

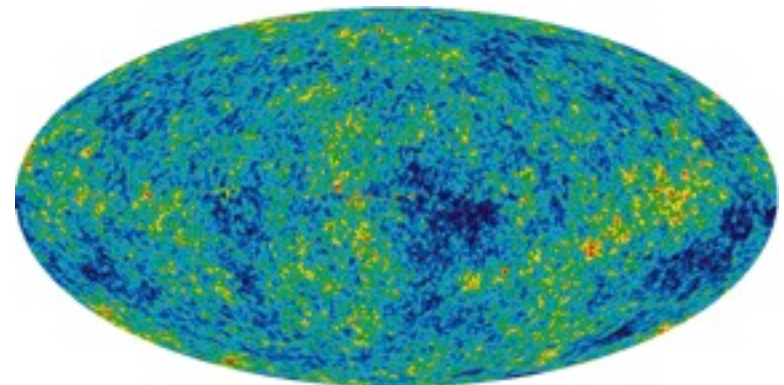
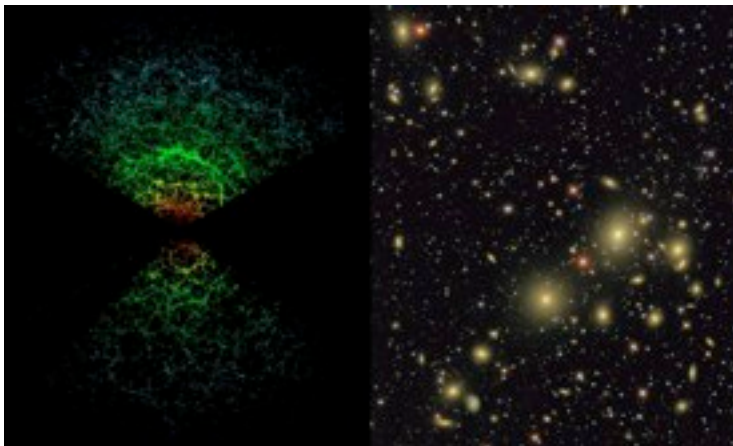
The Quest for Primordial Non-Gaussianity

Overview and some recent developments;
skewed toward observations

Dragan Huterer
University of Michigan

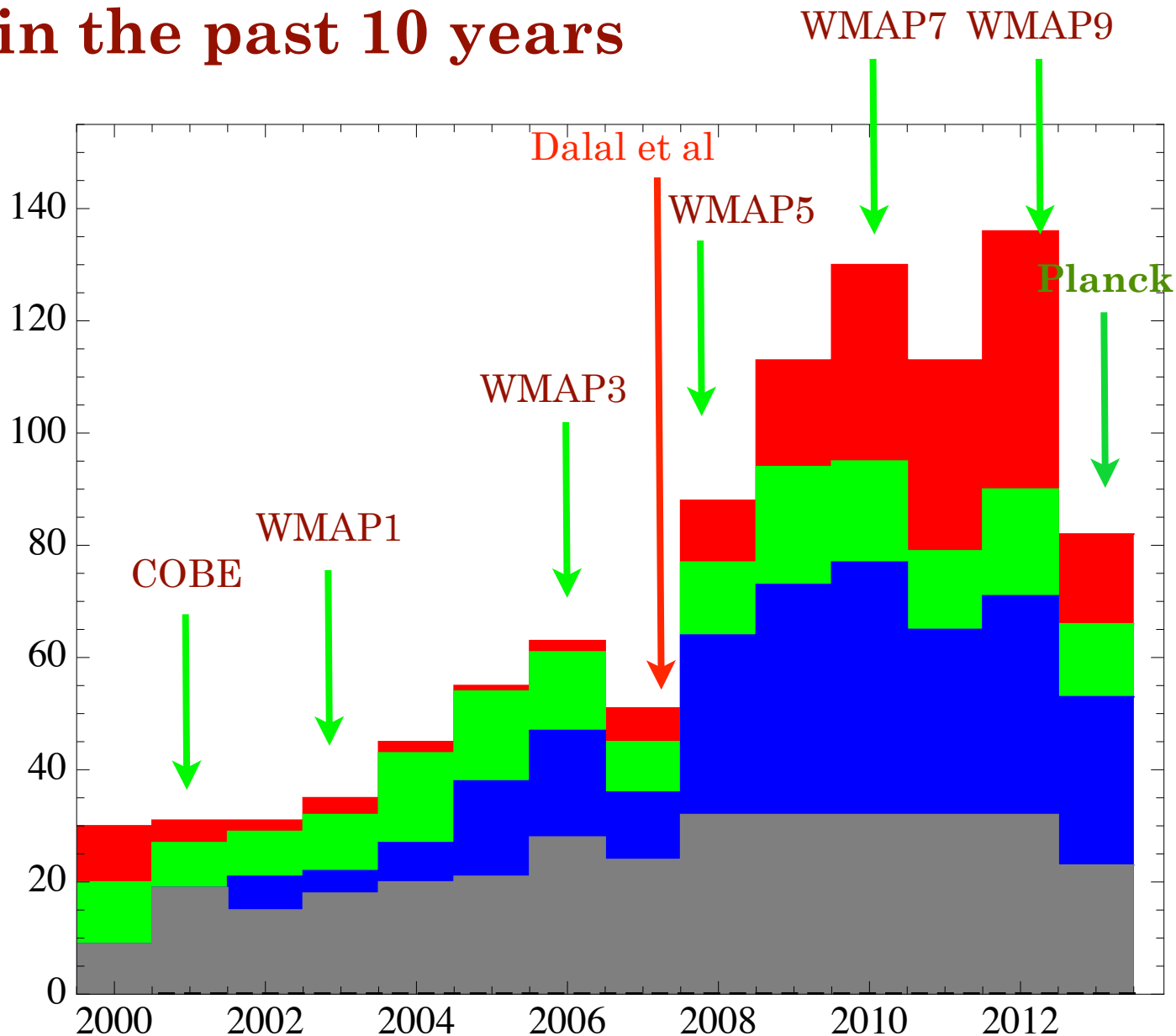
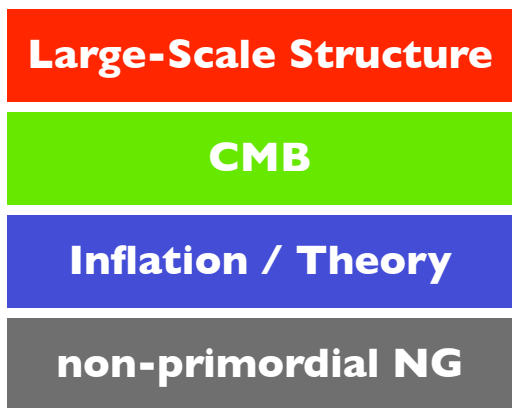
Why study non-Gaussianity (NG)?

1. NG presents a window to the very early universe. For example, NG can distinguish between physically distinct models of inflation.
2. Conveniently, NG can be constrained/measured using CMB anisotropy maps and LSS. In particular, there is a rich set of observable quantities that are sensitive to primordial NG.



