

# L2 Computational Physics

## Week 4 – 2<sup>nd</sup> order ODEs

# This week...

- 2<sup>nd</sup> order differential equations
  - Harmonic oscillator
  - Ballistics
  - Pendulum
- Black Box ODE solvers
- Phase space and Orbits

# 2<sup>nd</sup> order ODEs

# 2<sup>nd</sup> order ODEs

- 1<sup>st</sup> order (last week)

$$\frac{dx}{dt} = f(x, t)$$

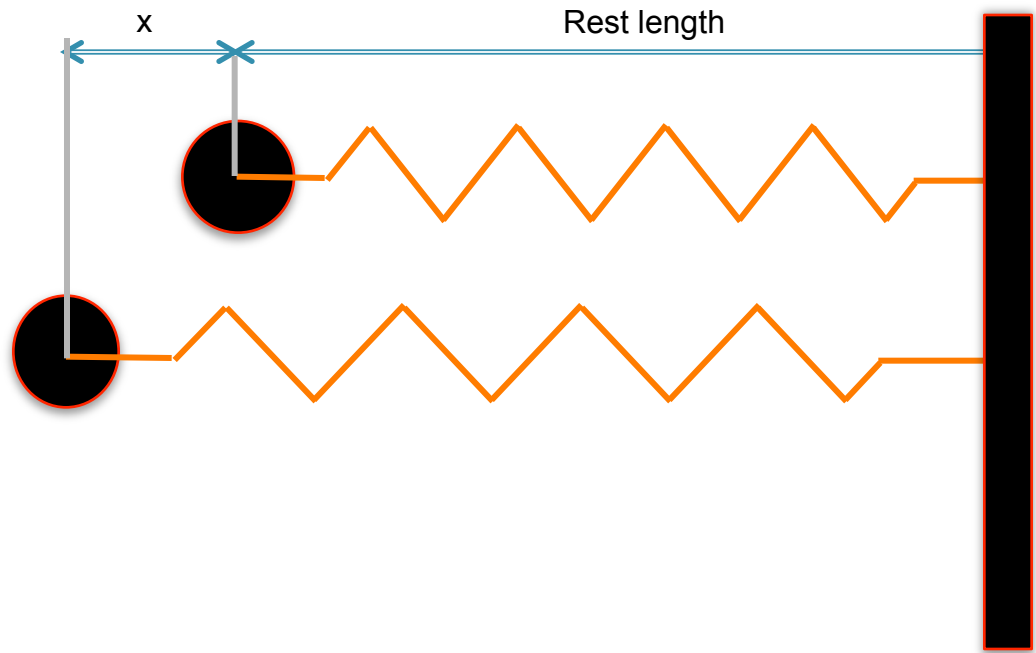
- 2<sup>nd</sup> order (this week)

$$\frac{d^2x}{dt^2} = f(x, t)$$

# Harmonic Oscillator

# Mass on a Spring

- Mass  $m$
- Spring constant  $k$

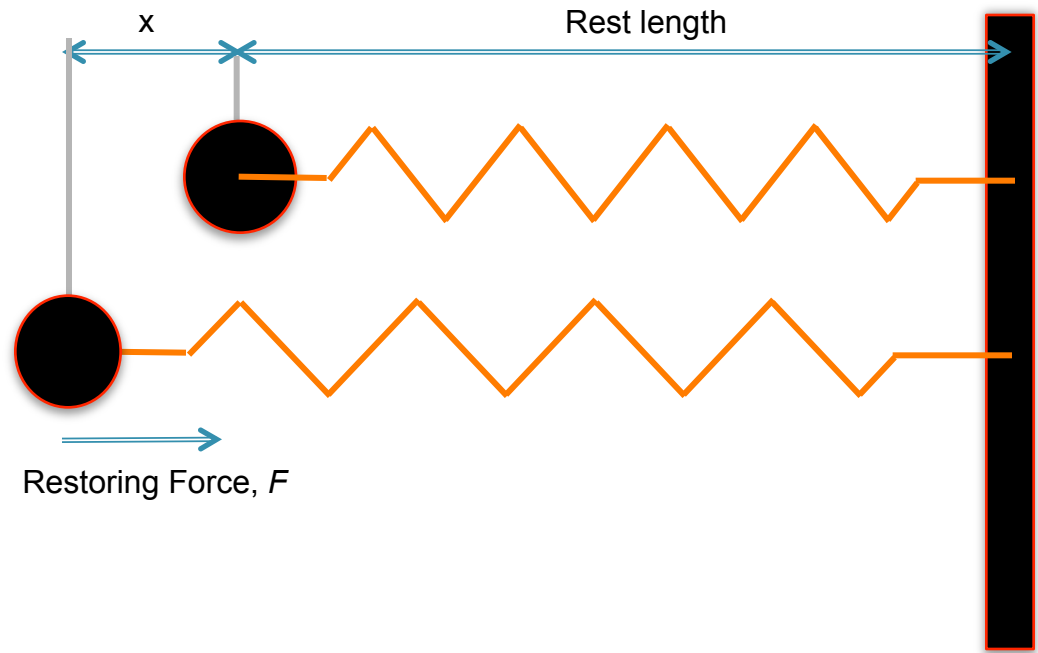


# Mass on a Spring

- Mass  $m$
- Spring constant  $k$
- Extension  $x$

$$F = -kx$$

$$\frac{d^2 x}{dt^2} = -\frac{k}{m} x$$



# Solve with Euler

- Euler, Heun, RK4 et al are for 1<sup>st</sup> order ODEs
- Rewrite our 2<sup>nd</sup> order DEQ in terms of 1<sup>st</sup> order equations
  - Introduce linear velocity,  $v$
- Now we have a pair of first order DEQs

$$\frac{d^2 x}{dt^2} = -kx$$



$$\begin{aligned}\frac{dx}{dt} &= v \\ \frac{dv}{dt} &= -\frac{k}{m}x\end{aligned}$$

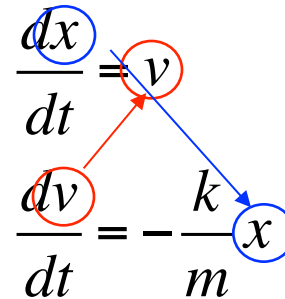


# Solve with Euler

- Euler, Heun, RK4 et al are for 1<sup>st</sup> order ODEs
- Rewrite our 2<sup>nd</sup> order DEQ in terms of 1<sup>st</sup> order equations
  - Introduce linear velocity,  $v$
- Now we have a pair of first order DEQs
  - The equations are ‘coupled’

$$\frac{d^2 x}{dt^2} = -kx$$

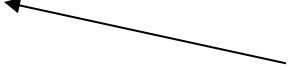


$$\begin{aligned} \frac{dx}{dt} &= v \\ \frac{dv}{dt} &= -\frac{k}{m}x \end{aligned}$$


# Solve with Euler

- Recap Euler:

$$x_{t+\Delta t} \approx x_t + f(x_t, t)\Delta t$$

$$f(x_t, t) \equiv \left. \frac{dx}{dt} \right|_{x_t, t}$$


# Solve with Euler

- Recap Euler:  $x_{t+\Delta t} \approx x_t + f(x_t, t)\Delta t$

- Harmonic Oscillator:  $\frac{dx}{dt} = v, \frac{dv}{dt} = -\frac{k}{m}x$

# Solve with Euler

- Recap Euler:  $x_{t+\Delta t} \approx x_t + f(x_t, t)\Delta t$

- Harmonic Oscillator:  $\frac{dx}{dt} = v, \frac{dv}{dt} = -\frac{k}{m}x$

$$x_{t1} = x_{t0} + v\Delta t$$

- Ta-da:

$$v_{t1} = v_{t0} - \frac{k}{m}x\Delta t$$

```
from __future__ import division
import numpy
import matplotlib.pyplot as pyplot
```

```
m=0.1
k= 0.4
```

