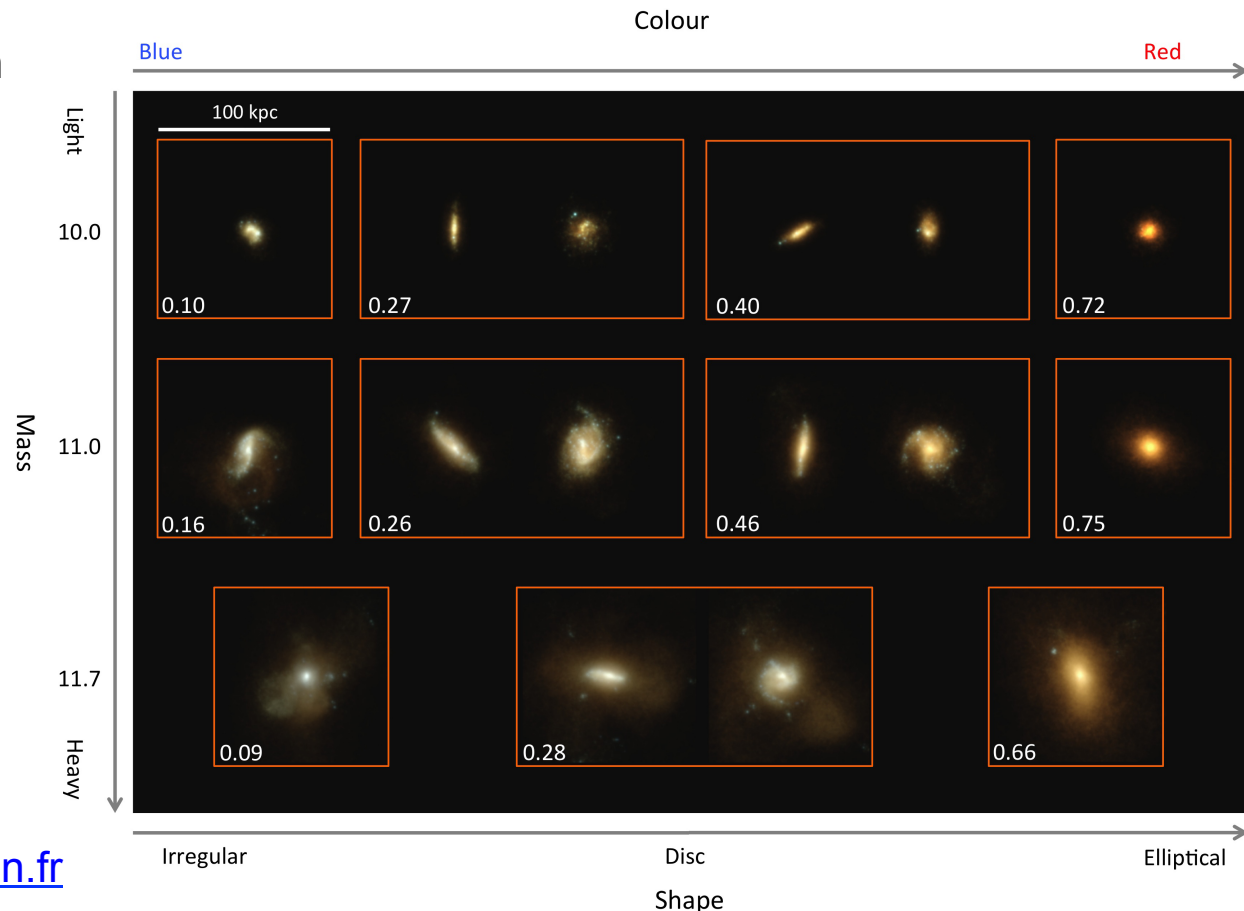
The Horizon-AGN simulation

- Ramses code (AMR) *Teyssier (2002)*
- $L_{\text{box}}=100 \text{ Mpc}/h$
- 1024^3 DM particles
- $M_{\text{DM,res}}=8 \times 10^7 M_{\text{sun}}$
- Finest cell resolution $dx=1 \text{ kpc}$
- Gas cooling & UV background heating
- Low efficiency star formation
- Stellar winds + SNII + SNIa
- O, Fe, C, N, Si, Mg, H
- AGN feedback radio/quasar

- $z=0.6$ using 6.7 Mhours on 4096 cores
- 150 000 galaxies per snapshot (> 50 part.)
- 7 billions leaf cells ($>$ Illustris or EAGLE)

- *Dubois et al, 2014*
- *Welker et al, 2014*
- *Codis et al, sub.*

<http://horizon-AGN.projet-horizon.fr>



Green: gas density / Red: temperature / Blue: metallicity

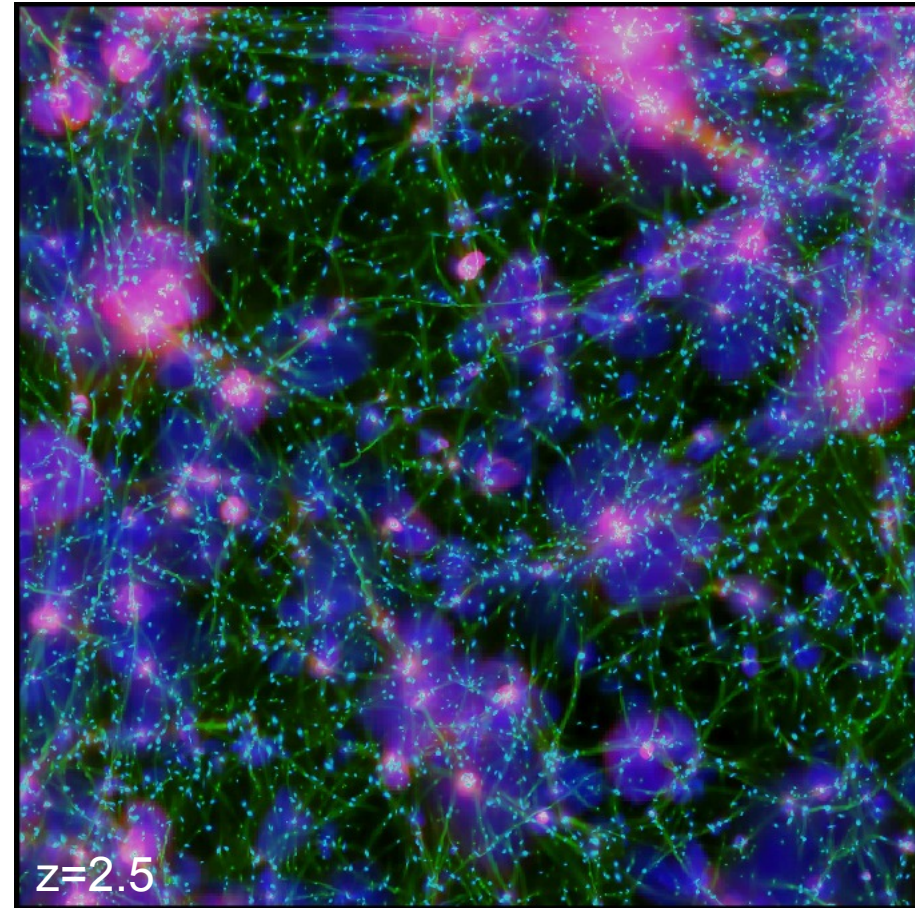
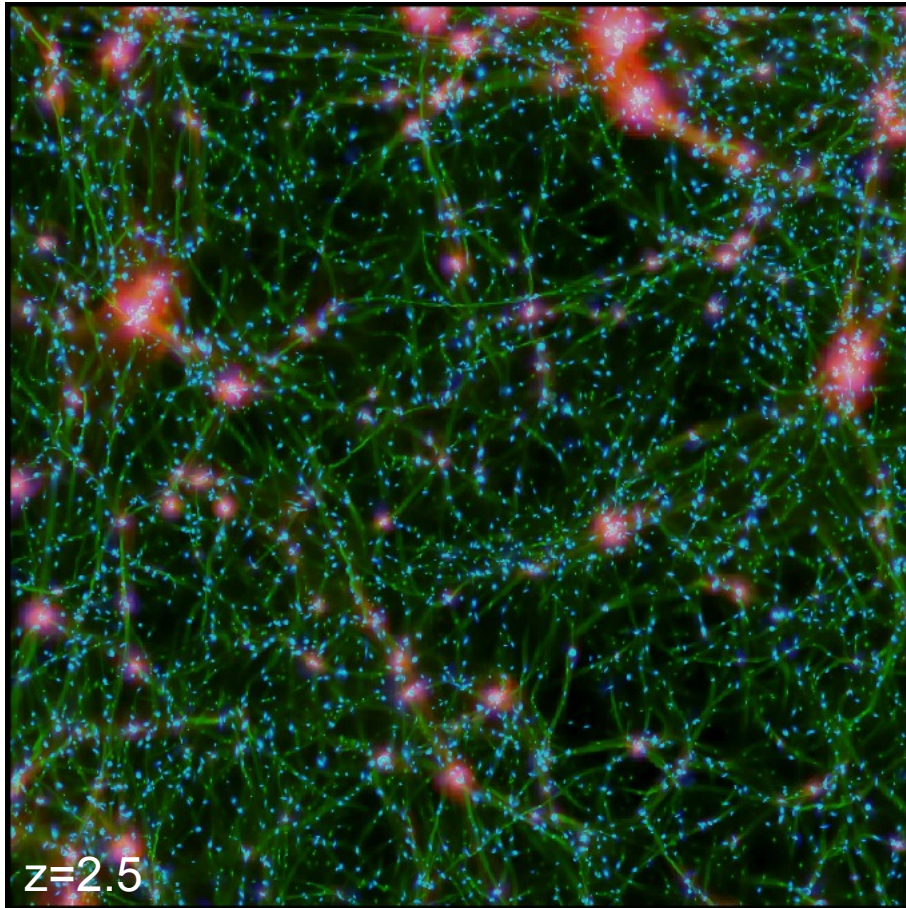
$z=38.305$