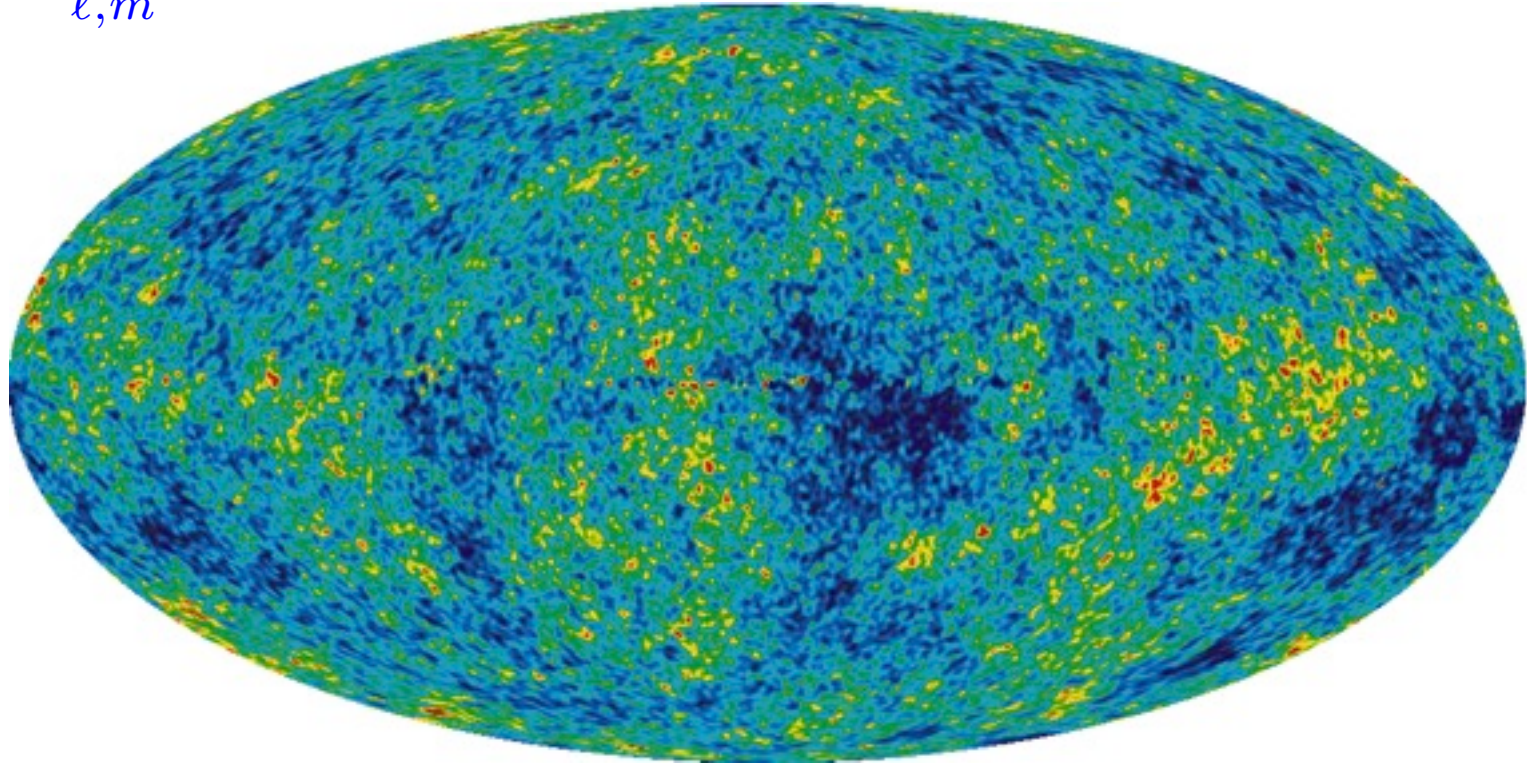
Non-Gaussianity papers in the past 10 years

of articles with
"Non-Gaussian"
in the title
on the ADS data base



Initial conditions in the universe

$$\frac{\delta T}{T}(\theta, \phi) = \sum_{\ell, m} a_{\ell m} Y_{\ell m}(\theta, \phi) \quad \ell \simeq \frac{180^\circ}{\theta}$$



Generic inflationary predictions: **Statistical Isotropy:**

$$\langle a_{\ell m} a_{\ell' m'} \rangle \equiv C_{\ell \ell' m m'} = C_\ell \delta_{\ell \ell'} \delta_{m m'}$$

- Nearly scale-invariant spectrum of density perturbations
- Background of gravity waves

Gaussianity:

- (Very nearly) gaussian initial conditions: $\langle a_{\ell m} a_{\ell' m'} a_{\ell'' m''} \rangle = 0$

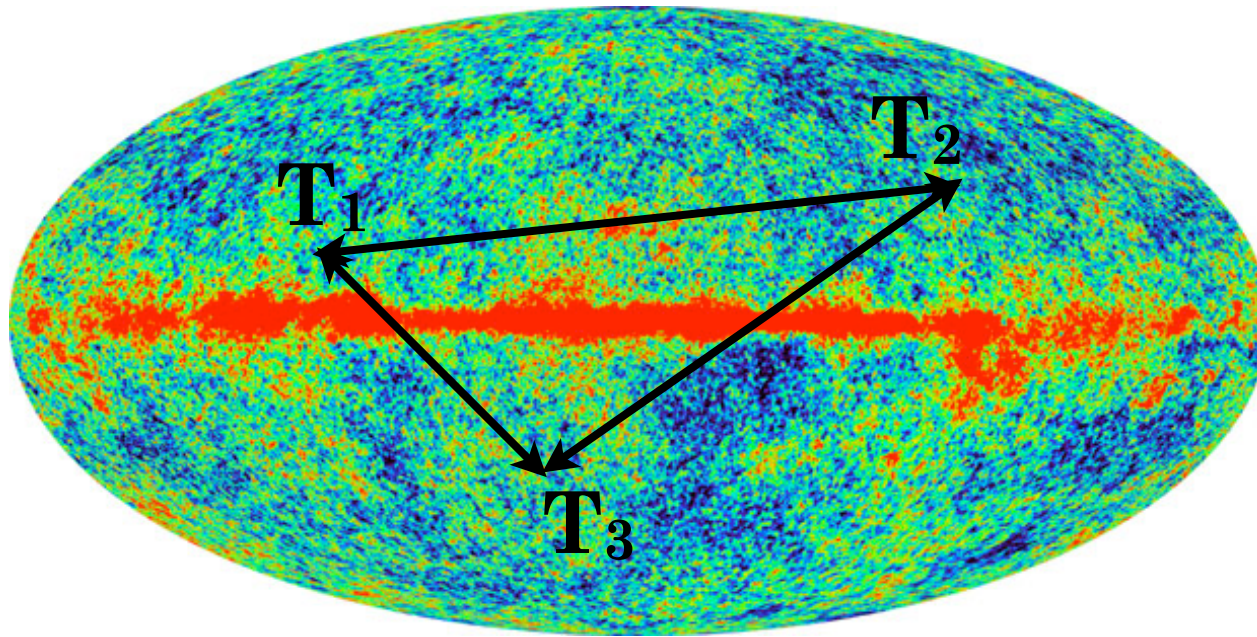
Standard Inflation, with...