**Model Parameters**

```
x0=0.3
v0=0.0
t0=0
```

**Initial conditions**

```
t1=35
n_panels=500
dt=(t1-t0)/n_panels
```

**timestep**

```
xs = numpy.zeros((n_panels),)
vs = numpy.zeros((n_panels),)
ts = numpy.zeros((n_panels),)
```

**arrays**

```
def dx_dt(x,v):
    return v
```

```
def dv_dt(x,v):
    return -(k/m)*x
```

**Differential equations!**

```
for i in range(n_panels-1):
    k0x=dx_dt(x0,v0)
    k0v=dv_dt(x0,v0)
    x1=x0+k0x(x0,v0)*dt
    v1=v0+k0v(x0,v0)*dt

    xs[i+1]=x1
    vs[i+1]=v1
    ts[i+1]=t0+i*dt
    x0=x1
    v0=v1
```

**Integrate via euler**

```
pyplot.figure(figsize=(8,8))
pyplot.subplot(211)
pyplot.plot(ts,xs,color='red')
pyplot.xlabel('time (s)')
pyplot.ylabel('position (m)')
pyplot.subplot(212)
pyplot.plot(xs,vs,color='red')
pyplot.xlabel('position (m)')
pyplot.ylabel('velocity (m/s)')

pyplot.show()
```

**Plot**

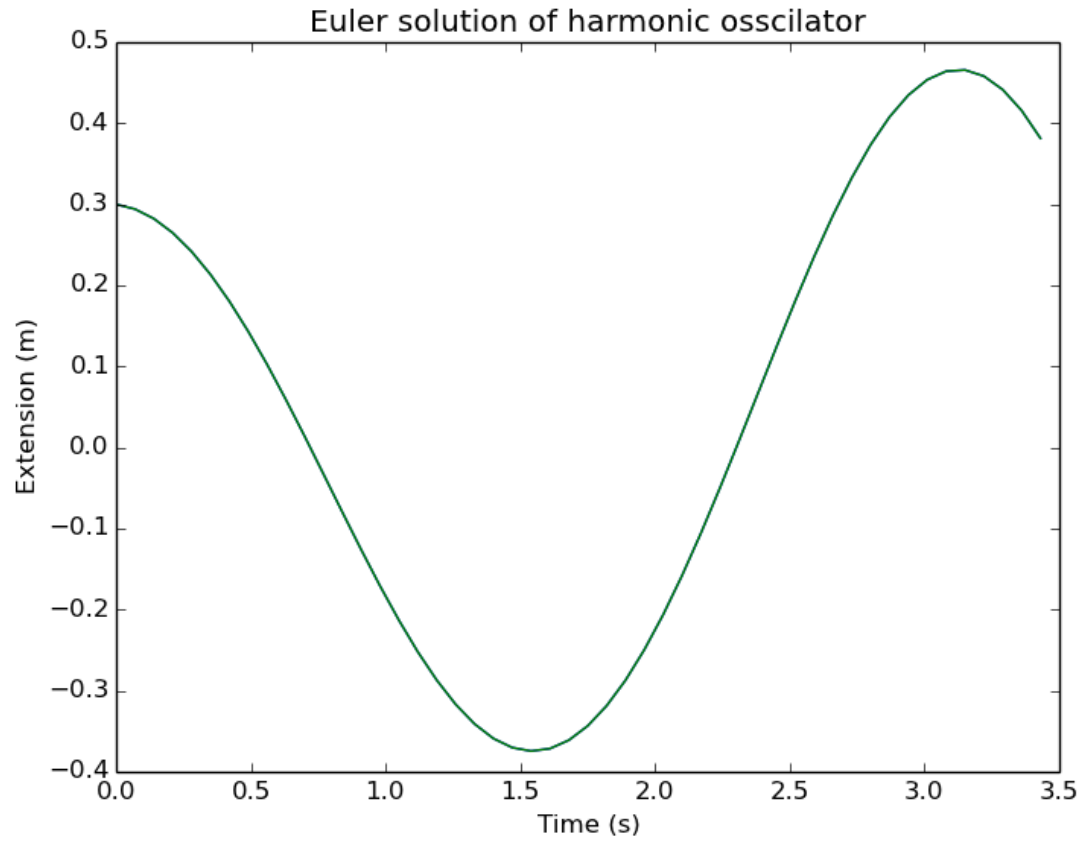
# Euler Timestep

- Within our code we have two states
  - Before the timestep
    - $x_0, v_0$
  - After the timestep
    - $x_1, v_1$
- Any simulation code (not just DEQs) has to carefully separate the present and the future
- ‘Tick Tock’

```
for i in range(n_panels-1):
    k0x=dx_dt(x0,v0)
    k0v=dv_dt(x0,v0)
    x1=x0+k0x(x0,v0)*dt
    v1=v0+k0v(x0,v0)*dt

    xs[i+1]=x1
    vs[i+1]=v1
    ts[i+1]=t0+i*dt
    x0=x1
    v0=v1
```

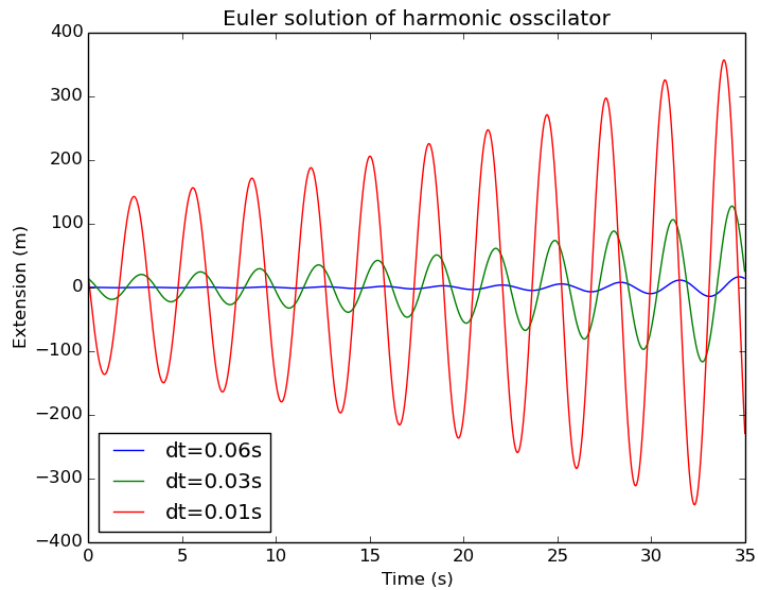
# Let's run it!