A visual inspection of the impact of AGN feedback on large-scale structures

Green: gas density / Red: temperature / Blue: metallicity

Without AGN

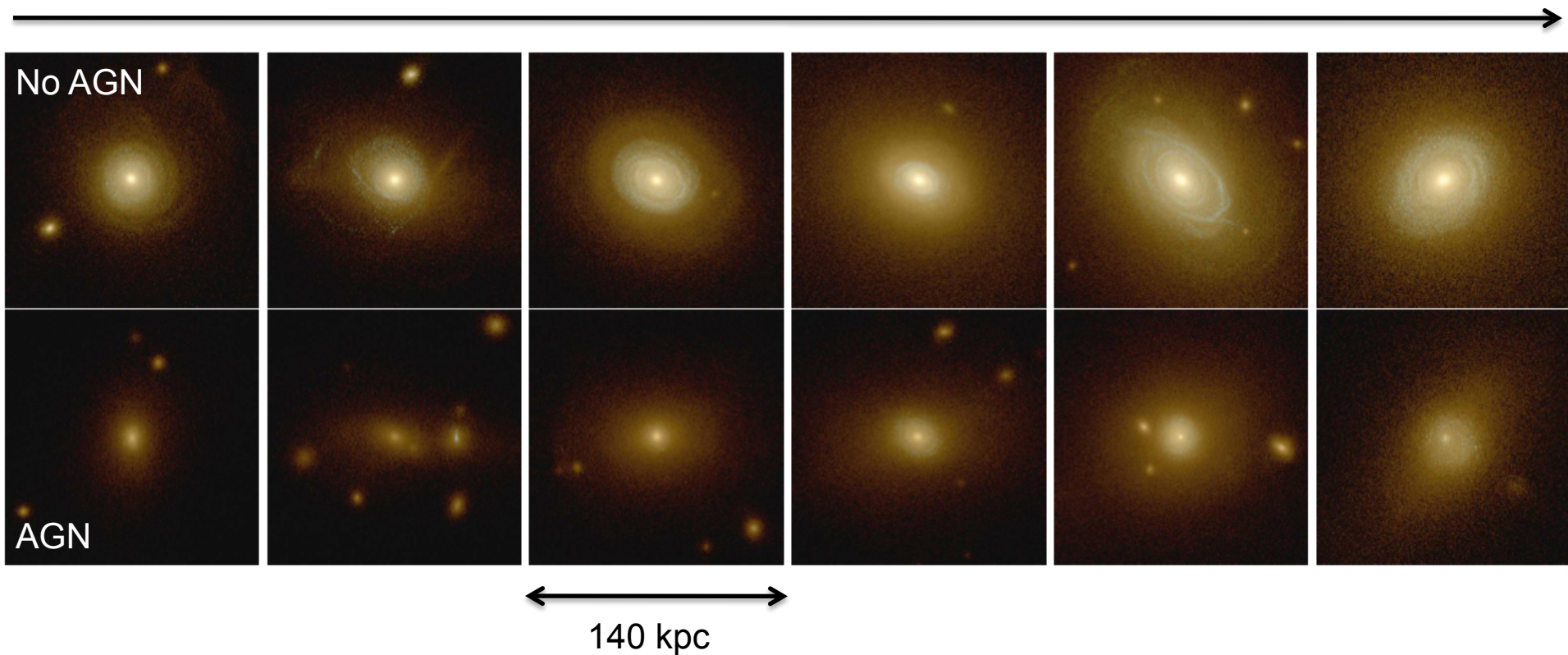
With AGN



AGN are responsible for turning discs into ellipticals

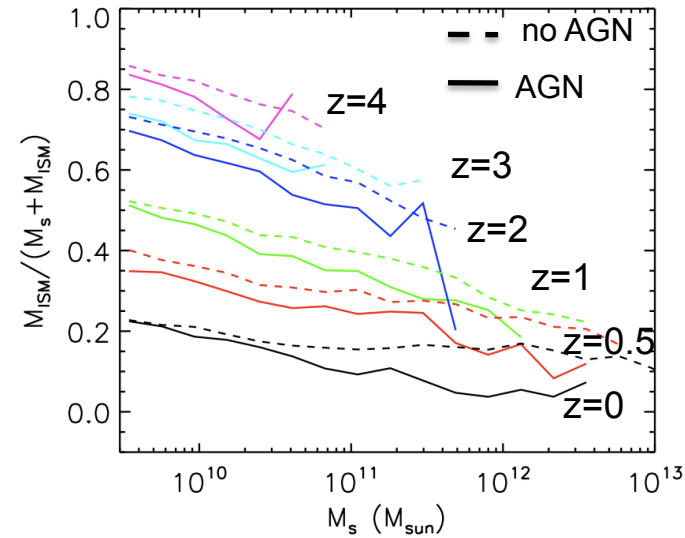
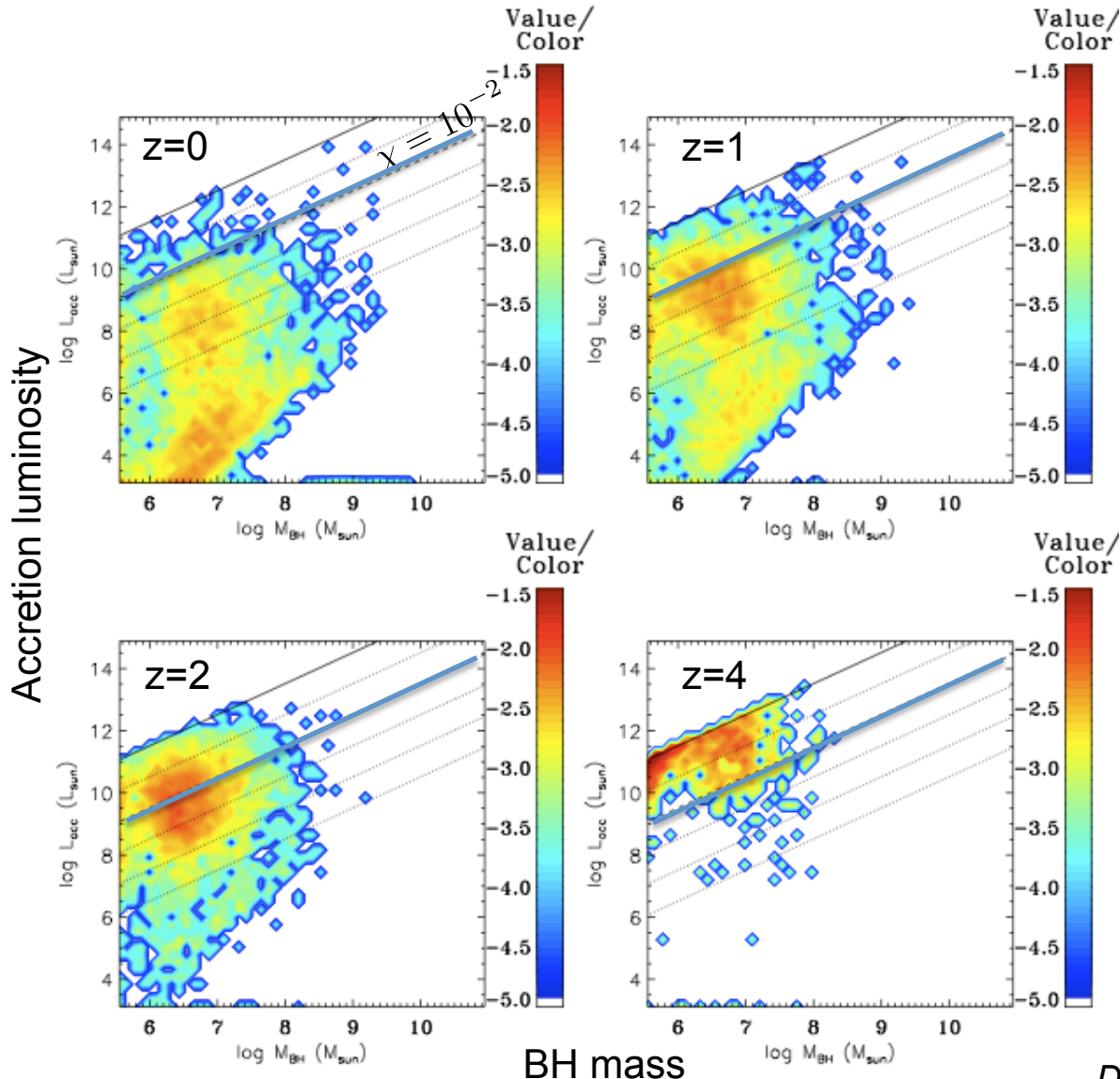
Central galaxies at $z=0$ in groups $10^{13} < M_{\text{halo}} < 10^{14} M_{\text{sun}}$

Increasing mass \rightarrow



Dubois, Gavazzi, Peirani, Silk, 2013

Radio mode or quasar mode ?



Galaxies are gas-rich at high-redshift
 Star formation and feedback
 consumes/removes/prevents from
 collapsing the gas

Dubois, Devriendt, Slyz, Teyssier, 2012

Growing the first bright quasars

Observational facts:

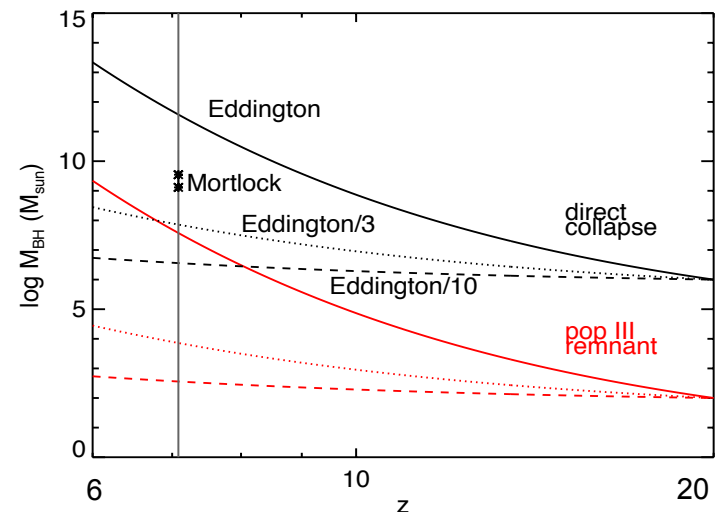
- Very bright quasars in the SDSS with $z > 6$
(Willott et al., 2003; Fan et al., 2006; Jiang et al., 2009)
- Detection of a $2 \cdot 10^9 M_{\text{sun}}$ BH at $z=7$
(Mortlock et al., 2011)

Requirement:

- Need to grow from 10^5 - $10^6 M_{\text{sun}}$ up to $10^9 M_{\text{sun}}$ in less than 700 Myrs ! Eddington limit provides an e-folding time = 45 Myr

Question:

- How to bring gas sufficiently rapidly into the bulge of the galaxy ?



Growing the first bright quasars

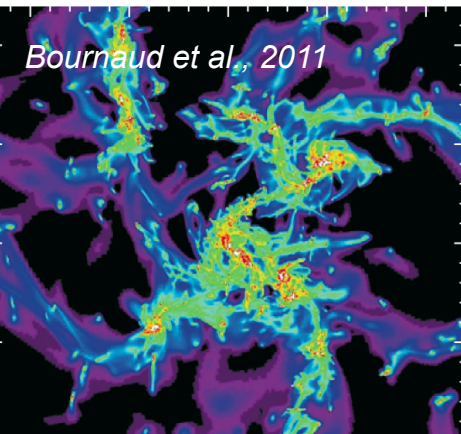
- Direct accretion from the cosmic cold flows (*Di Matteo et al., 2012*)

Cosmological context with large statistics but low resolution ($\sim 1\text{kpc}$)

Versus

- Violent disc instabilities (*Bournaud et al., 2011*)

High resolution (1pc) but isolated disc



Observational facts:

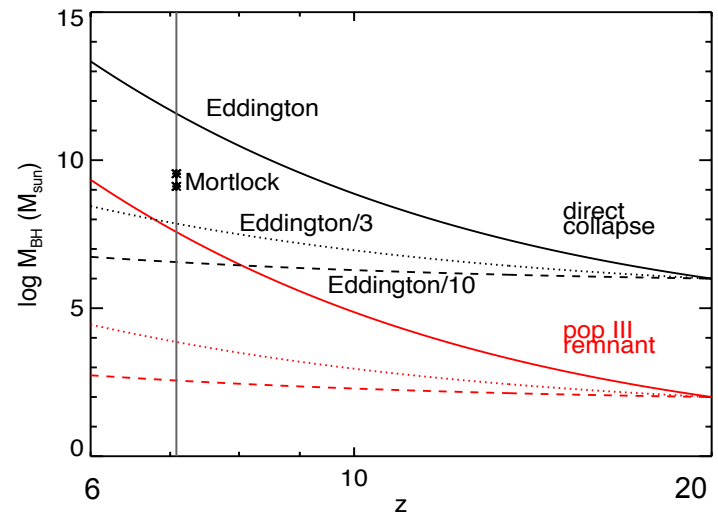
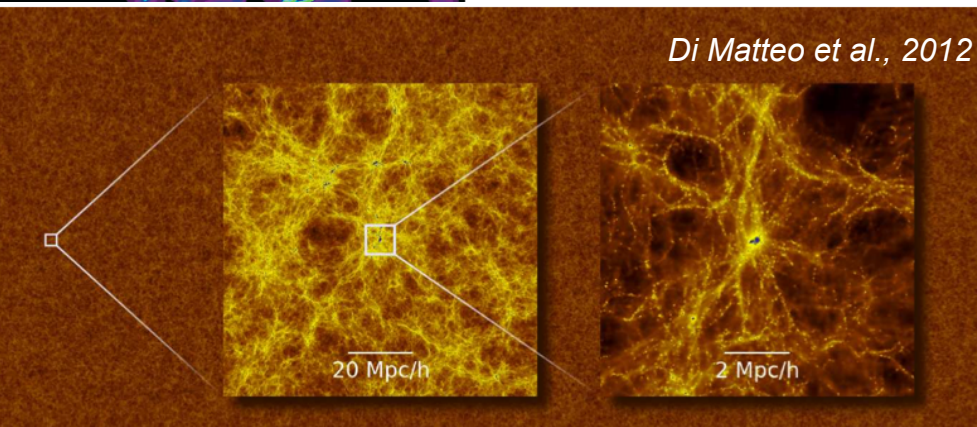
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Growing the first bright quasars

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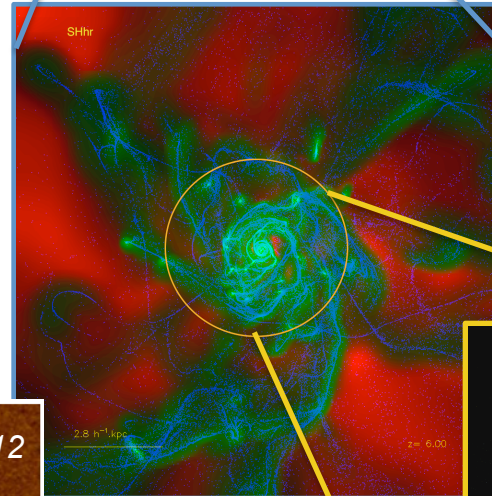
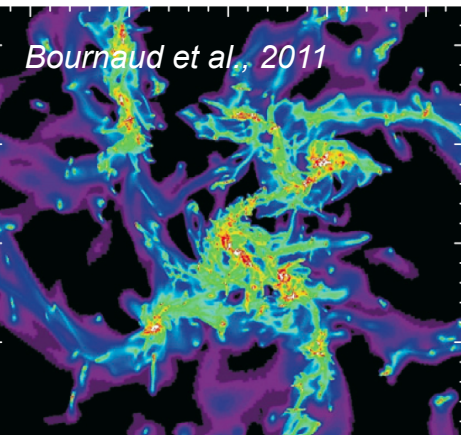
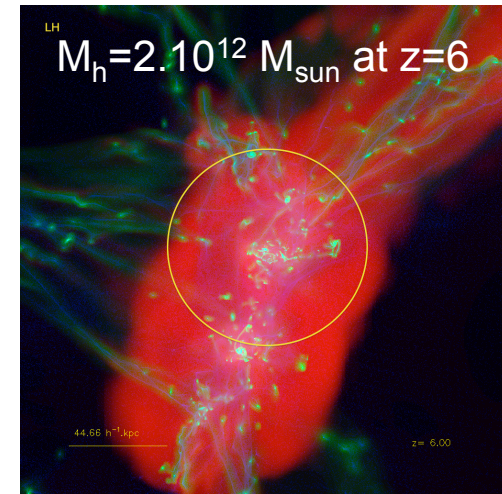
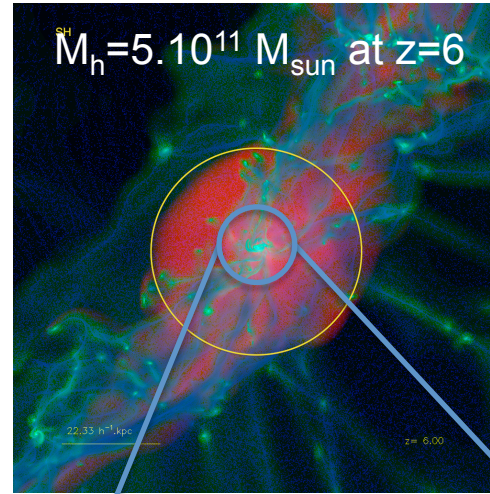
Cosmological context with large statistics but low resolution ($\sim 1\text{kpc}$)