1. a single scalar field
2. the canonical kinetic term
3. always slow rolls
4. in Bunch-Davies vacuum
5. in Einstein gravity

produces **unobservable** NG

Therefore, measurement of nonzero NG would point to a **violation** of one of the assumptions above

NG from 3-point correlation function



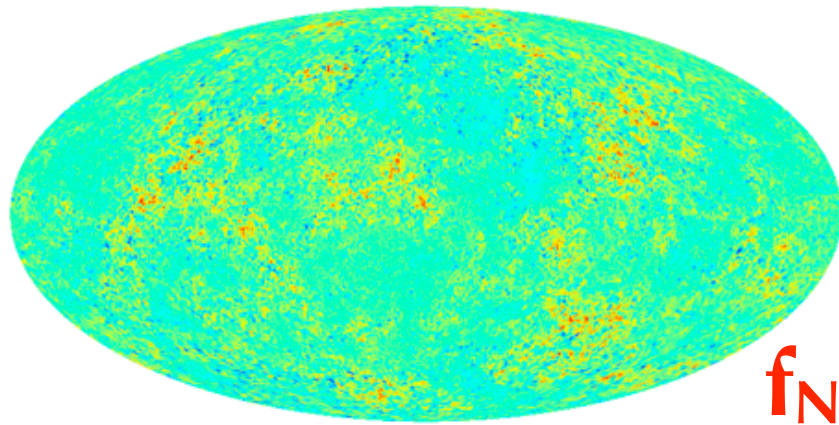
Commonly used “local” model of NG

$$\Phi = \Phi_G + f_{\text{NL}} (\Phi_G^2 - \langle \Phi_G^2 \rangle)$$

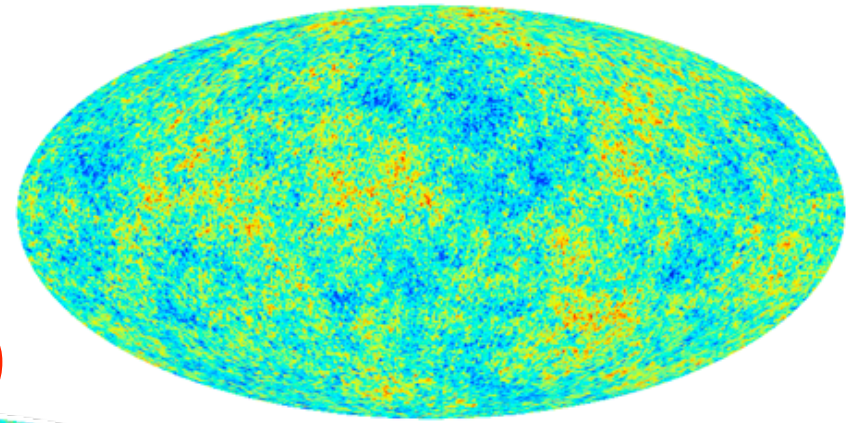
Salopek & Bond 1990; Verde et al 2000; Komatsu & Spergel 2001; Maldacena 2003

Then the 3-point function is related to f_{NL} via (in k-space)

$$B(k_1, k_2, k_3) \sim f_{\text{NL}} [P(k_1)P(k_2) + \text{perm.}]$$

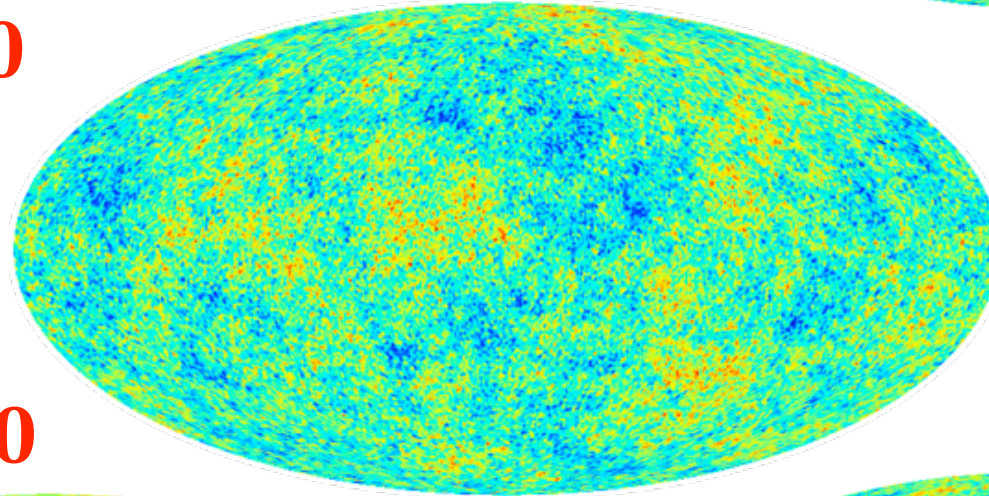


$f_{\text{NL}} = -5000$



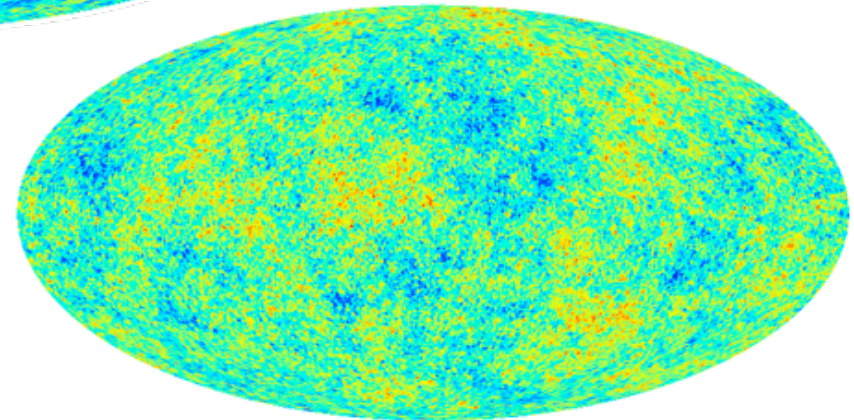
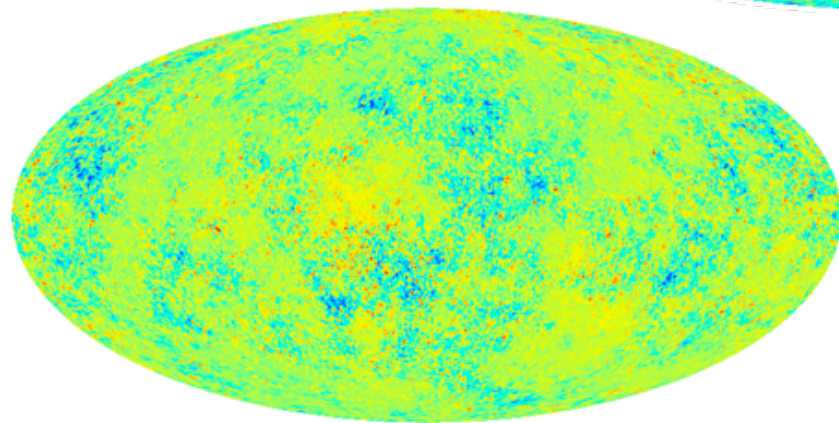
$f_{\text{NL}} = -500$