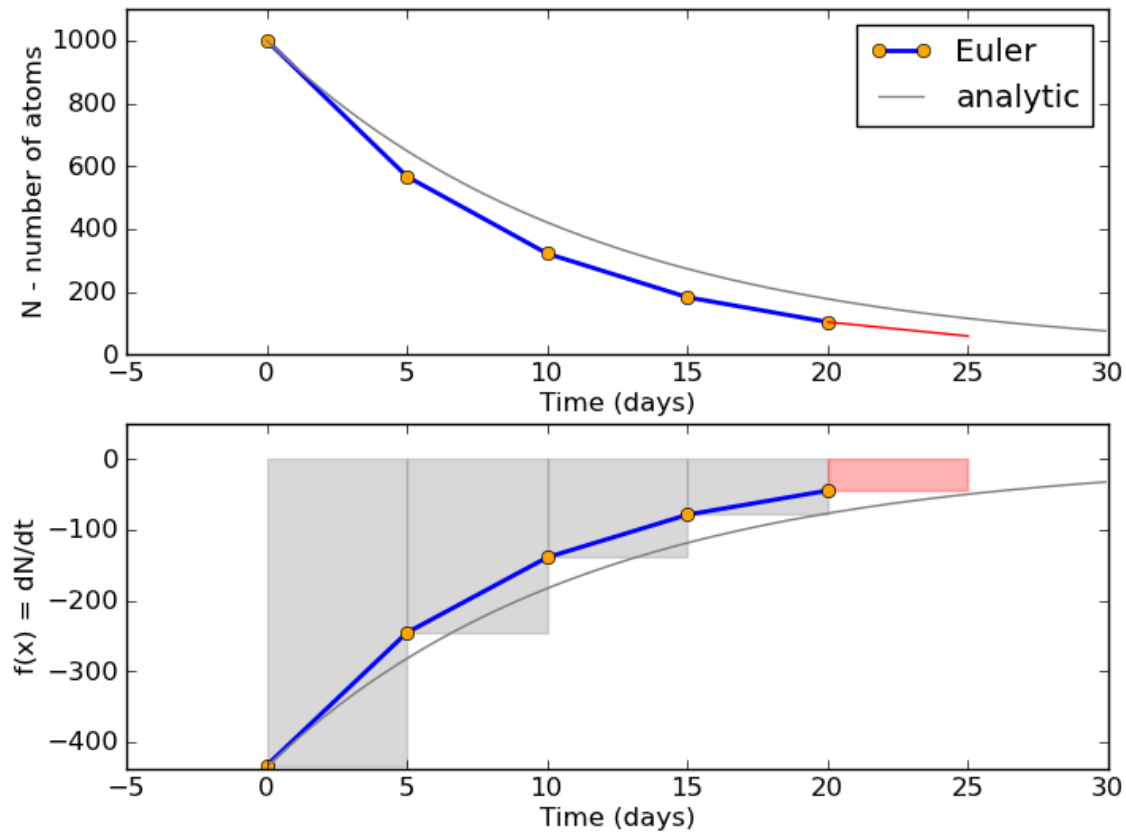


# Euler is failing us

- Euler method has an error in it

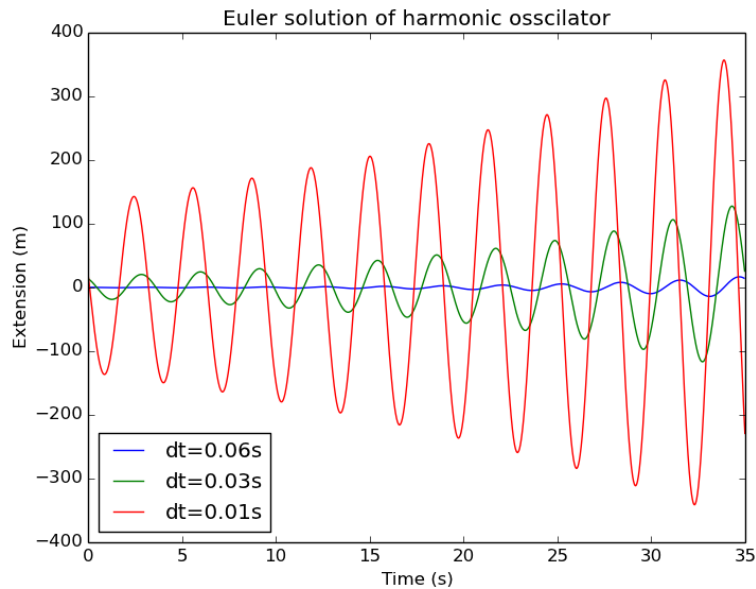


# Flashback – nuclear decay



# Euler is failing us

- Euler method has an error in it
- The error accumulates over time
  - Problematic for decay
  - Catastrophic for periodic systems!



# Energy Conservation

- **Physics 101 – Energy is conserved**
- The error in the Euler timestep violates conservation of energy
- There's an error in total energy in each timestep:

$$0.5 \frac{k}{m} (kx^2 + mv^2) (\Delta t)^2$$

- We need a better method!

# Energy Conservation

There's an error total energy in each timestep:

$$0.5 \frac{k}{m} (kx^2 + mv^2) (\Delta t)^2$$

- The error itself is always positive and is proportional to KE ( $v^2$ ) and PE ( $x^2$ )
- Over one full cycle of the pendulum:

$$E(t + \Delta t) = E(t) + \alpha E(t)$$

$$E(t + \Delta t) = (1 + \alpha)E(t)$$

# Energy Conservation

There's an error total energy in each timestep:

$$0.5 \frac{k}{m} (kx^2 + mv^2) (\Delta t)^2$$

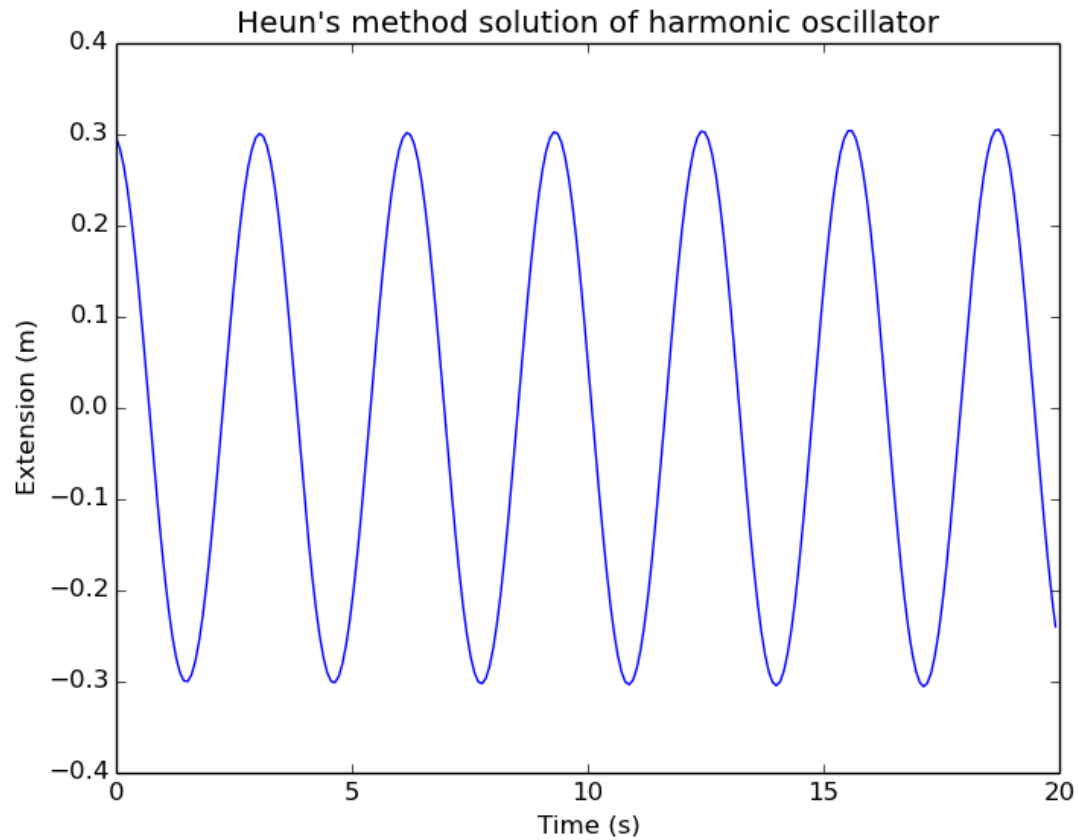
- The error itself is always positive and is proportional to KE ( $v^2$ ) and PE ( $x^2$ )
- Over one full cycle of the pendulum:

$$E(t + \Delta t) = E(t) + \alpha E(t)$$

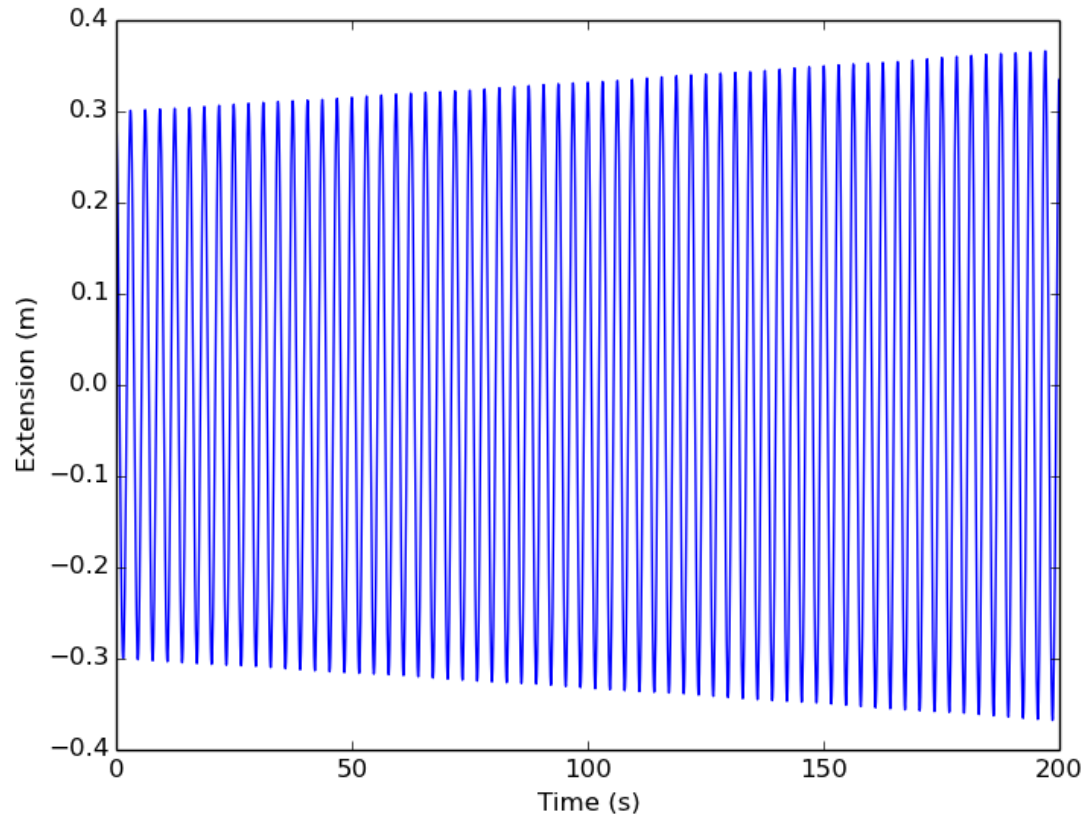
$$E(t + \Delta t) = (1 + \alpha)E(t)$$

**Exponential growth!**

# Let's try Heun

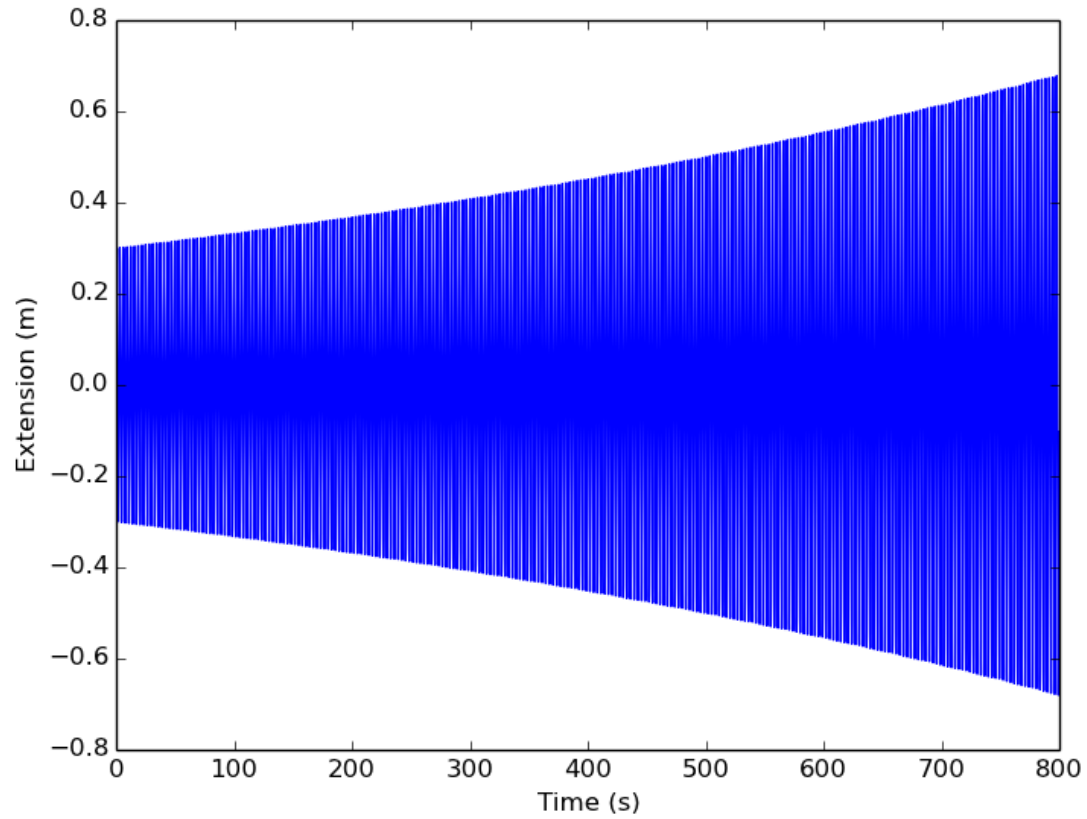


# Close...





... but no cigar



EnergyCons\_Euler

IP[y]: Notebook EnergyCons\_Euler

File Edit View Insert Cell Kernel Help

Code Cell Toolbar: None

```
In [1]: from __future__ import division
import sympy
from IPython.display import display
sympy.init_printing()
```

```
In [2]: v0 = sympy.Symbol('v_0') # Velocity at start of timeste
x0 = sympy.Symbol('x_0') # Extension at start of timestep
dt = sympy.Symbol('dt') # Timestep
k = sympy.Symbol('k') # Spring constant
m = sympy.Symbol('m') # Mass
```

```
In [3]: v1 = v0 - (k/m) * x0 * dt # Velocity after timestep
x1 = x0 + v0 * dt # Extension after timestep
```

```
In [4]: def spring_energy(vel, extension):
return 0.5*m*vel**2 + 0.5*k*extension**2
```

```
In [5]: e0 = spring_energy(v0, x0)
e1 = spring_energy(v1, x1)
display(e0)
display(e1)
```

$$0.5kx_0^2 + 0.5mv_0^2$$

$$0.5k(dt v_0 + x_0)^2 + 0.5m\left(-\frac{dk}{m}x_0 + v_0\right)^2$$

```
In [7]: |(e1-e0)
```

$$\text{Out}[7]: -0.5kx_0^2 + 0.5k(dt v_0 + x_0)^2 - 0.5mv_0^2 + 0.5m\left(-\frac{dk}{m}x_0 + v_0\right)^2$$

```
In [8]: (e1 - e0).expand().simplify()
```

$$\text{Out}[8]: \frac{0.5k}{m} dt^2 (kx_0^2 + mv_0^2)$$

EnergyCons\_Heun

IP[y]: Notebook EnergyCons\_Heun

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Code Cell Toolbar: None

```
In [10]: # k0 = dx/dt at t0=
k0_v = -(k/m) * x0
k0_x = v0
# x' = euler prediction
vp = v0 + k0_v * dt
xp = x0 + k0_x * dt
# k1 = dx/dt at euler prediction
k1_v = -(k/m) * xp
k1_x = vp
# Averaged rate of change
k_v = (k0_v + k1_v)/2
k_x = (k0_x + k1_x)/2
# x1
v1 = v0 + k_v * dt
x1 = x0 + k_x * dt
#
v1 = v1.simplify()
x1 = x1.simplify()
```

```
In [11]: display(v1)
display(x1)
```

$$\frac{1}{m} \left( -\frac{dk}{2} (dt v_0 + 2x_0) + mv_0 \right)$$

$$-\frac{dt^2 kx_0}{2m} + dt v_0 + x_0$$

```
In [12]: e0 = spring_energy(v0, x0)
e1 = spring_energy(v1, x1)
(e1-e0).expand().simplify()
```

$$\text{Out}[12]: \frac{0.125dt^4}{m^2} k^2 (kx_0^2 + mv_0^2)$$

Euler-Cromer

## Euler

$$v_{t1} = v_{t0} - \frac{k}{m} x_0 \Delta t$$

$$x_{t1} = x_{t0} + v_0 \Delta t$$

## Euler-Cromer

$$v_{t1} = v_{t0} - \frac{k}{m} x_0 \Delta t$$

$$x_{t1} = x_{t0} + v_1 \Delta t$$

## Euler

$$v_{t1} = v_{t0} - \frac{k}{m} x_0 \Delta t$$

$$x_{t1} = x_{t0} + v_0 \Delta t$$

## Euler-Cromer

$$v_{t1} = v_{t0} - \frac{k}{m} x_0 \Delta t$$

$$x_{t1} = x_{t0} + v_1 \Delta t$$

Euler-Cromer reorders the time-step slightly.