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- Violent disc instabilities

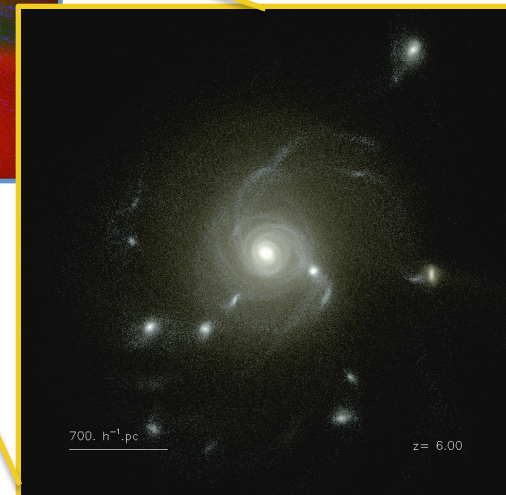
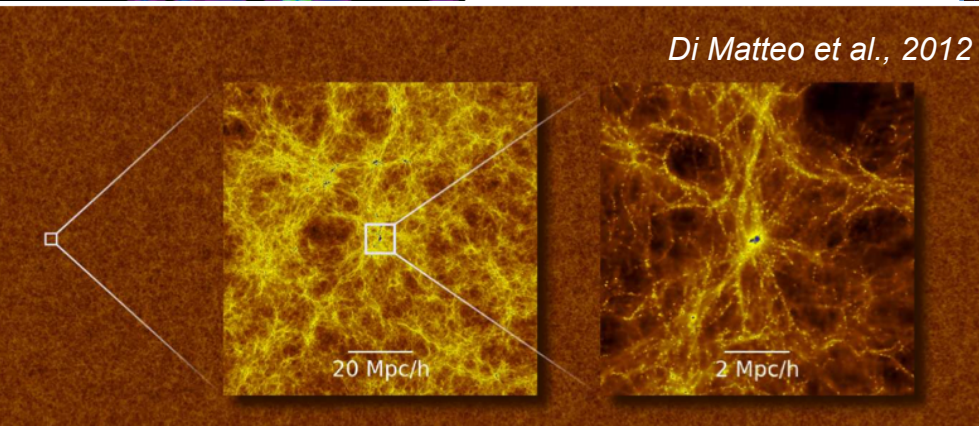
(*Bournaud et al., 2011*)

High resolution (1pc) but isolated disc



Cosmological zooms
10 pc resolution

*Dubois, Pichon,
Haehnelt et al., 2012*



Growing the first bright quasars

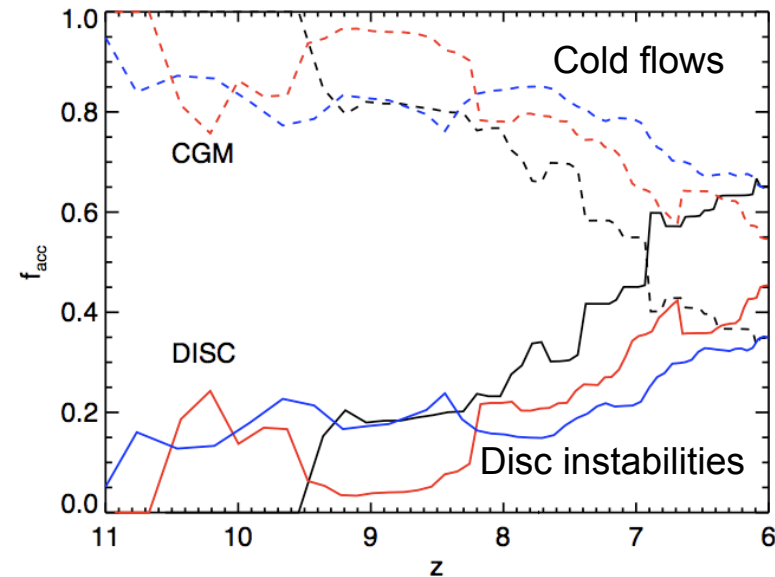
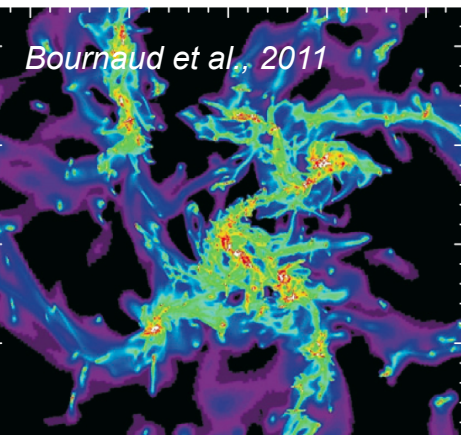
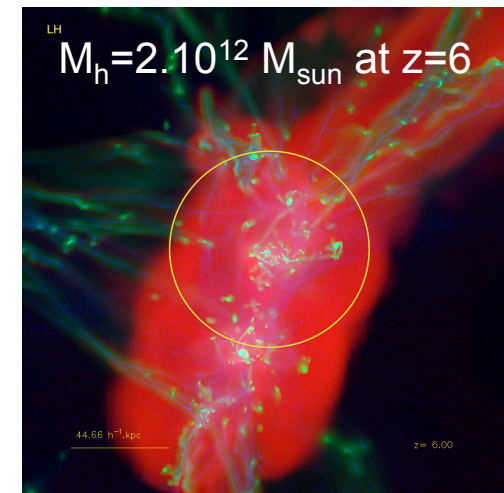
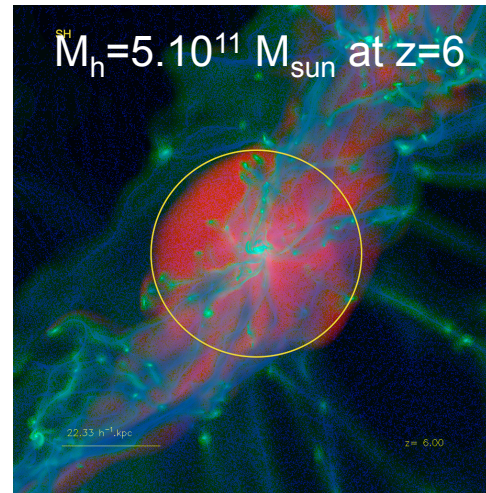
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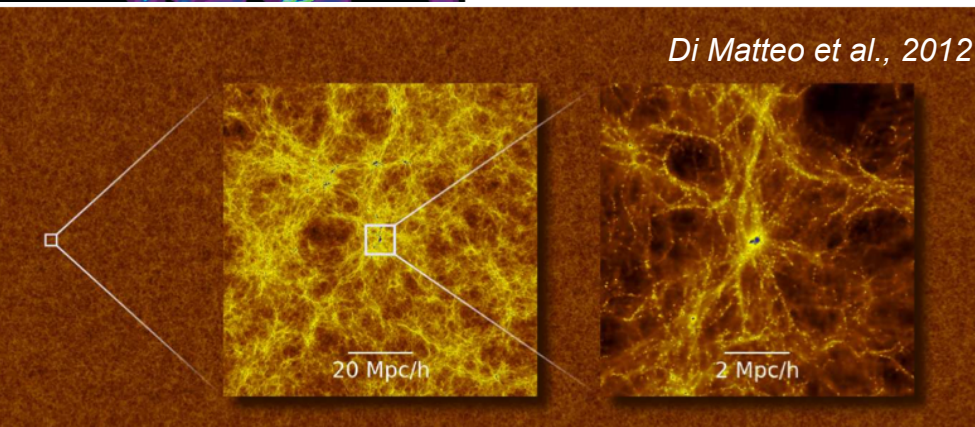
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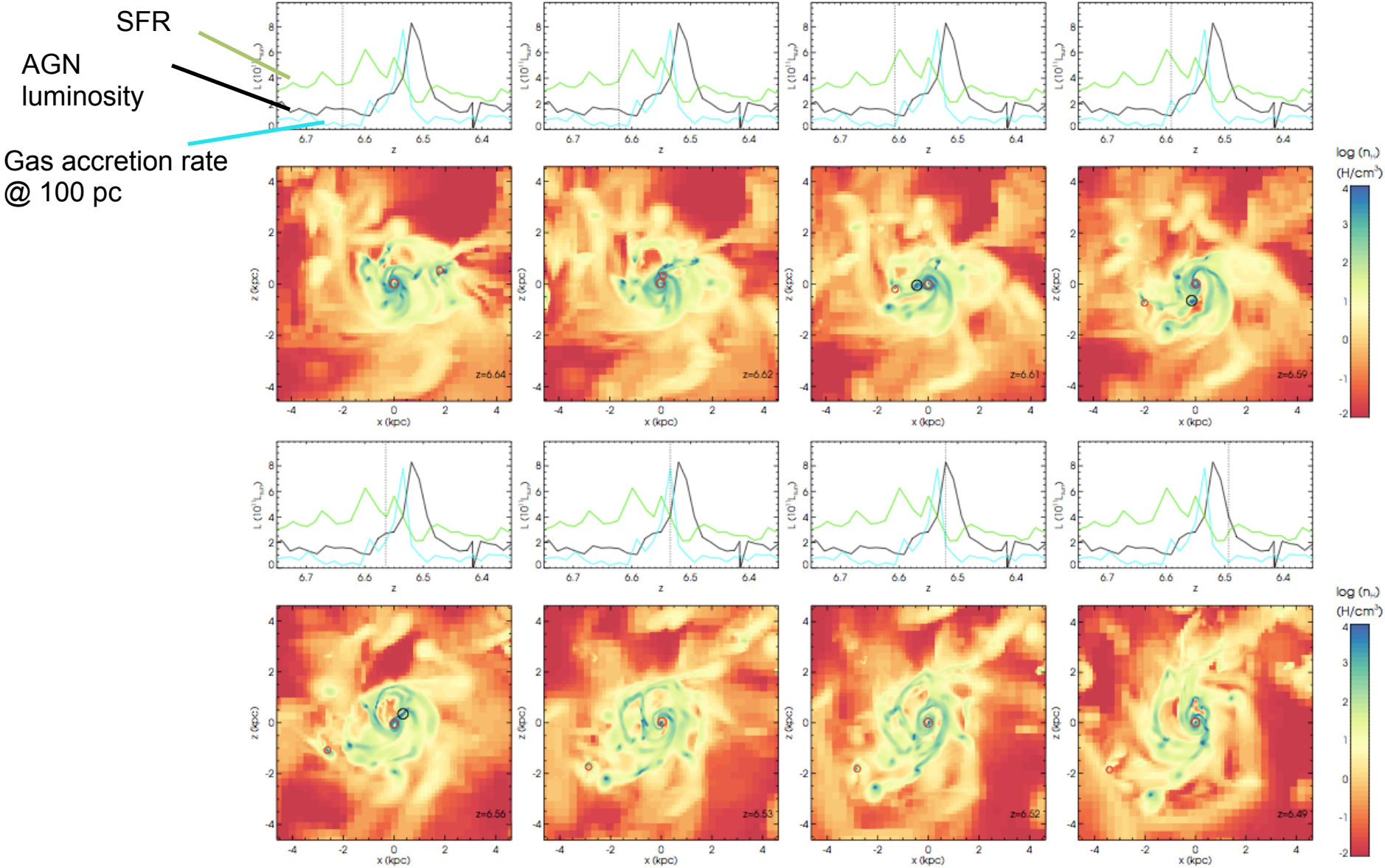
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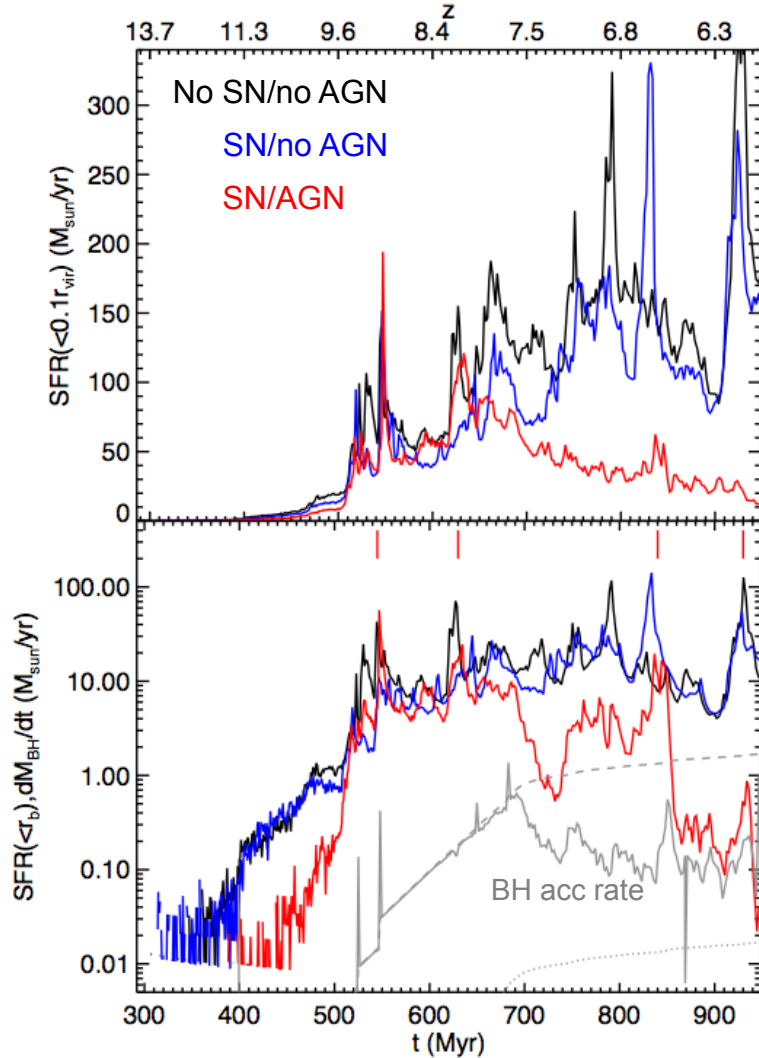
Dubois, Pichon, Haehnelt et al., 2012