$f_{\text{NL}} = 0$

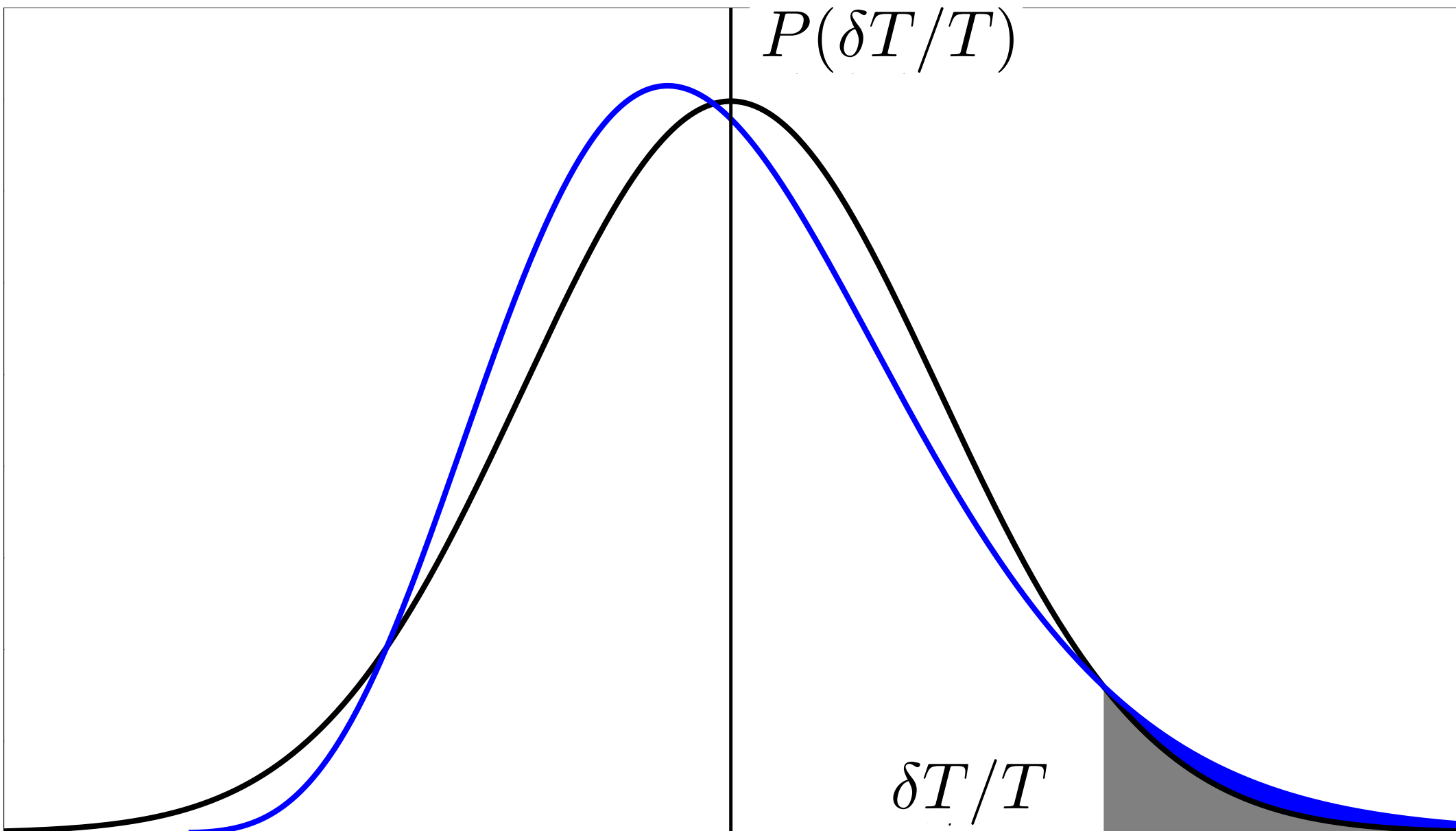


$f_{\text{NL}} = +5000$

$f_{\text{NL}} = +500$



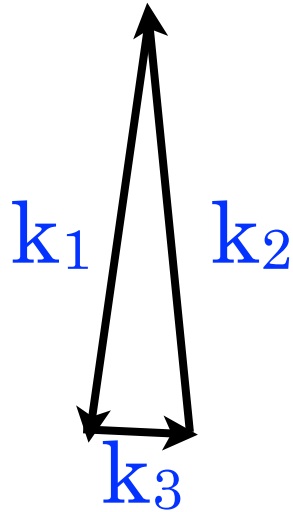
Current upper bound on NG is
~1000 times smaller than **this**:



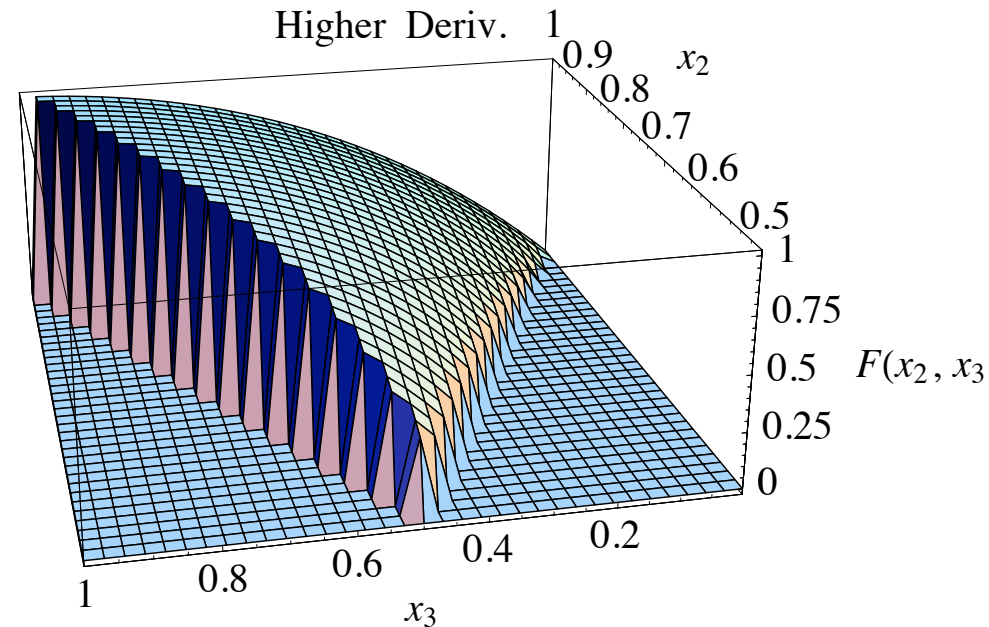
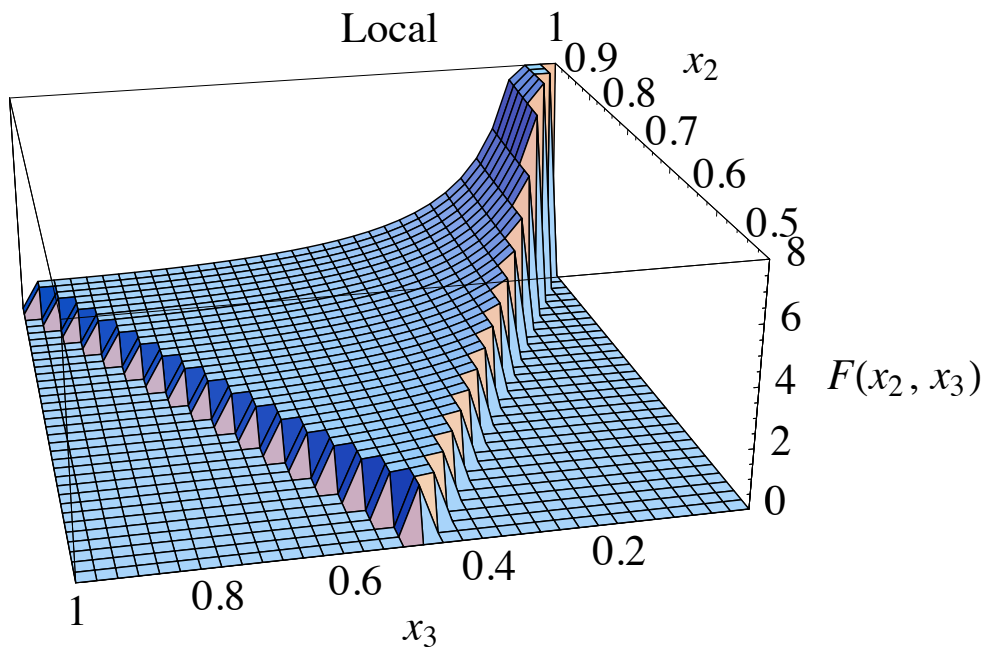
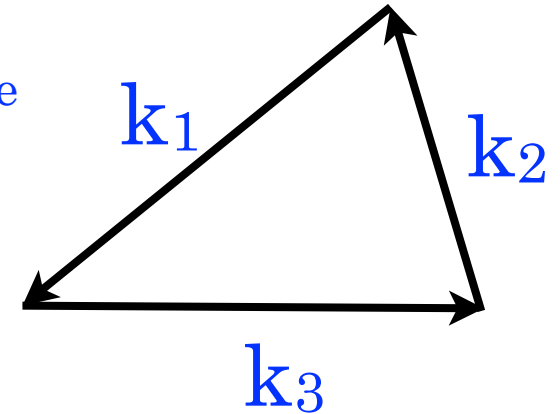
3-pt correlation function of CMB anisotropy ⇒ direct window into inflation

e.g. Luo & Schramm 1993

“local”
(eg. multi-field)



“equilateral”
(eg. higher-derivative
action; interactions)



Babich, Creminelli & Zaldarriaga 2004

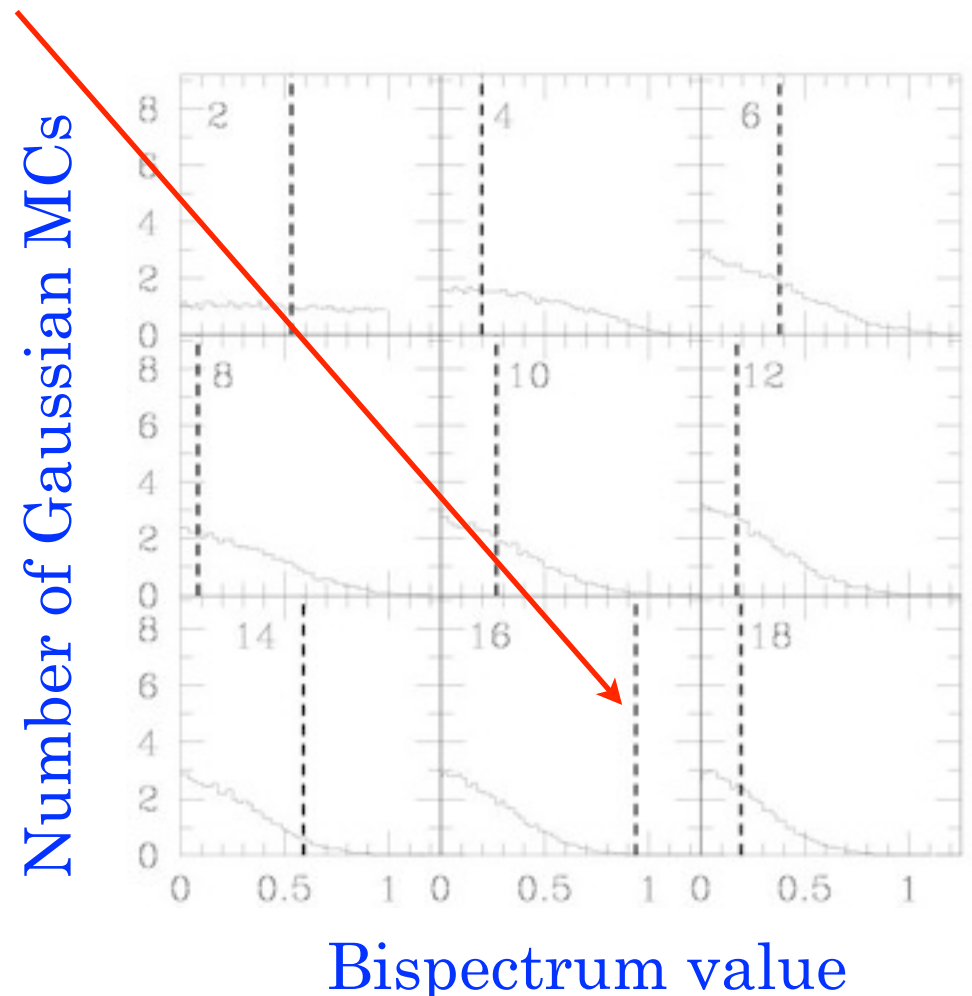
Brief history of NG measurements: 1990's

Early 1990s; COBE: Gaussian CMB sky (Kogut et al 1996)

1998; COBE: claim of NG at $l=16$ equilateral bispectrum (Ferreira, Magueijo & Gorski 1998)

but explained by a known systematic effect! (Banday, Zaroubi & Gorski 1999)

(and anyway isn't unexpected given all bispectrum configurations you can measure; Komatsu 2002)



Brief history of NG measurements: 2000's

Pre-WMAP CMB: all is gaussian (e.g. MAXIMA; Wu et al 2001)