This change leads to energy conservation

There is still an error, but it cancels over one full cycle

Very useful for studying periodic systems

```
for i in range(n_panels-1):
    k0x=dx_dt(x0,v0)
    k0v=dv_dt(x0,v0)
    v1=v0+k0v(x0,v0)*dt
    x1=x0+k0x(x0,v1)*dt

    xs[i+1]=x1
    vs[i+1]=v1
    ts[i+1]=t0+i*dt
    x0=x1
    v0=v1
```

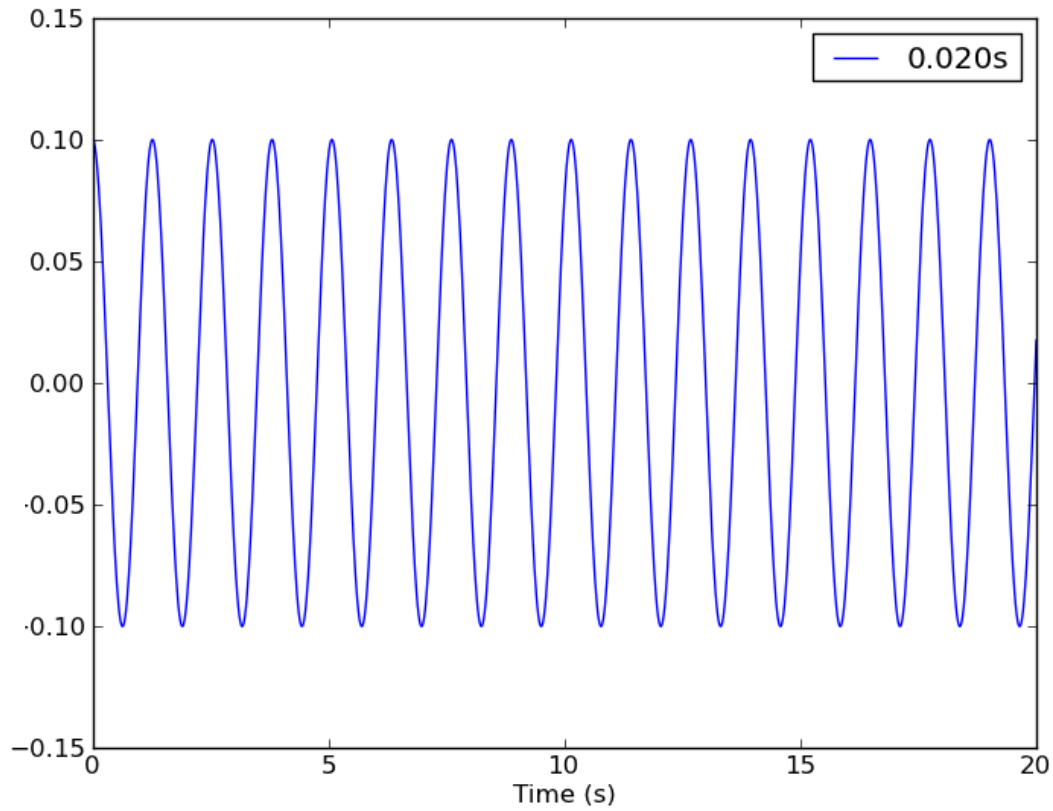
**Integrate via euler –  
cromer!!**

```
pyplot.figure(figsize=(8,8))
pyplot.subplot(211)
pyplot.plot(ts,xs,color='red')
pyplot.xlabel('time (s)')
pyplot.ylabel('position (m)')
pyplot.subplot(212)
pyplot.plot(xs,vs,color='red')
pyplot.xlabel('position (m)')
pyplot.ylabel('velocity (m/s)')

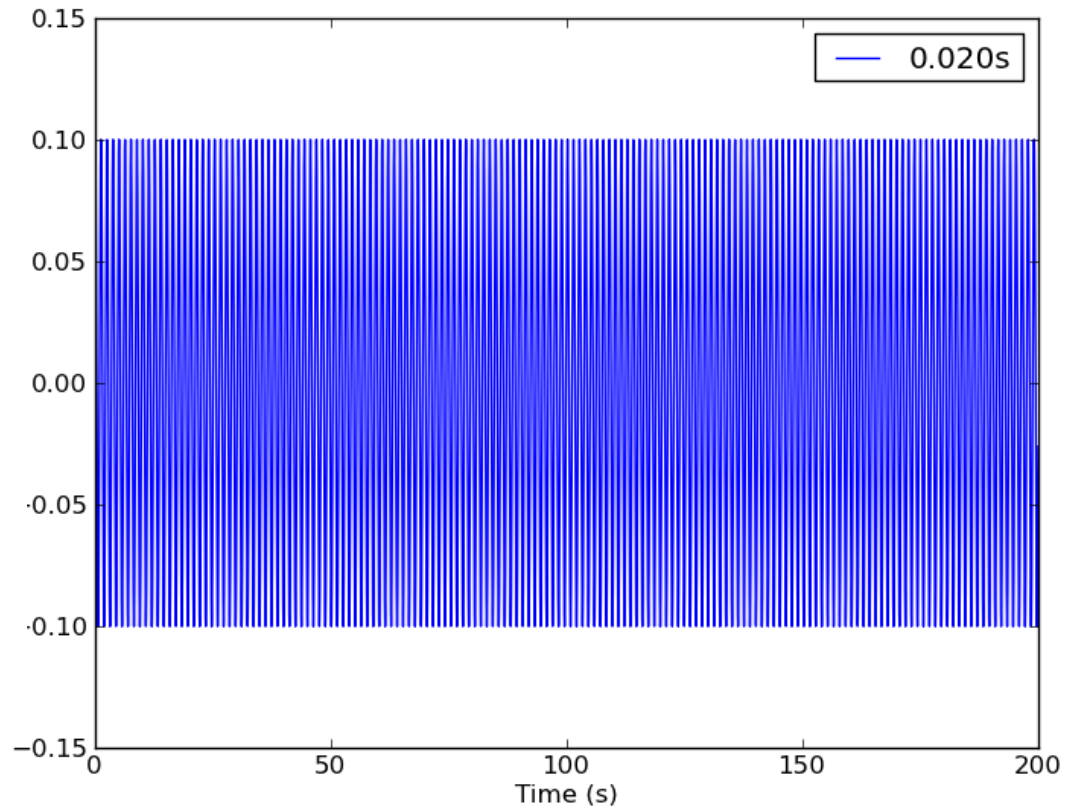
pyplot.show()
```