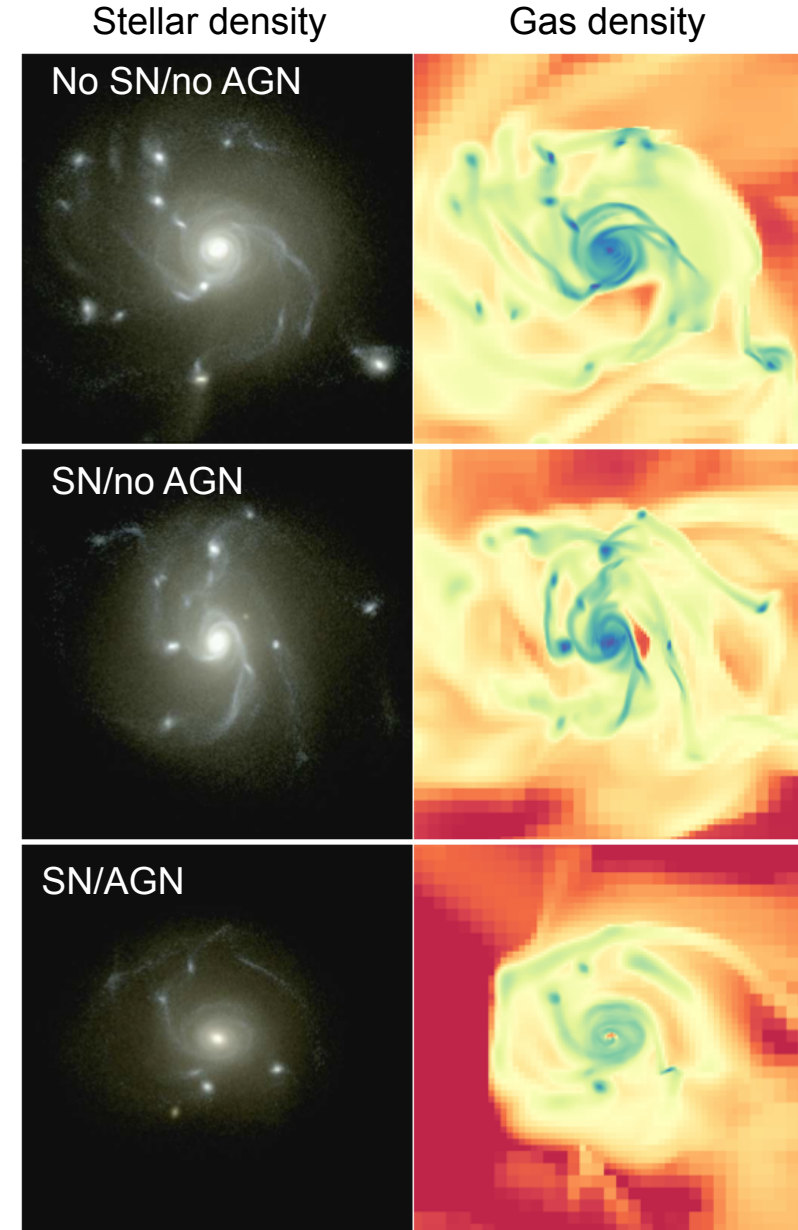
A rapid clump migration to trigger « late-time » AGN bursts



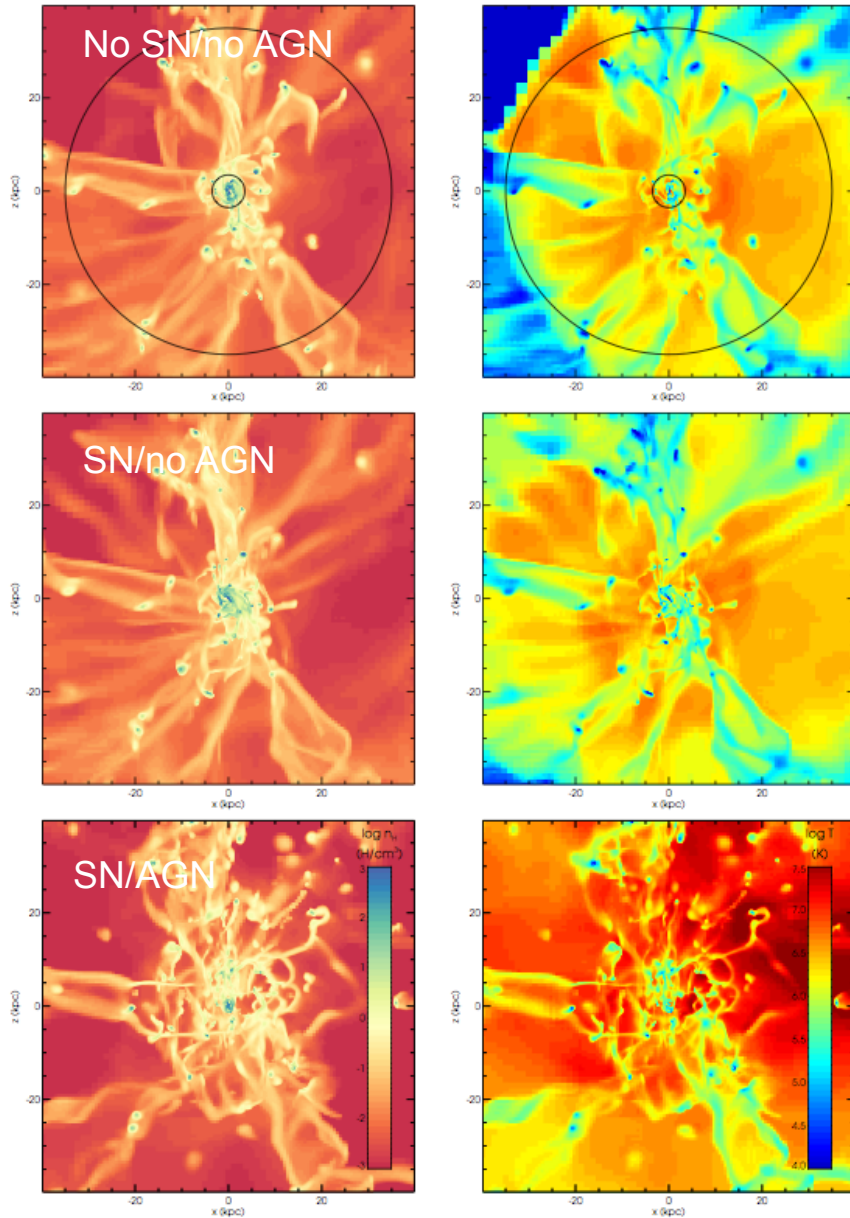
AGN quench star formation



Dubois, Pichon et al., 2013

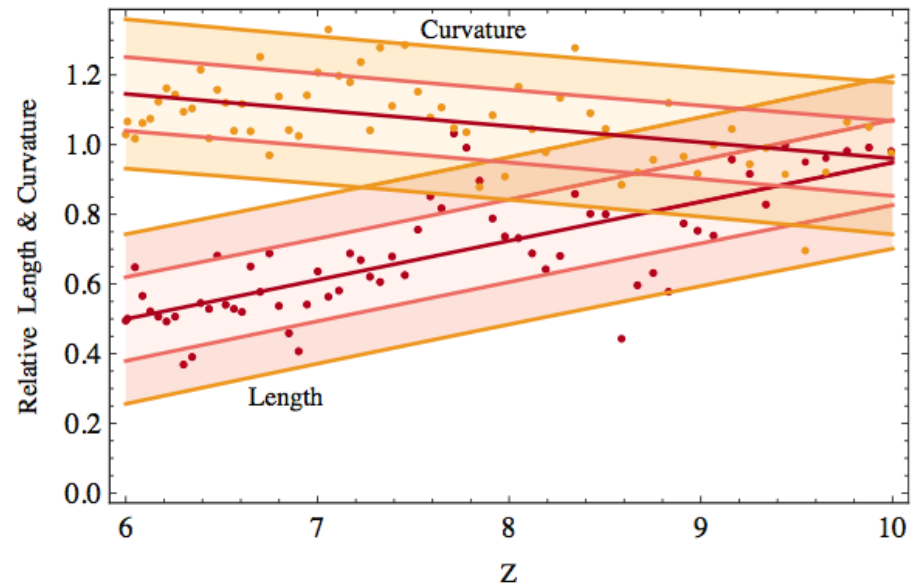


AGN blow cold flows away



Gas is driven out hot from the central galaxy due to AGN.

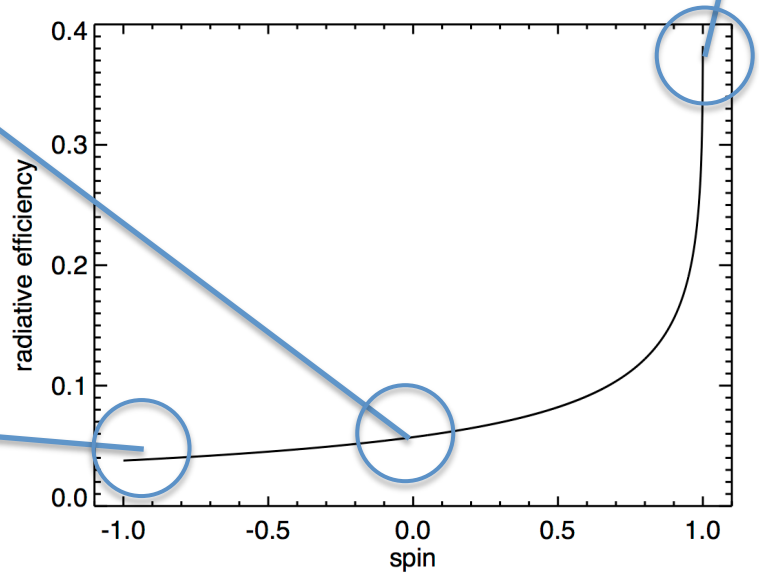
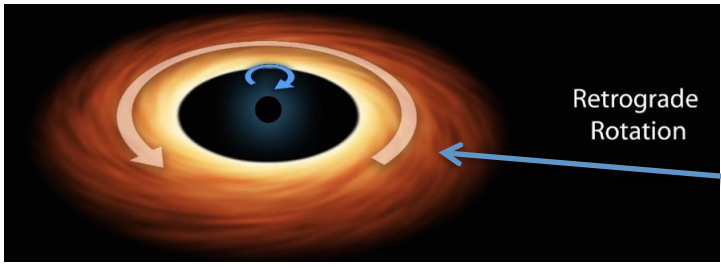
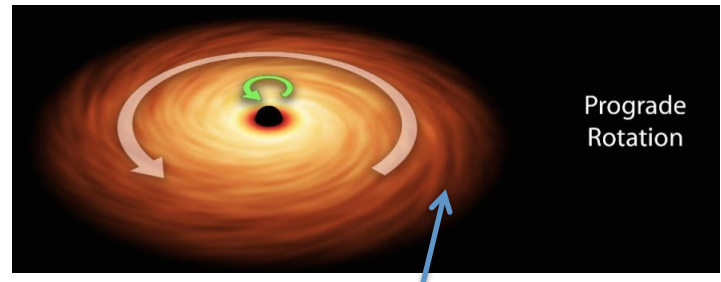
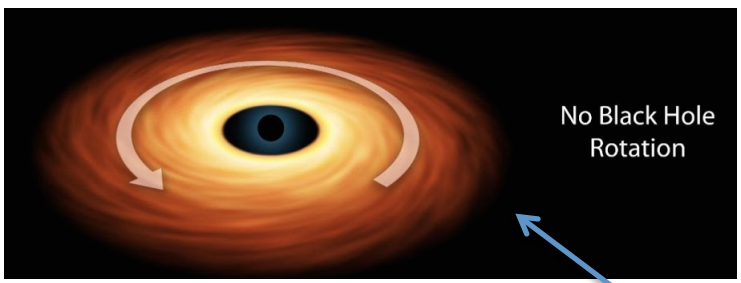
Cold filaments are repelled from the halo. Their structure is strongly perturbed