WMAP pre-2008: all is gaussian

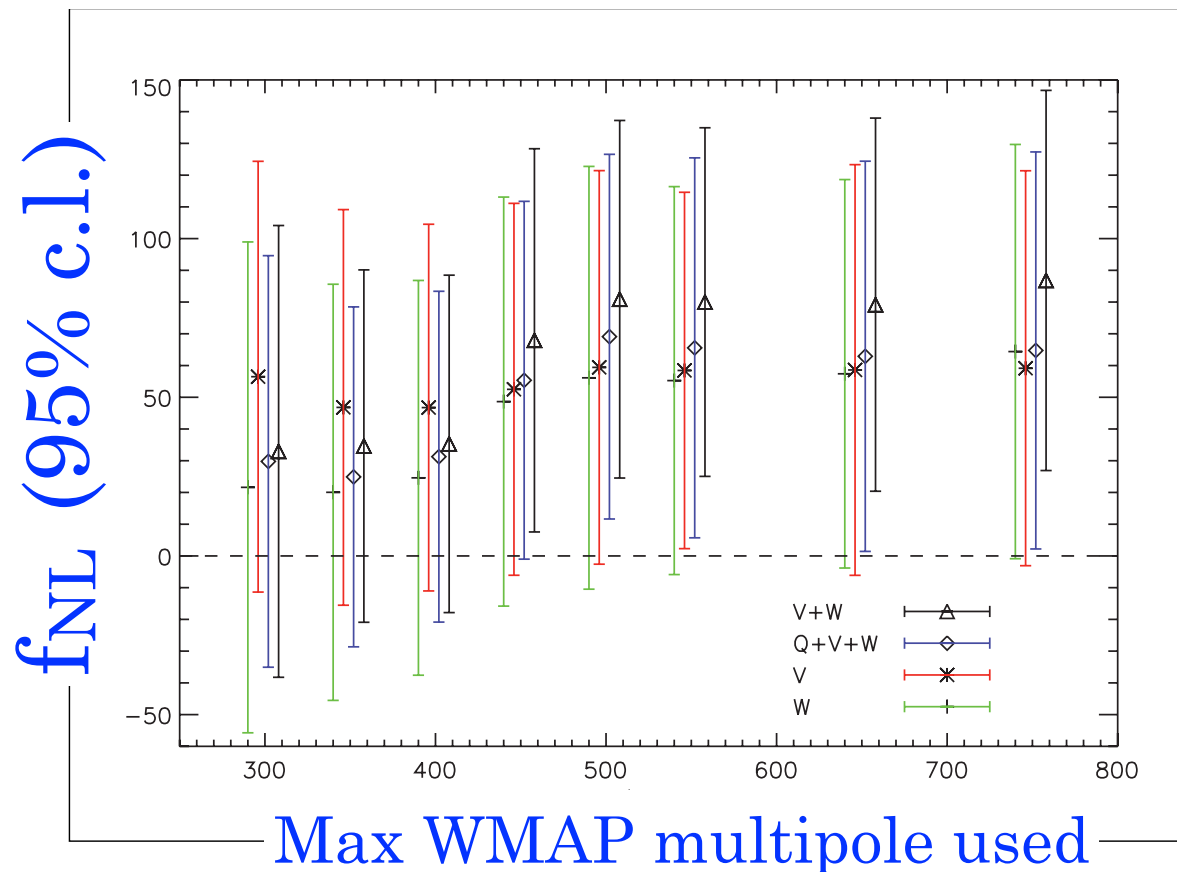
(Komatsu et al. 2003; Creminelli, Senatore, Zaldarriaga & Tegmark 2007)

$$-36 < f_{\text{NL}} < 100 \quad (95\% \text{ CL})$$

Dec 2007, claim of NG in WMAP

(Yadav & Wandelt arXiv:0712.1148)

$$27 < f_{\text{NL}} < 147 \quad (95\% \text{ CL})$$



Constraints from WMAP

Band	Foreground ^b	f_{NL}^{local}	f_{NL}^{equil}	f_{NL}^{orthog}	b_{src}
V+W	Raw	59 ± 21	33 ± 140	-199 ± 104	N/A
V+W	Clean	42 ± 21	29 ± 140	-198 ± 104	N/A
V+W	Marg. ^c	32 ± 21	26 ± 140	-202 ± 104	-0.08 ± 0.12
V	Marg.	43 ± 24	64 ± 150	-98 ± 115	0.32 ± 0.23
W	Marg.	39 ± 24	36 ± 154	-257 ± 117	-0.13 ± 0.19