**Plot**

# Euler-Cromer

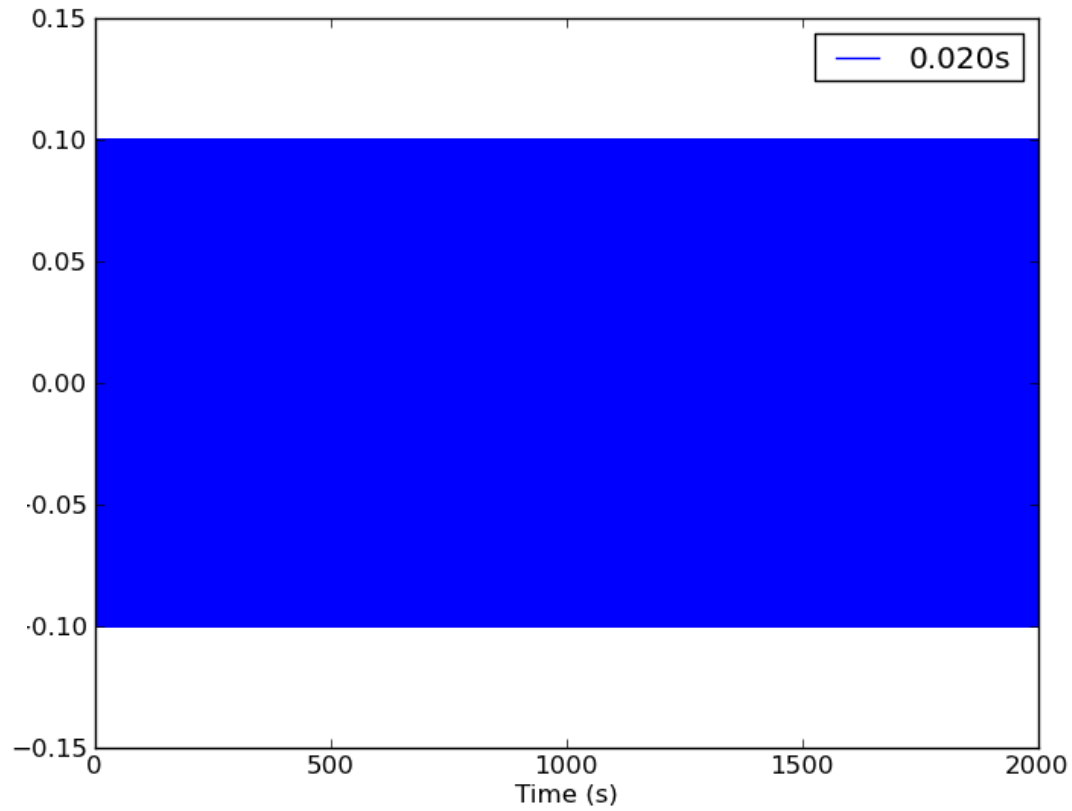


# Euler-Cromer

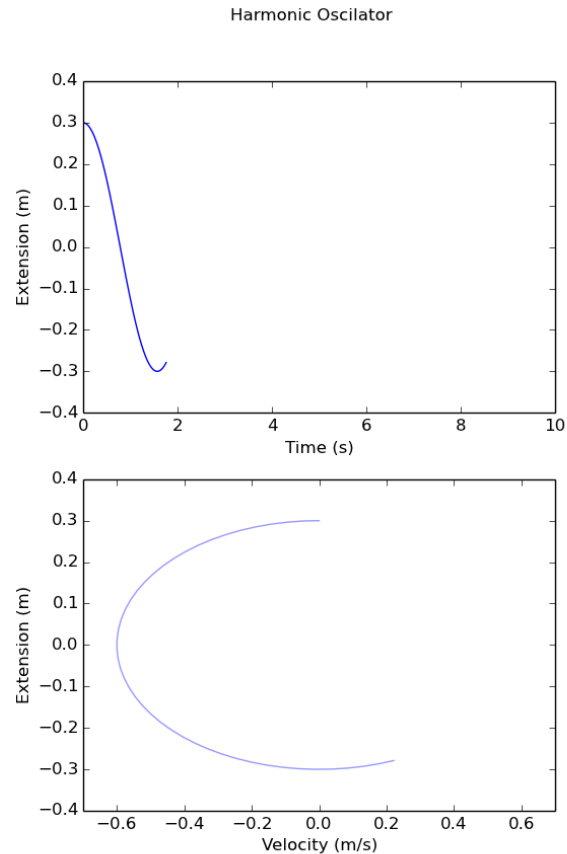




# Euler-Cromer



# Phase space and orbits...



(movie on DUO)

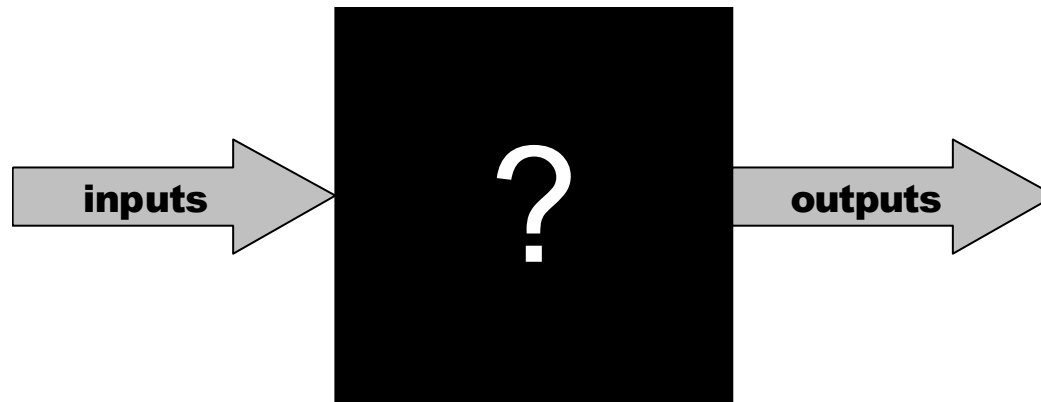
# Moral of the story

- Numerical methods are imprecise
  - Errors introduced
- Can distort results for a damped or bounded system
- Can be catastrophic for a periodic system

# Black Box ODE solver

# Black Box

- As with plain integration, lots of black box solvers for DEQs exist
- `scipy.integrate.odeint`



# scipy.integrate.odeint

```
Python Shell
File Edit Shell Debug Options Windows Help
IDLE 1.2.4
>>> import scipy.integrate
>>> help(scipy.integrate.odeint)
Help on function odeint in module scipy.integrate.odepack:

odeint(func, y0, t, args=(), Dfun=None, col_deriv=0, full_output=0, ml=None,
mu=None, rtol=None, atol=None, tcrit=None, h0=0.0, hmax=0.0, hmin=0.0, ixpr
=0, mxstep=0, mxhnil=0, mxordn=12, mxords=5, printmessg=0)
    Integrate a system of ordinary differential equations.

    Solve a system of ordinary differential equations using lsoda from the
    FORTRAN library odepack.

    Solves the initial value problem for stiff or non-stiff systems
    of first order ode-s::

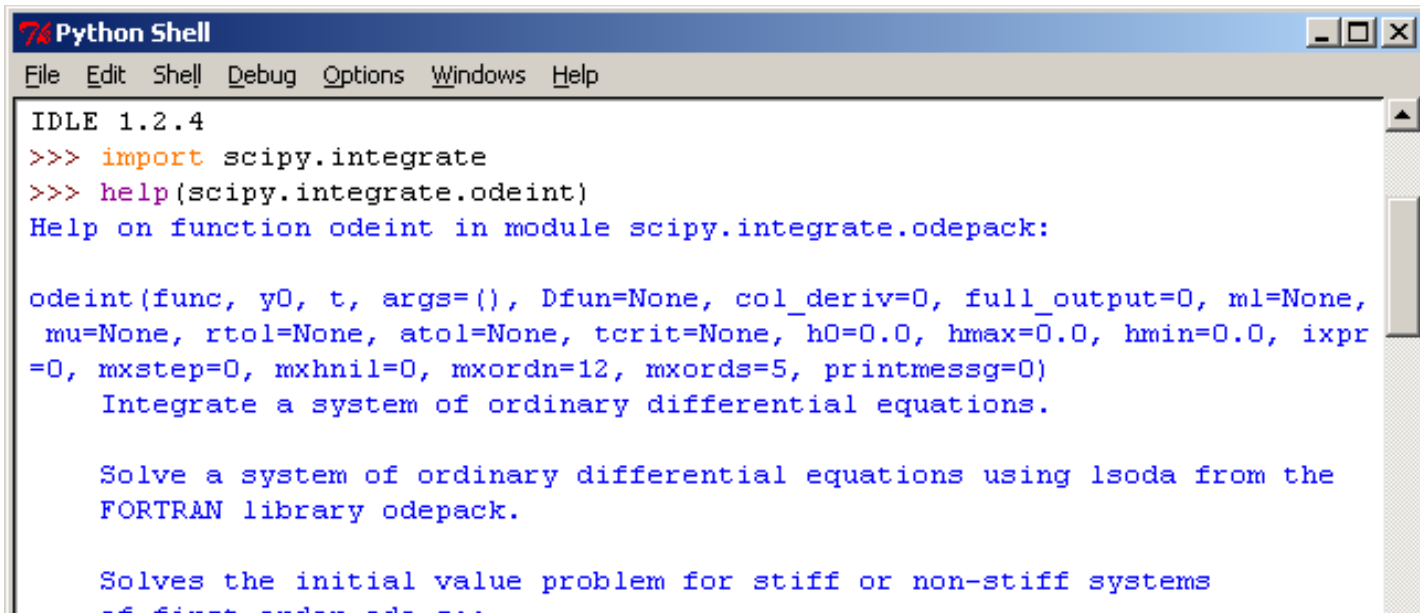
        dy/dt = func(y,t0,...)

    where y can be a vector.

Parameters
-----
func : callable(y, t0, ...)
    Computes the derivative of y at t0.
y0 : array
```