Dubois, Pichon et al., 2013

BH spin and its consequence on BH growth $spin = a = \frac{J_{BH}}{(GM_{BH}^2/c)}$

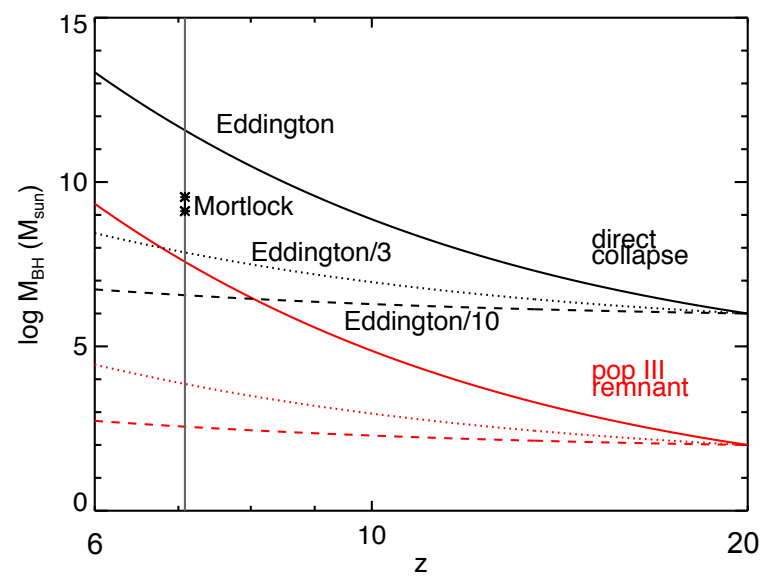
- Radiative efficiency depends on the BH spin parameter
- A non-spinning BH has a low radiative efficiency $e_r=0.057$. For a maximally spinning BH with $a=0.998$ $e_r=0.321$ ($e_r=0.038$ if $a=-0.998$)



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- The spin of BH is inherited from the history of gas accretion and successive mergers.
- Potential issue here: if BHs are maximally spinning then $t_{Edd}(a=0.998)=144$ Myr. Only possible to grow a $10^5 M_{sun}$ seed BH up to $10^8 M_{sun}$ in a Gyr ($z=6$)

$$t_{Edd} = \frac{M_{BH}}{\dot{M}_{Edd}} = \frac{\epsilon_r \sigma_{TC}}{4\pi G m_p}$$



BH spin evolution through accretion and binary coalescence (in a nutshell)

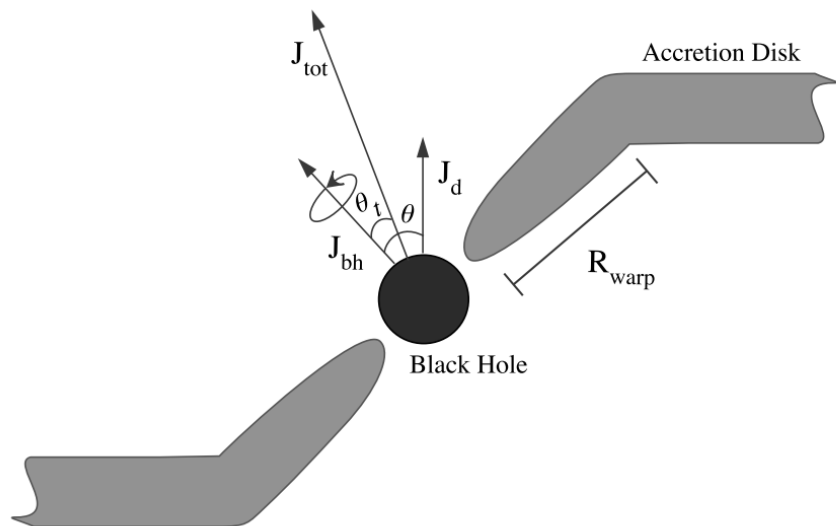
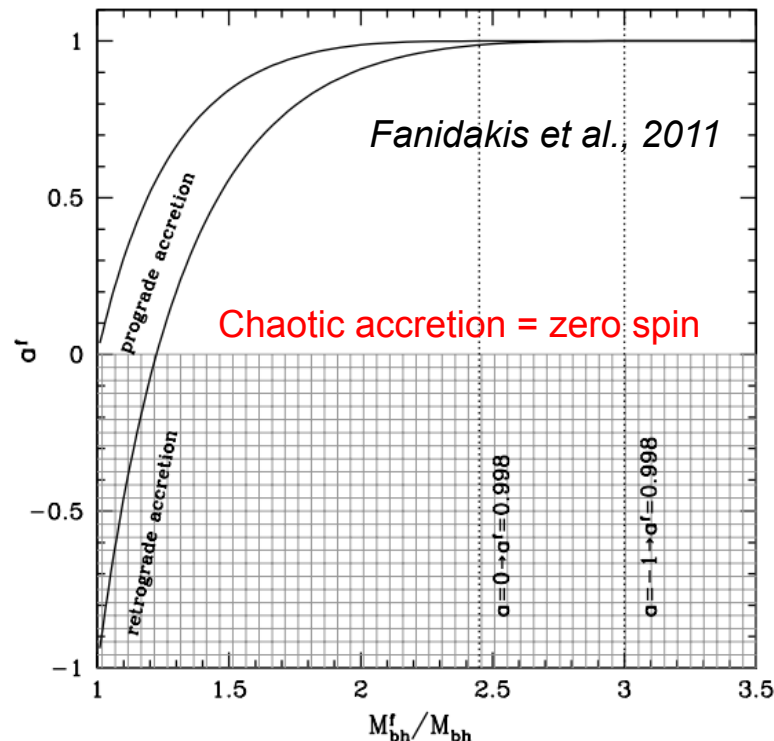


Figure 2. Schematic illustration of a warped accretion disc. \mathbf{J}_{bh} is the angular momentum of the BH ($|\mathbf{J}_{\text{bh}}| = |a|GM_{\text{bh}}^2/c$), \mathbf{J}_{d} is the angular momentum of the disc given by Eq. (20) and \mathbf{J}_{tot} represents the total angular momentum of the system, $\mathbf{J}_{\text{bh}} + \mathbf{J}_{\text{d}}$.

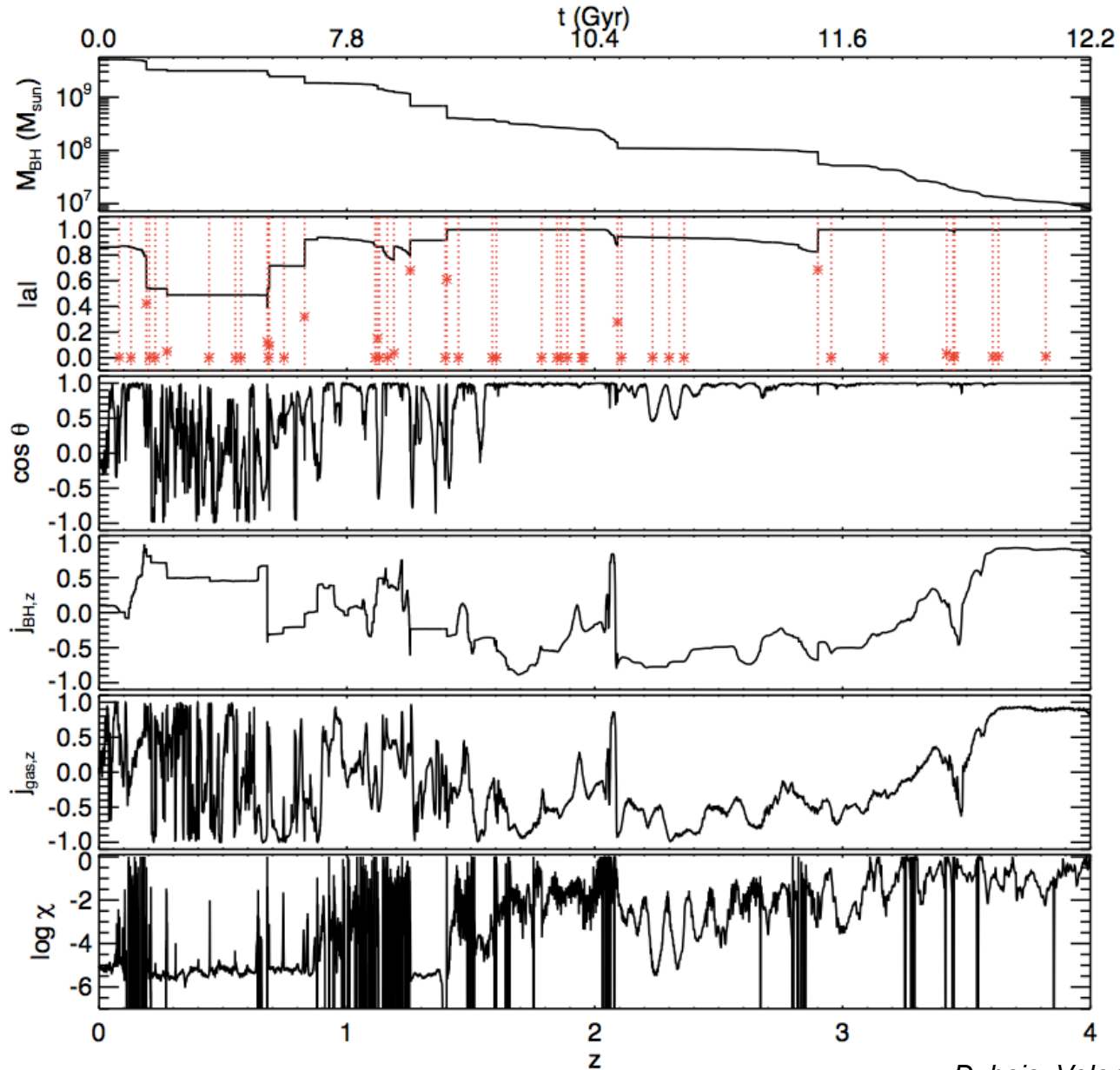


+

- BH-BH coalescence re-orientates the BH spin and change its amplitude if the mass ratio is large enough.
- The final value depends on the angle between the initial spin and orbital AM of the binary

Two types of evolution: 1- the typical SMBH in a BCG

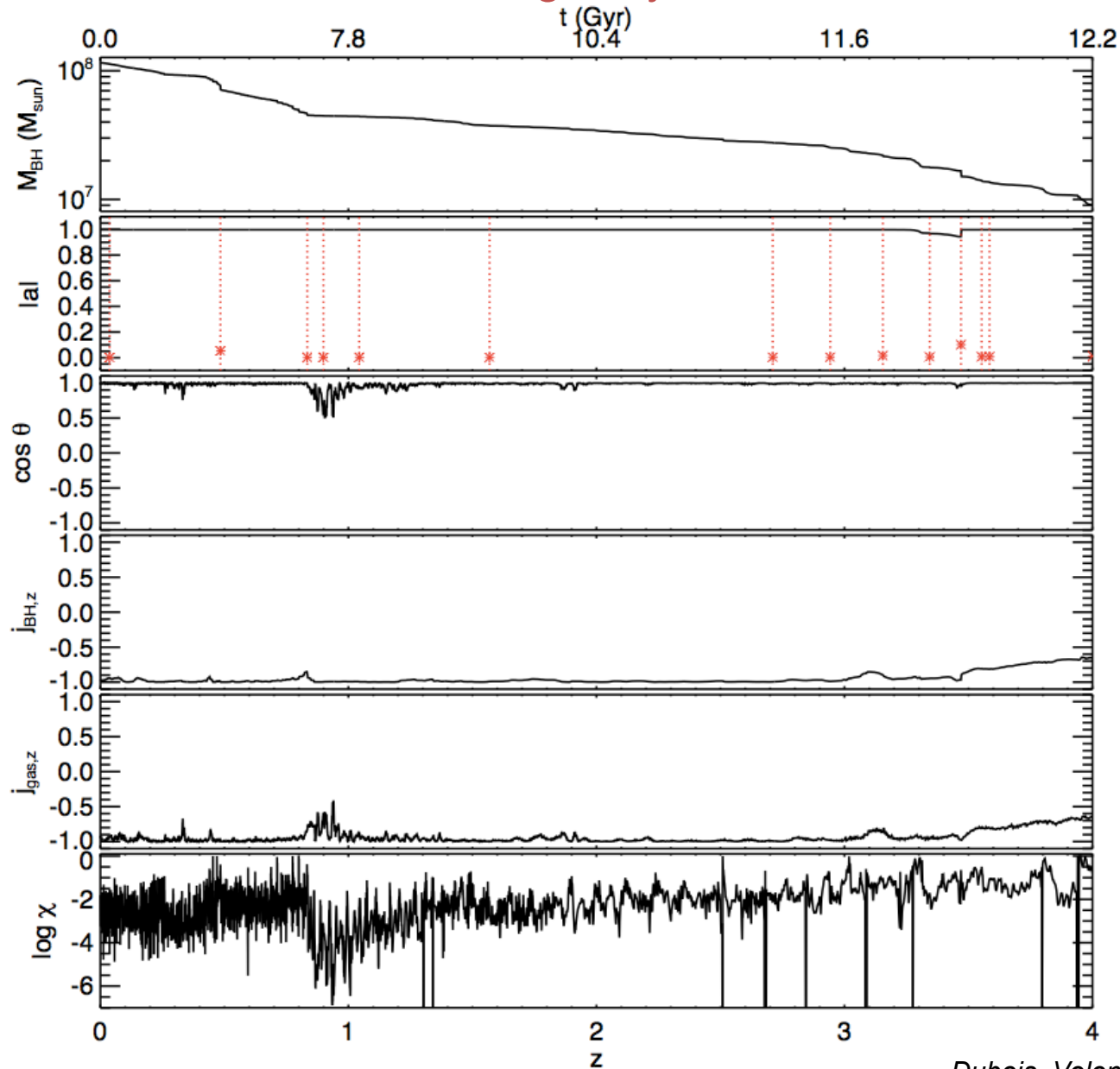
Large volume cosmological simulation
kpc resolution