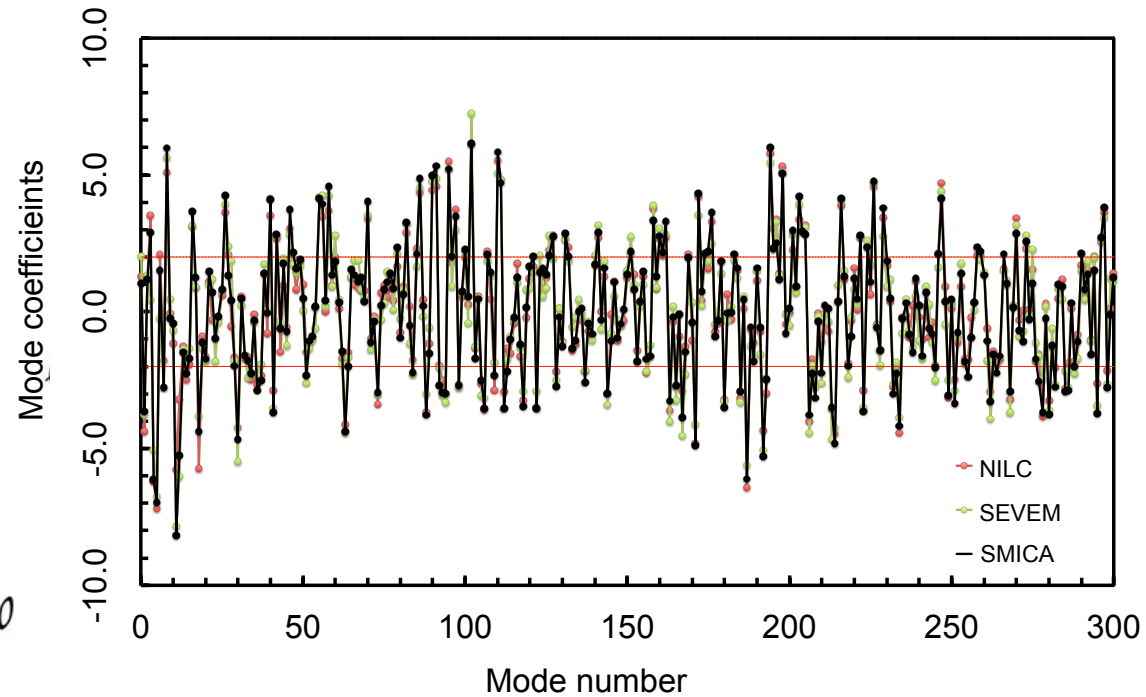
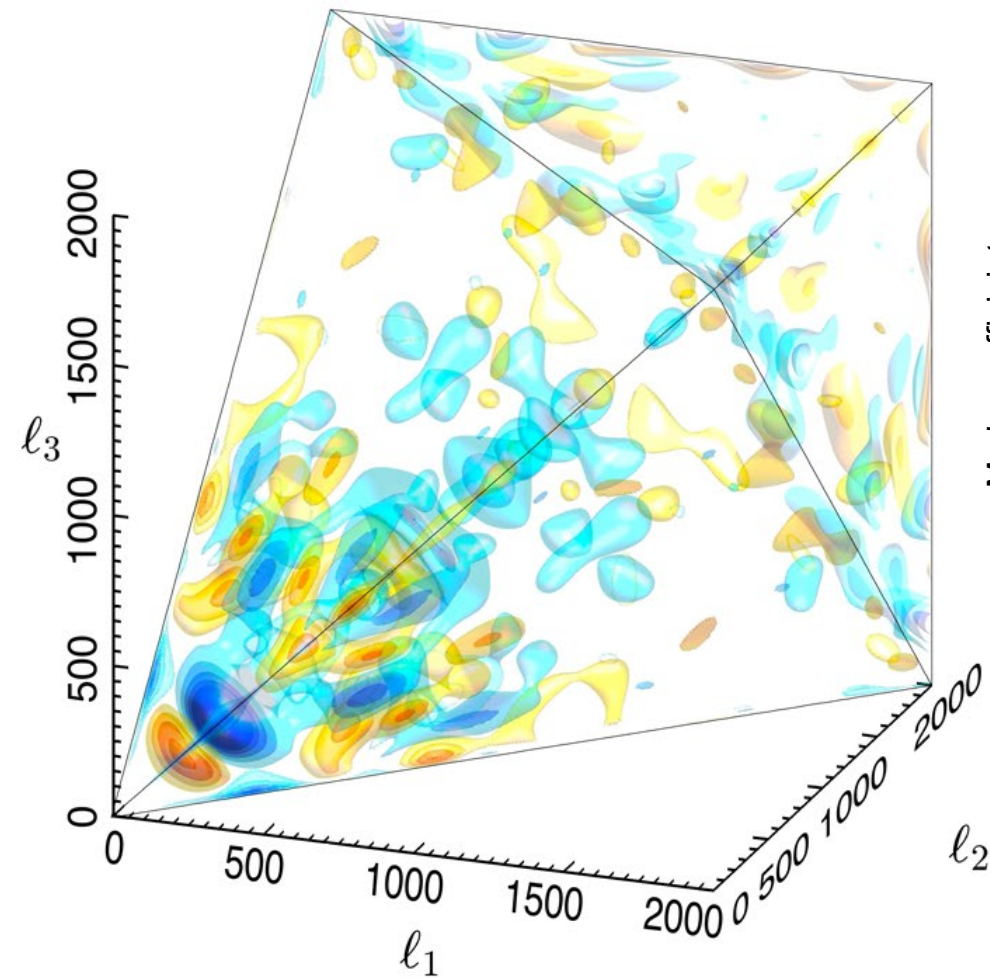
Komatsu et al. 2010

Constraints from Planck

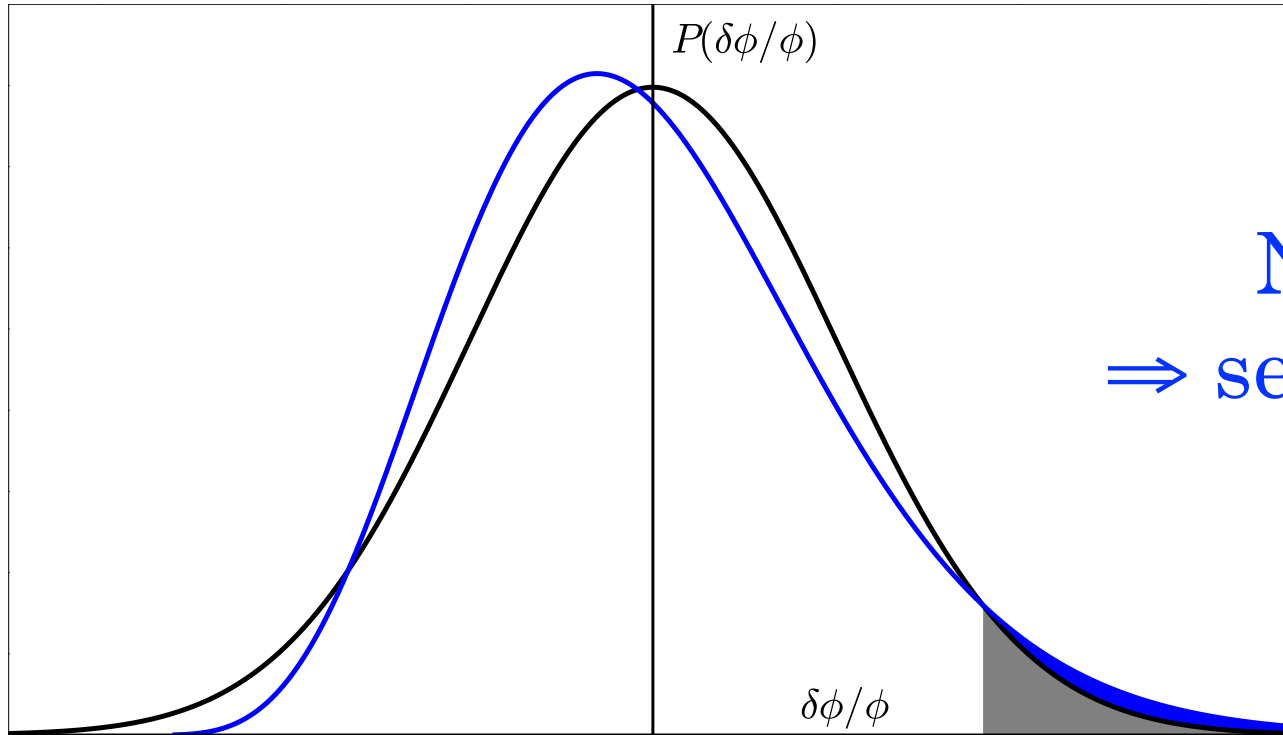
	ISW-lensing subtracted		
	KSW	Binned	Modal
SMICA			
Local	2.7 ± 5.8	2.2 ± 5.9	1.6 ± 6.0
Equilateral	-42 ± 75	-25 ± 73	-20 ± 77
Orthogonal	-25 ± 39	-17 ± 41	-14 ± 42
NILC			
Local	4.5 ± 5.8	3.6 ± 5.8	2.7 ± 6.0
Equilateral	-48 ± 76	-38 ± 73	-20 ± 78
Orthogonal	-53 ± 40	-41 ± 41	-37 ± 43
SEVEM			
Local	3.4 ± 5.9	3.2 ± 6.2	2.6 ± 6.0
Equilateral	-36 ± 76	-25 ± 73	-13 ± 78
Orthogonal	-14 ± 40	-9 ± 42	-2 ± 42
C-R			
Local	6.4 ± 6.0	5.5 ± 5.9	5.1 ± 5.9
Equilateral	-62 ± 79	-55 ± 74	-32 ± 78
Orthogonal	-57 ± 42	-41 ± 42	-42 ± 42

Constraints from Planck: modal expansion

$$B(k_1, k_2, k_3) = \sum_{p,r,s} \alpha_{prs} q_p(k_1) q_r(k_2) q_s(k_3)$$



Galaxy cluster counts' sensitivity to NG



NG initial PDF
⇒ sensitivity to counts
“on the tail”

(amount of NG shown is $>100\times$ bigger than allowed by data!)