Ln: 35 Col: 0

# scipy.integrate.odeint



```
Python Shell
File Edit Shell Debug Options Windows Help
IDLE 1.2.4
>>> import scipy.integrate
>>> help(scipy.integrate.odeint)
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    Integrate a system of ordinary differential equations.

    Solve a system of ordinary differential equations using lsoda from the
    FORTRAN library odepack.

    Solves the initial value problem for stiff or non-stiff systems
    of first order ode...
```

**Wow, that's a lot of options**

**It's a lot more than for `scipy.integrate.quadrature` (week 2)**

**Numerically solving ODEs is a very complex field!**

# Using odeint for the harmonic oscillator

- We need to follow a series of steps
  1. Represent the system state  $(x,v)$  as a vector
  2. Make a function that takes a system and returns a vector of  $(dx/dt, dv/dt)$
  3. Make a series of timepoints for odeint to solve
  4. Invoke odeint
- Let's look at points 1-2 first



# Code tidy-up

Using arrays

```
xs = numpy.zeros((n_panels),)
vs = numpy.zeros((n_panels),)

ts = numpy.linspace(t0,t1,n_panels)
```

```
xs_ode=numpy.zeros((n_panels),)
vs_ode=numpy.zeros((n_panels),)
```

**arrays**

```
# required format for ODEINT
def f((x,v),time):
    dx_dt=v
    dv_dt=-(k/m)*x
    return numpy.array((dx_dt,dv_dt))
```

**Call with a pair of x,v values – pass in as array for a single time value**

**Differential equations  
- combined**

```
# state is a single pair of x,v values
state=numpy.array((x0,v0))
```

```
values=scipy.integrate.odeint(f, state, ts)
```

```
xs_ode=values[:,0]
vs_ode=values[:,1]
```

```
#compare to euler
for i in range(n_panels-1):
    k0x,k0v=f( (x0,v0),ts[i])
    xs_eul[i+1]=x0+k0x*dt
    vs_eul[i+1]=v0+k0v*dt
    x0=xs_eul[i+1]
    v0=vs_eul[i+1]
```

**Integrate**

**Via python**

**Returns 2D array**

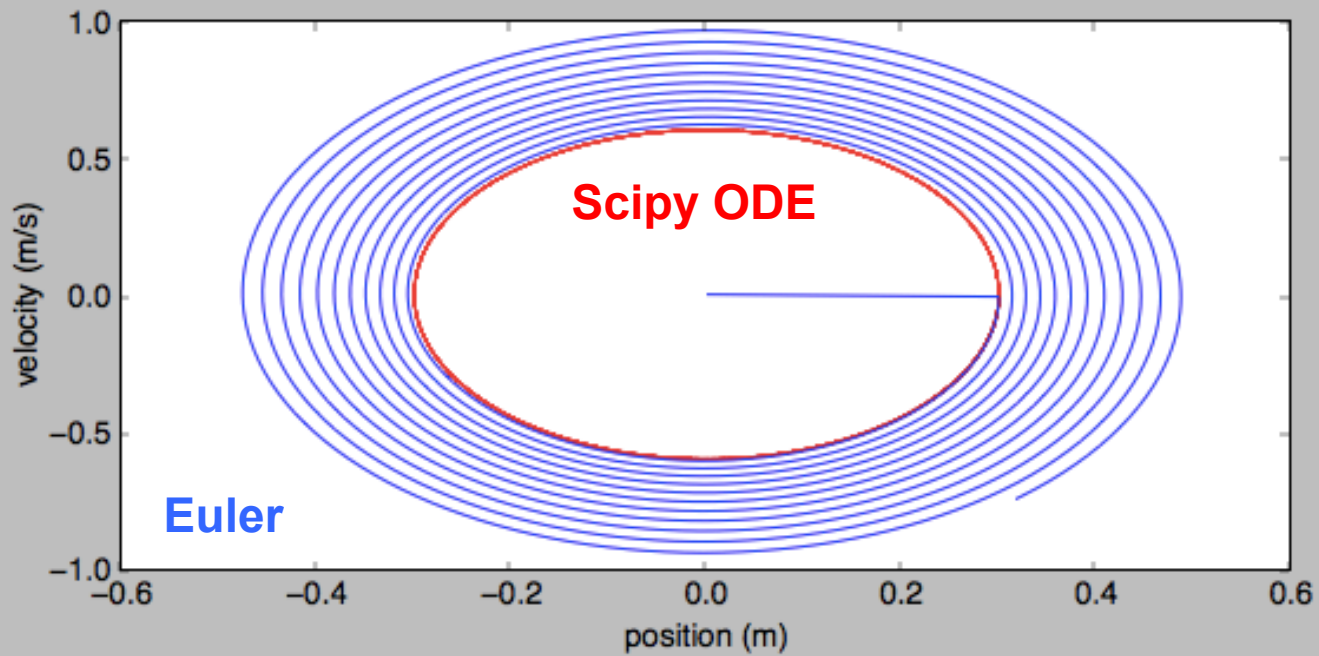
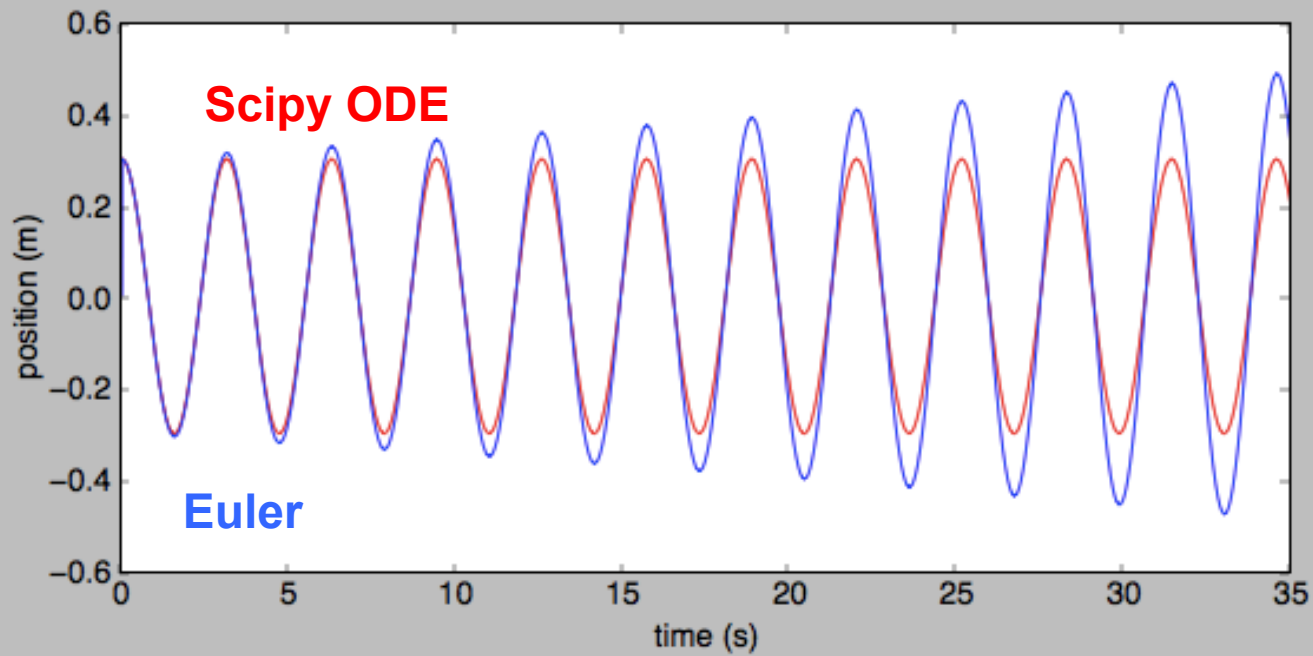
**So split into x and v**

```
pyplot.figure(figsize=(8,8))
pyplot.subplot(211)
pyplot.plot(ts,xs_ode,color='red')
pyplot.plot(ts,xs_eul,color='blue')
pyplot.xlabel('time (s)')
pyplot.ylabel('position (m)')
```