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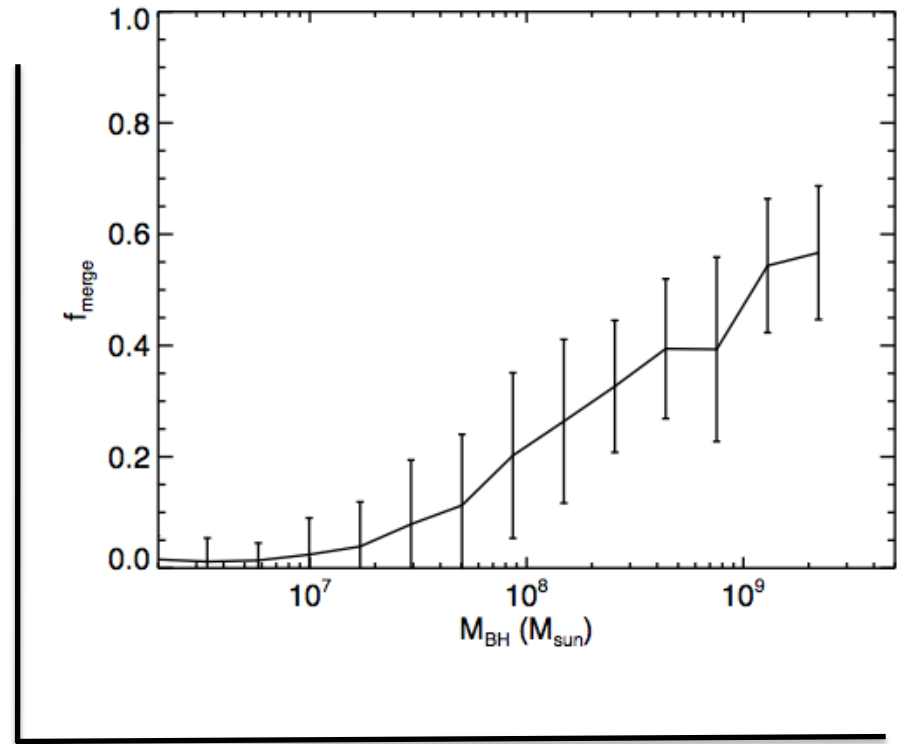
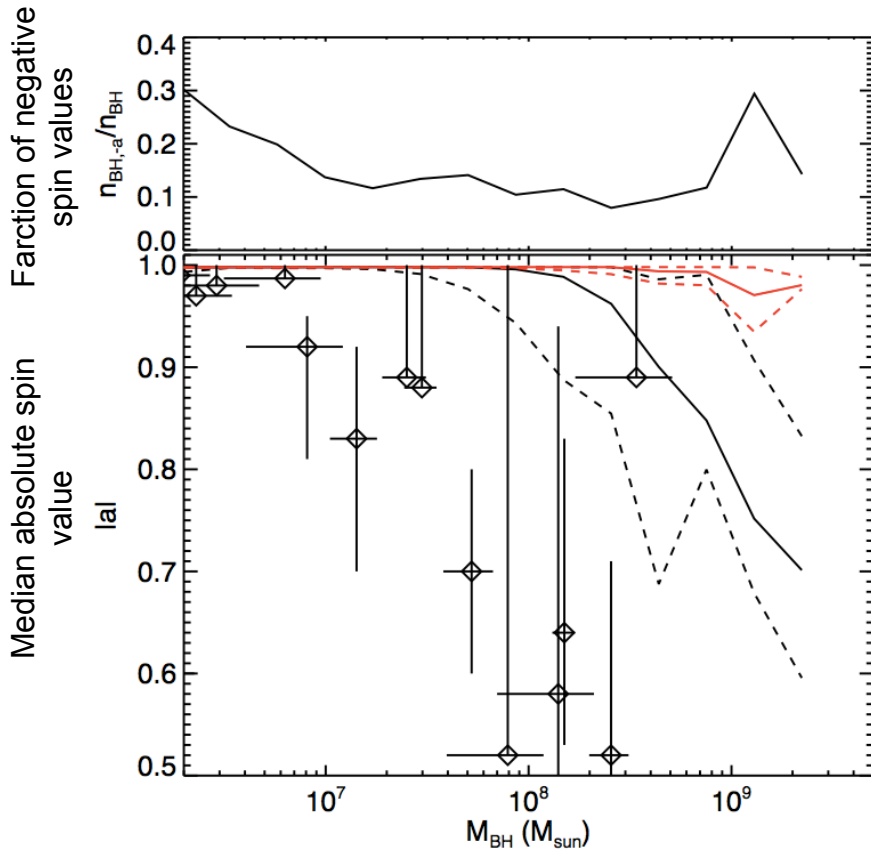
2- the typical SMBH in a massive disc galaxy

Large volume cosmological simulation
kpc resolution

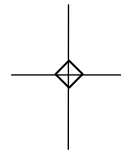


Spins versus Masses and the role of mergers

Observations



Data (Reynolds, 2013)



Numerical simulation:

Fiducial model



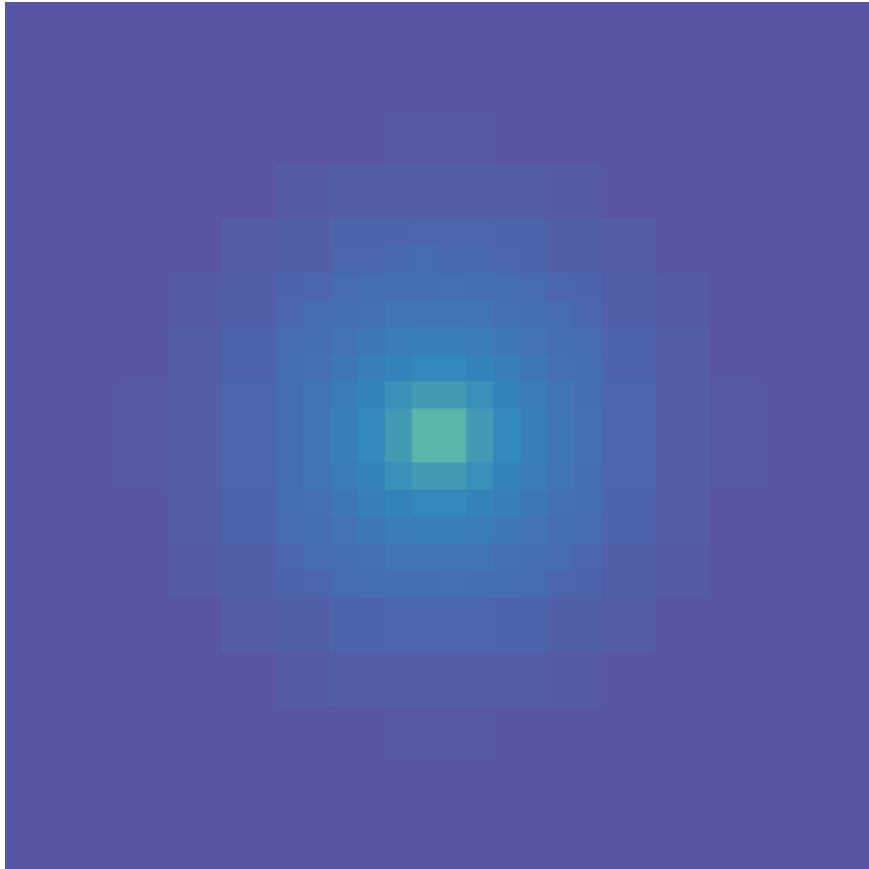
Model w/o mergers to change spins



Cranking up the resolution

Isolated halo $M_{\text{vir}} = 10^{12} M_{\text{sun}}$ & 10 pc resolution

Edge-on

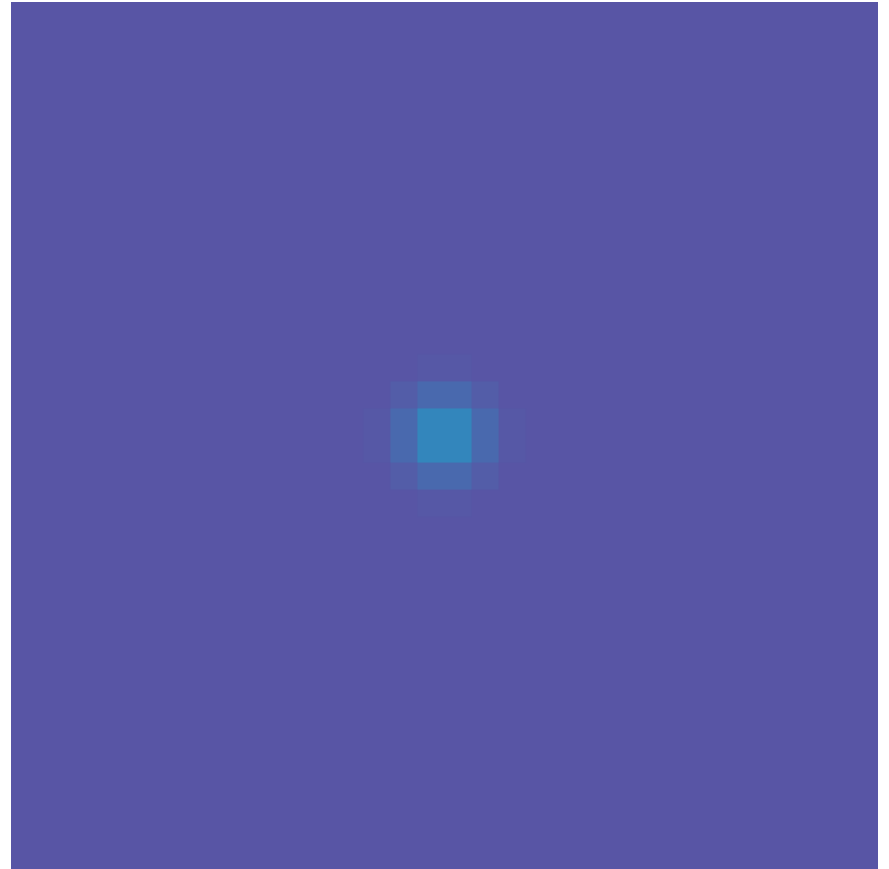


5 kpc

29/07/14

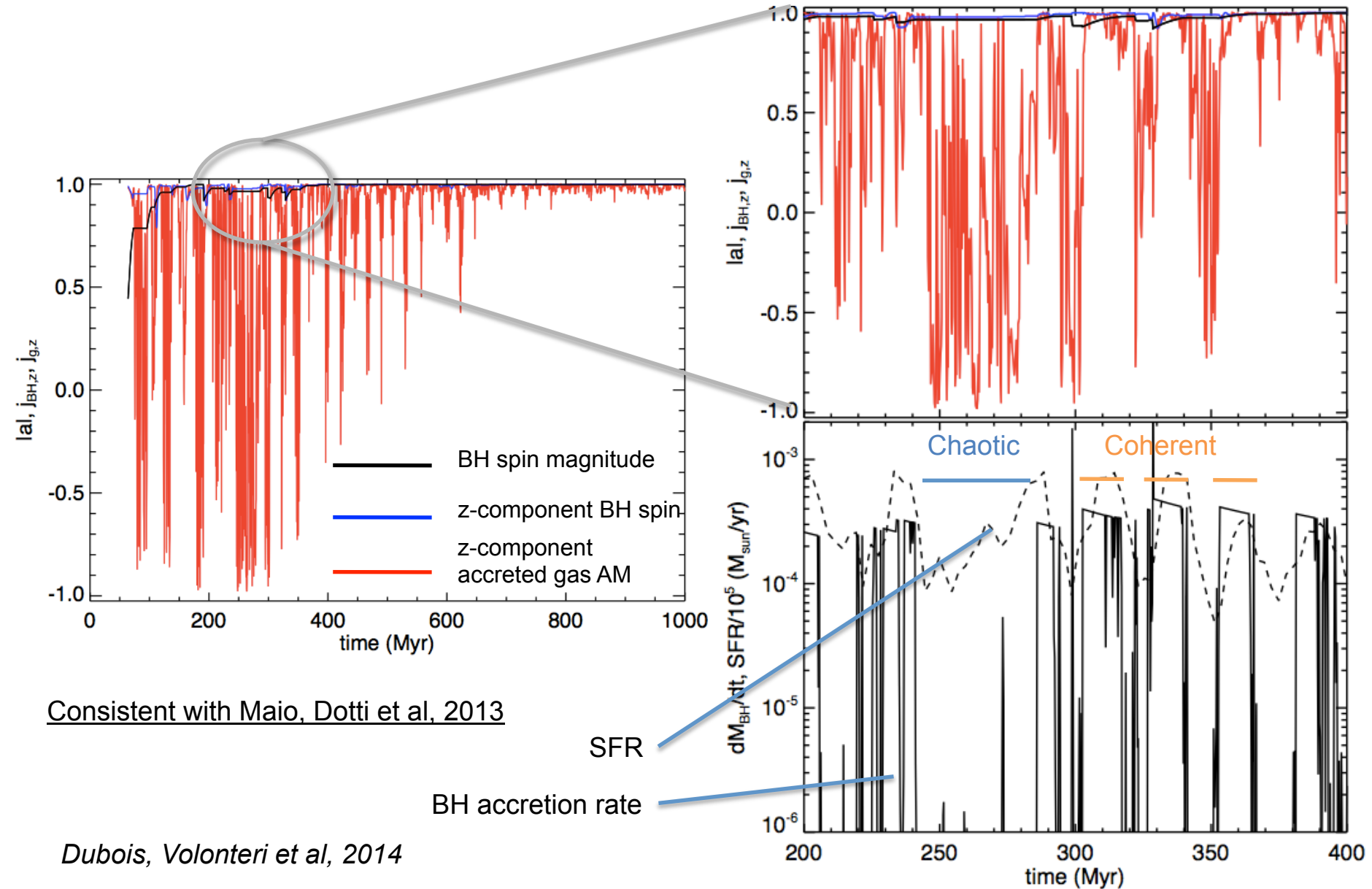
AGN vs SF

Face-on



5 kpc

Strong coherence of gas accretion

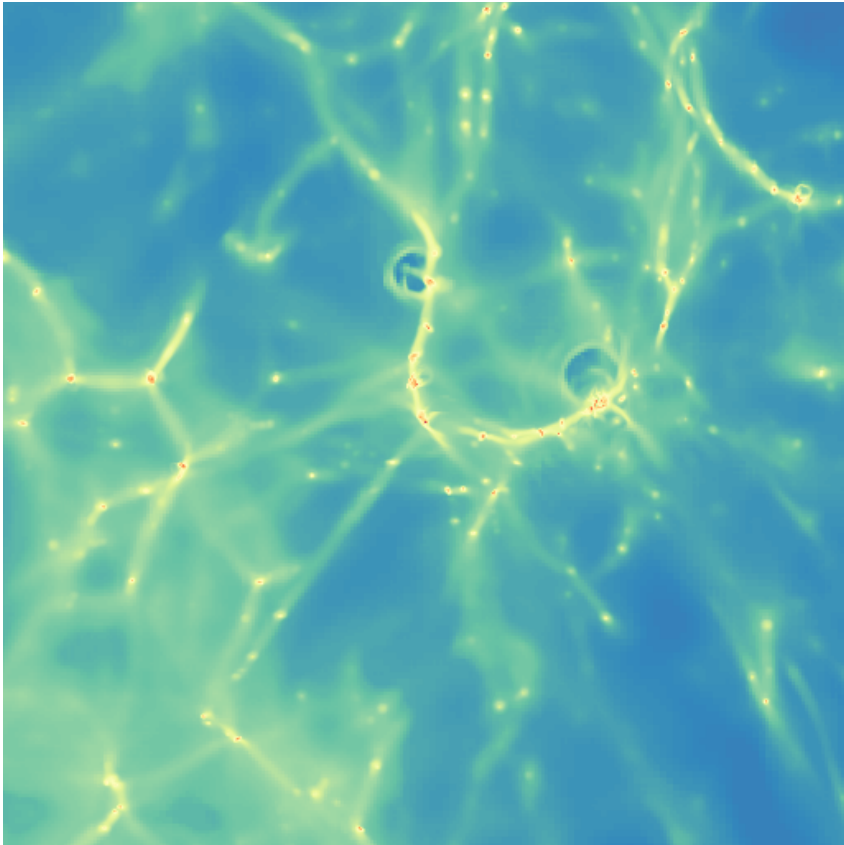


Consistent with Maio, Dotti et al, 2013

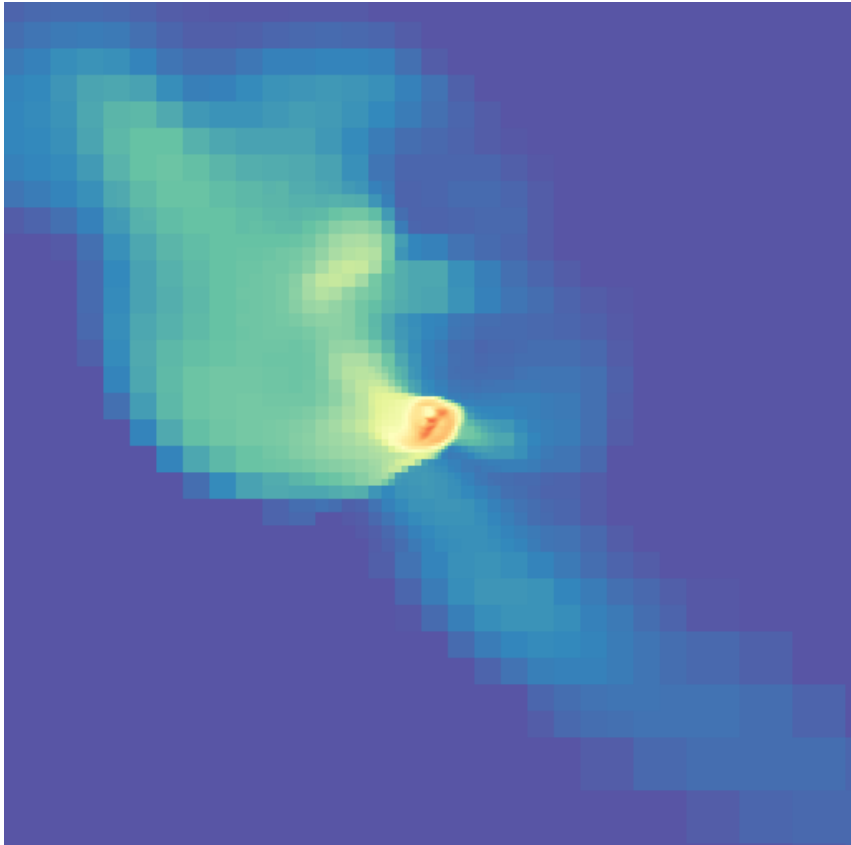
Dubois, Volonteri et al, 2014

Going cosmological

$M_{\text{vir}} = 10^{12} M_{\text{sun}}$ @ $z=2$ & 10 pc resolution

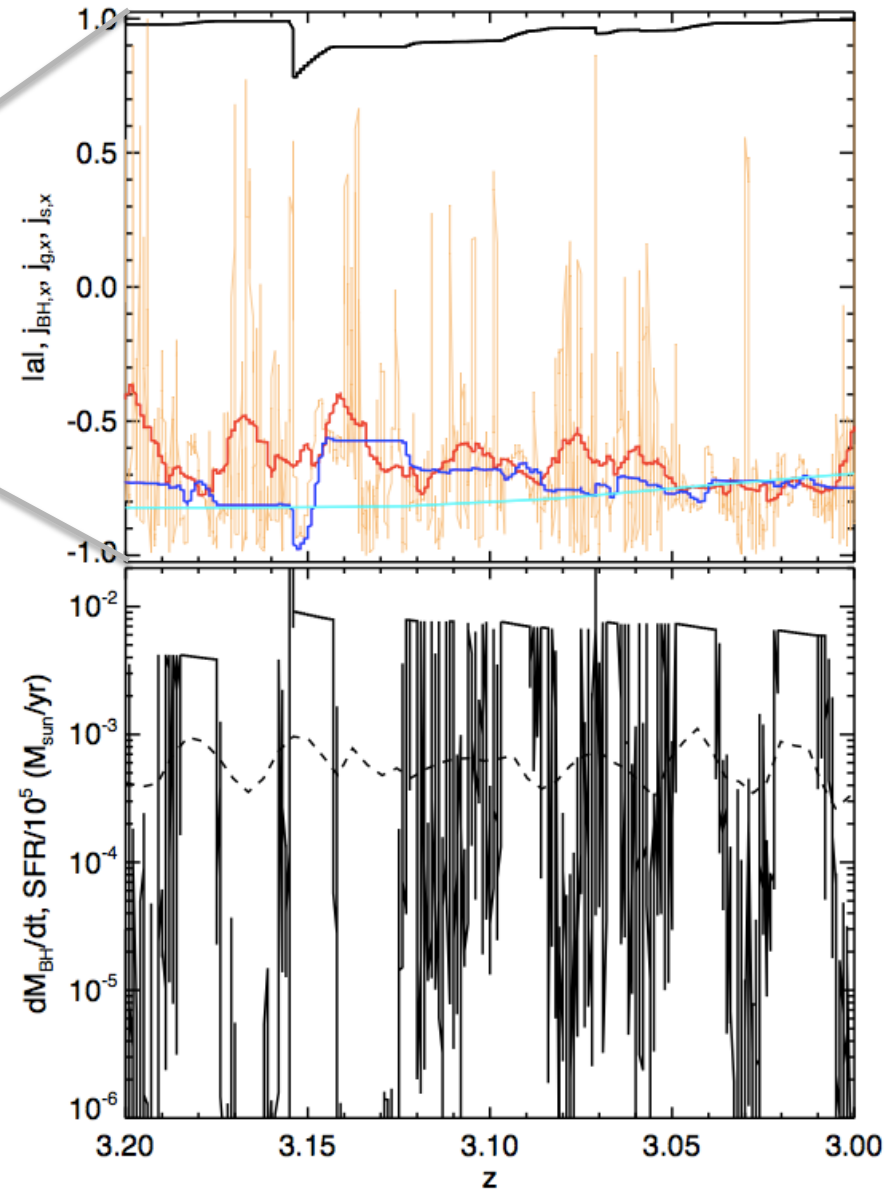
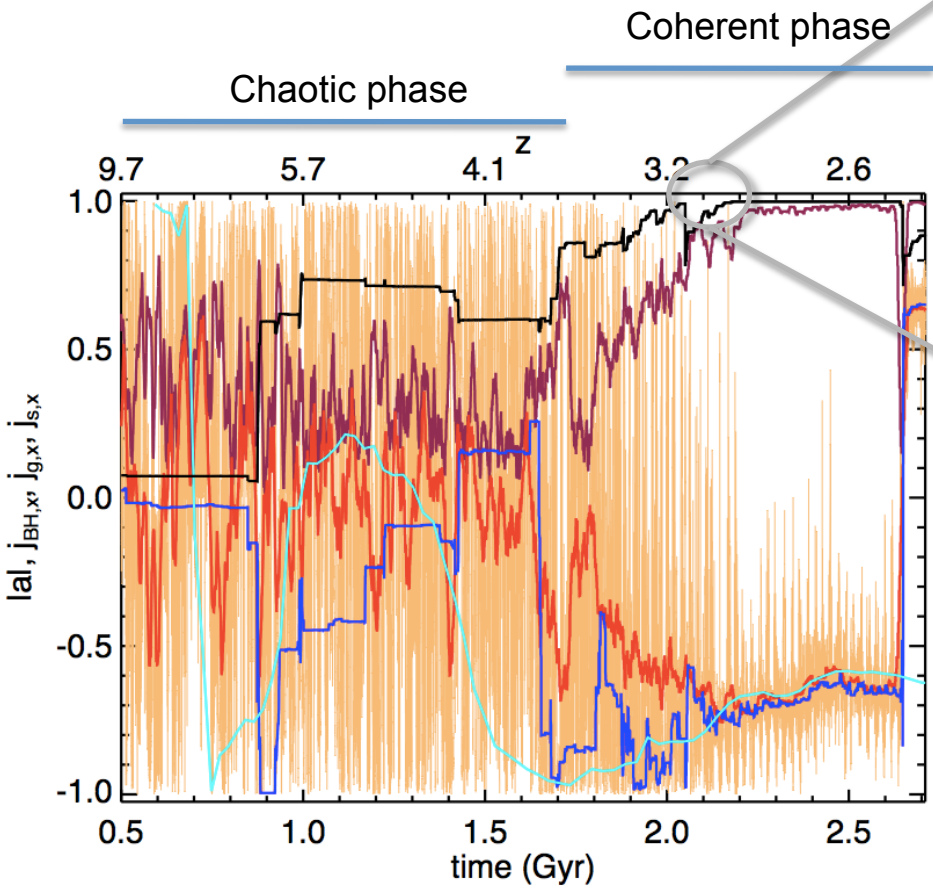