**Compare with plot**

```
pyplot.subplot(212)
pyplot.plot(xs_ode,vs_ode,color='red')
pyplot.plot(xs_eul,vs_eul,color='blue')
pyplot.xlabel('position (m)')
pyplot.ylabel('velocity (m/s)')
```

```
pyplot.show()
```

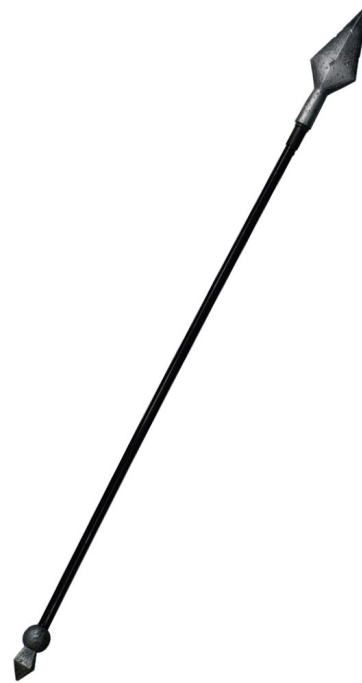


# Ballistics assessment

# Ballistics throughout history

## Pre-historic Physics

- “I’ve got a spear”
- “How fast do I have to throw it, and at what angle, for it to hit my prey?”



# Ballistics throughout history

## Modern Physics

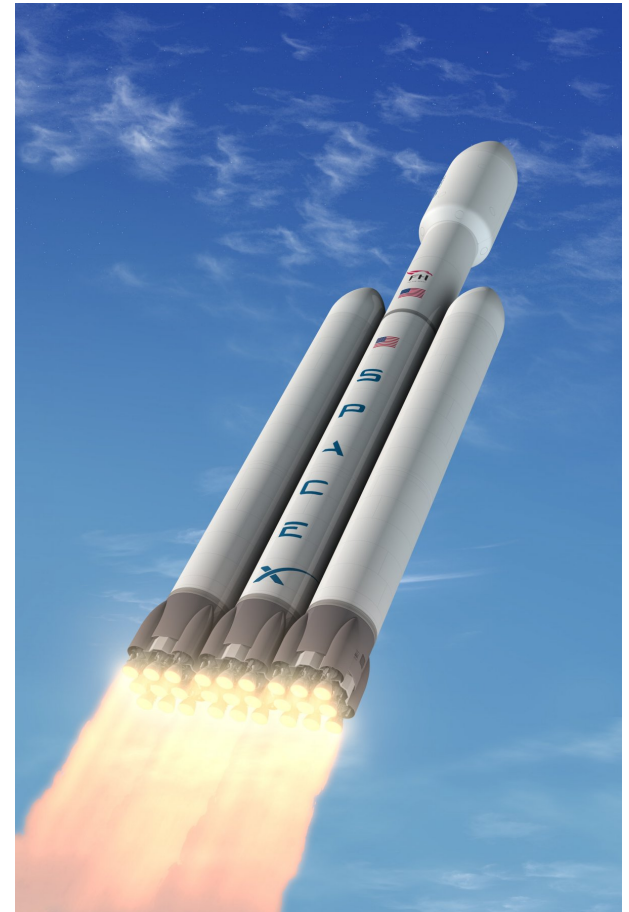
- “I’ve got Vickers and Armstrong 38.1cm gun. The wind speed is such, the air pressure is such and the temperature is such”
- “What azimuth elevation should I fire at to hit that ship 22 miles away?”





# Ballistics throughout history

- Exciting Physics!
  - I've got a rocket ship
  - How do I get to the Moon?



# How to answer the question...

- Write a differential equation for the forces on the projectile
- Solve it – numerically for realistic problems

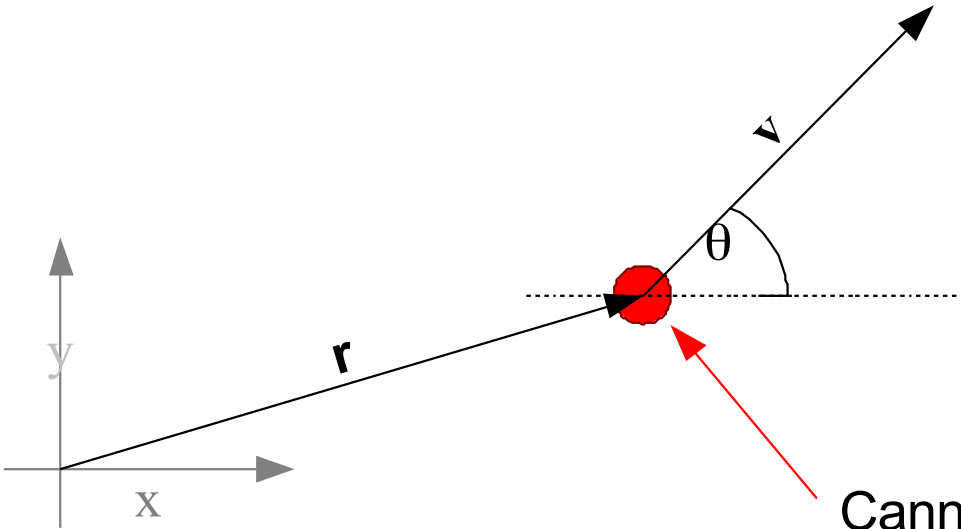
# Make some assumptions

- Height of projectile small compared to radius of earth ( $mg$  not  $GMm/r^2$ )
- Air pressure invariant of projectile height

# Assumptions

- It is standard to make simplifying assumptions when making a model
- Remember *GIGO* – *Garbage In, Garbage Out*
- Check your assumptions
  - especially that they remain valid as your model grows in complexity
- Document your assumptions

# Formulate the problem



Cannonball:

$m$  mass

$\mathbf{r}$  Position =  $(x, y)$

$\mathbf{v}$  Velocity =  $(v_x, v_y)$

# Forces

$$F_{grav} = \begin{pmatrix} F_{gx} \\ F_{gy} \end{pmatrix} = -mg \begin{pmatrix} 0 \\ 1 \end{pmatrix}$$

$$F_{drag} = -\kappa \rho A v \begin{pmatrix} v_x \\ v_y \end{pmatrix}$$

$$F = F_{grav} + F_{drag}$$

$$v^2 = (v_x^2 + v_y^2)$$

- $\kappa$  drag coefficient
- $\rho$  air density
- $A$  cross section area
- $v$   $(v_x^2 + v_y^2)^{.5}$

# Differential Equations

**2<sup>nd</sup> order**

$$\frac{d^2 x}{dt^2} = F_x / m$$

$$\frac{d^2 y}{dt^2} = F_y / m$$

**Turn into first order DEQs  
by introducing velocity**

# Differential Equations

**2<sup>nd</sup> order**

$$\frac{d^2 x}{dt^2} = F_x / m$$
$$\frac{d^2 y}{dt^2} = F_y / m$$



**1<sup>st</sup> order**

$$\frac{dv_x}{dt} = F_x / m$$
$$\frac{dv_y}{dt} = F_y / m$$
$$\frac{dx}{dt} = v_x$$
$$\frac{dy}{dt} = v_y$$



# Substitute in the forces

$$\frac{dv_x}{dt} = (-\kappa\rho Avv_x) / m$$

$$\frac{dv_y}{dt} = (-\kappa\rho Avv_y - mg) / m$$

$$\frac{dx}{dt} = v_x$$

$$\frac{dy}{dt} = v_y$$