← 1 Mpc comoving →

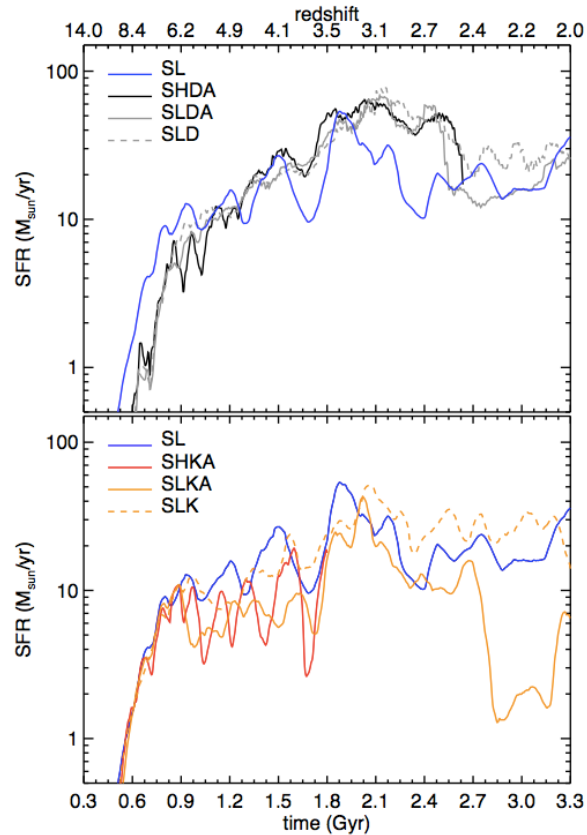
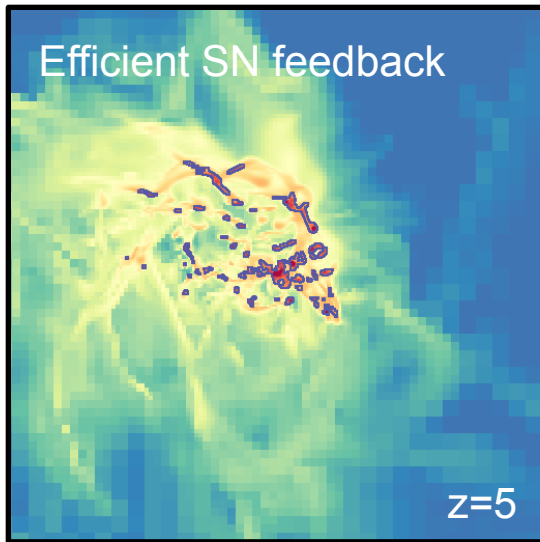
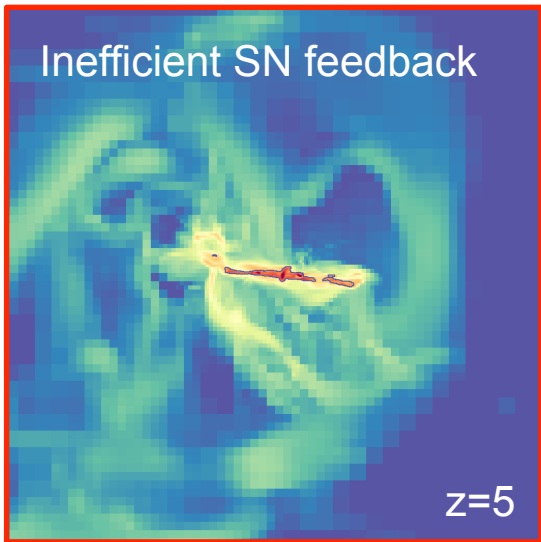


← 35 kpc comoving →

Going cosmological

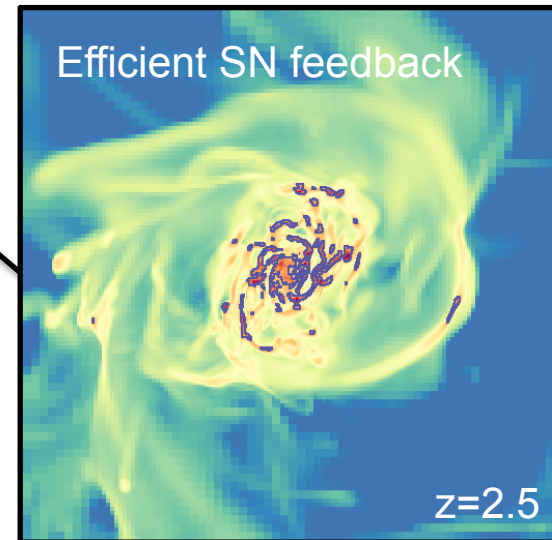
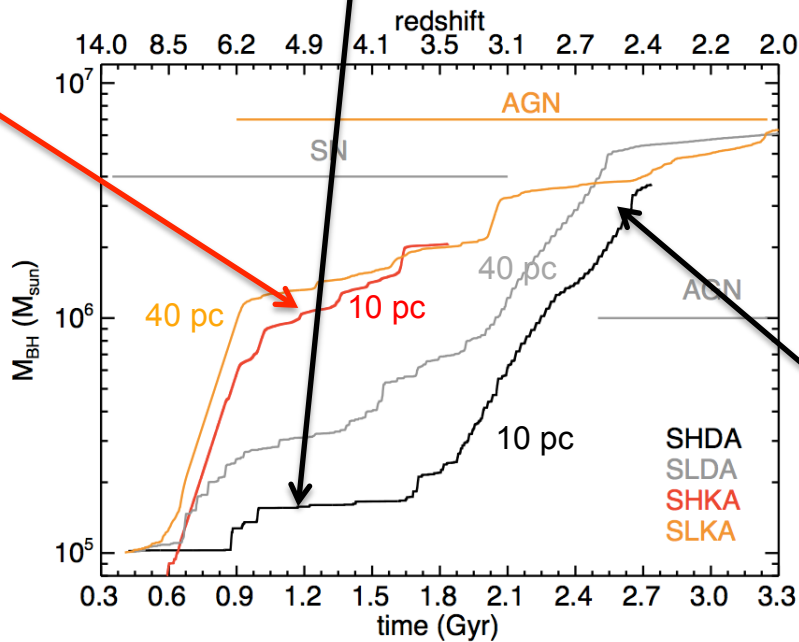


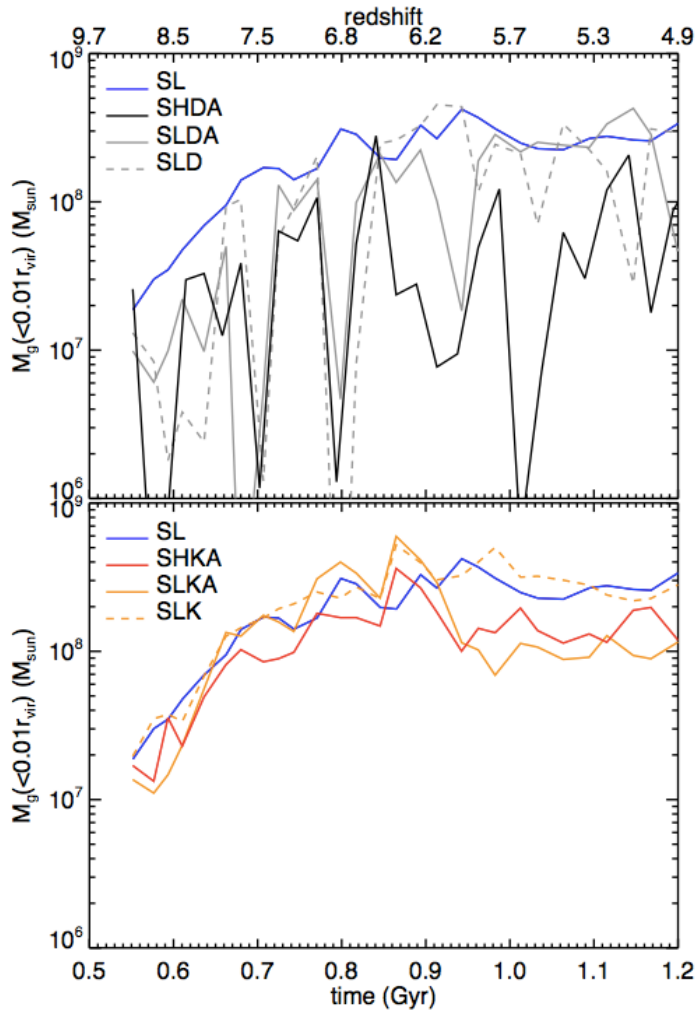
BH growth delayed by (too?) efficient SN feedback



“Inefficient”: kinetic blast wave model (Dubois & Teyssier, 2008)

“Efficient”: non-thermal component (CR, turbulence, magnetic fields) that delays gas cooling (Teyssier et al, 2013)





$$u_{\text{SN}} = 1.2 \sqrt{\frac{m_{\text{new,s}} \eta_{\text{SN}} e_{\text{SN}}}{m_{\text{g}}}}$$

$$\simeq 270 \sqrt{\frac{\eta_{\text{SN}}}{0.1}} \sqrt{\frac{(m_{\text{new,s}}/m_{\text{g}})}{0.1}} \text{ km s}^{-1}$$

$$u_{\text{esc}} = \sqrt{\frac{2Gm_{\text{cl}}}{r_{\text{cl}}}} \sim 300 \text{ km s}^{-1}$$

For $m_{\text{cl}} = 10^9 M_{\text{sun}}$ and $r_{\text{cl}} = 100 \text{ pc}$

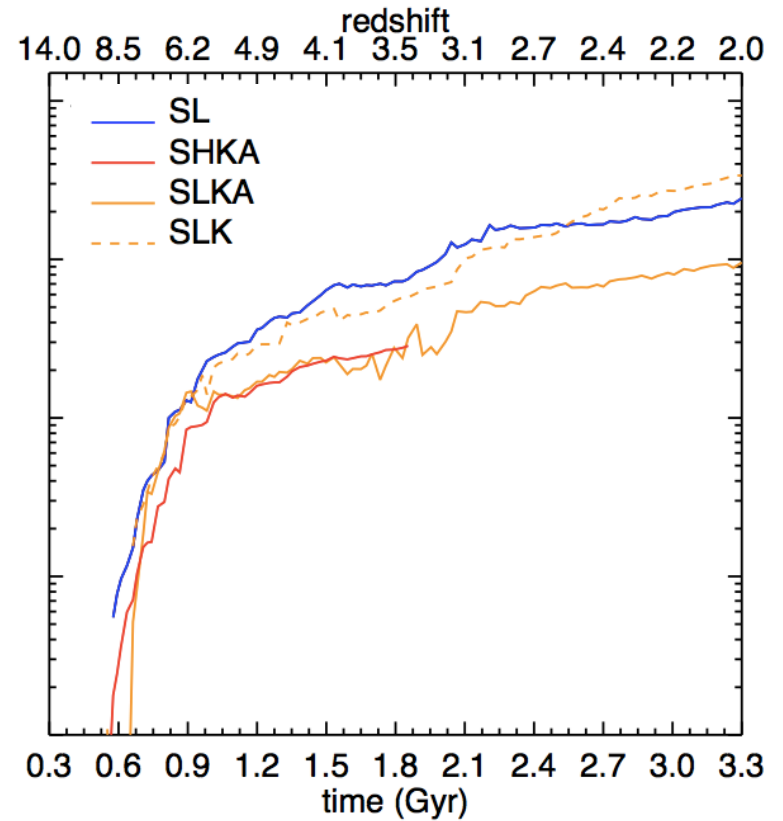
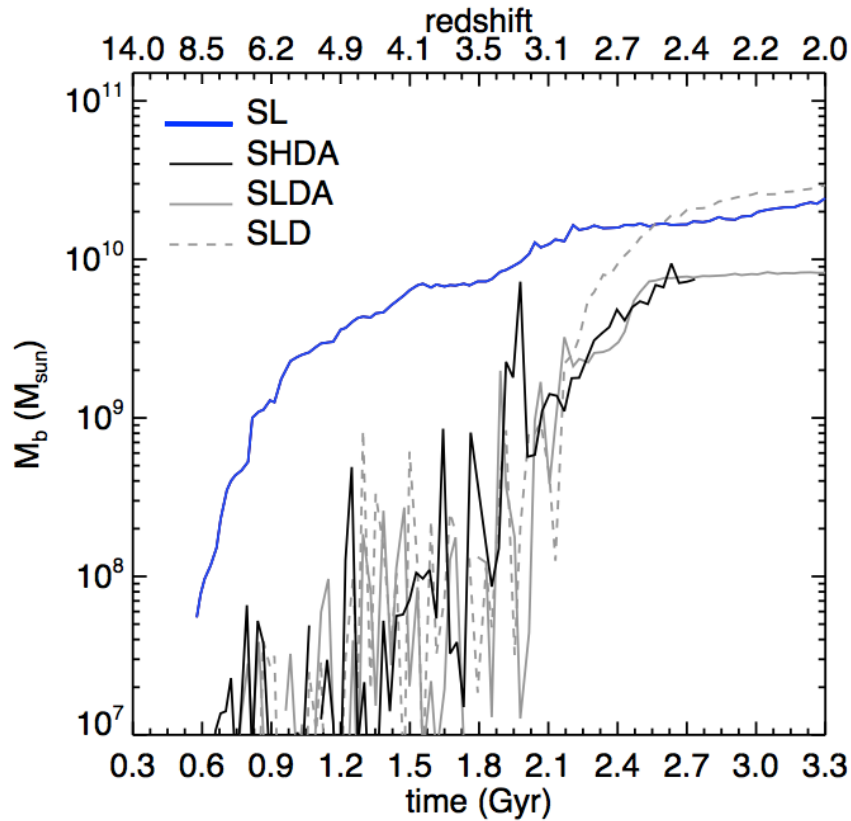
Dubois et al., in prep

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Dubois et al., in prep

Thank you for your attention

1 Mpc

$z = 38.305$