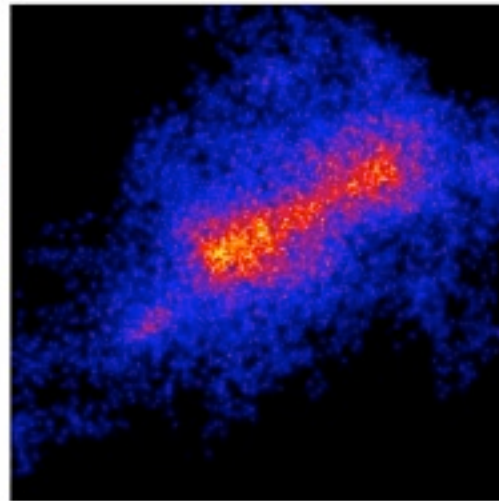
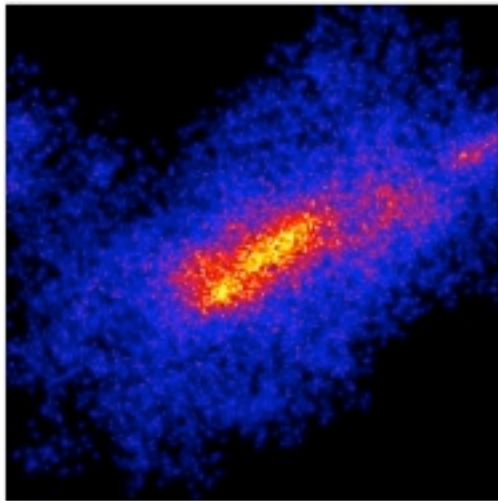
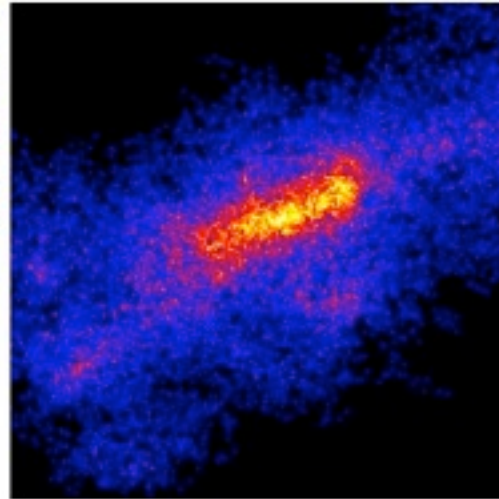
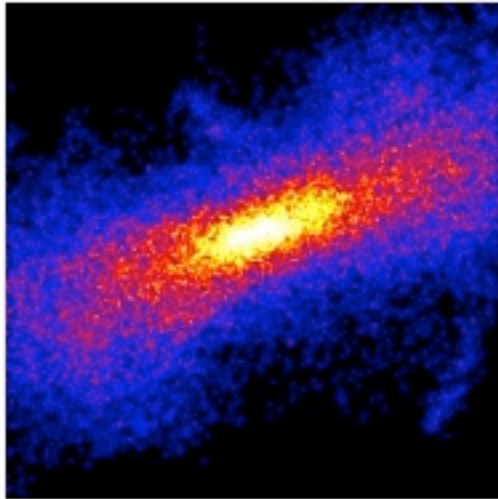
Lots of effort in the community to calibrate
the **non-Gaussian mass function** -
 $dn/d\ln M(M, z)$ - of DM halos

(analytic extensions of Press-Schechter + simulations)

DM halo gets more massive with $f_{\text{NL}} > 0$ (and v.v.)

$f_{\text{NL}} = +5000$
 $M = 1.2 \cdot 10^{16} M_{\odot}$

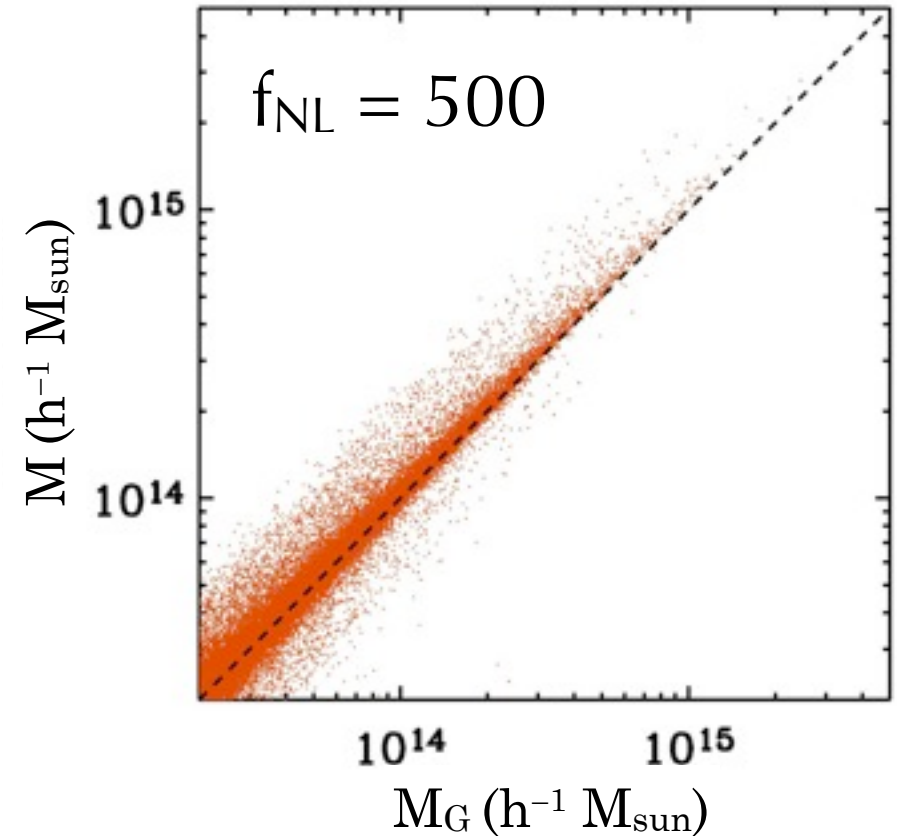
$f_{\text{NL}} = +500$
 $M = 5.9 \cdot 10^{15} M_{\odot}$



$f_{\text{NL}} = 0$
 $M = 5.1 \cdot 10^{15} M_{\odot}$

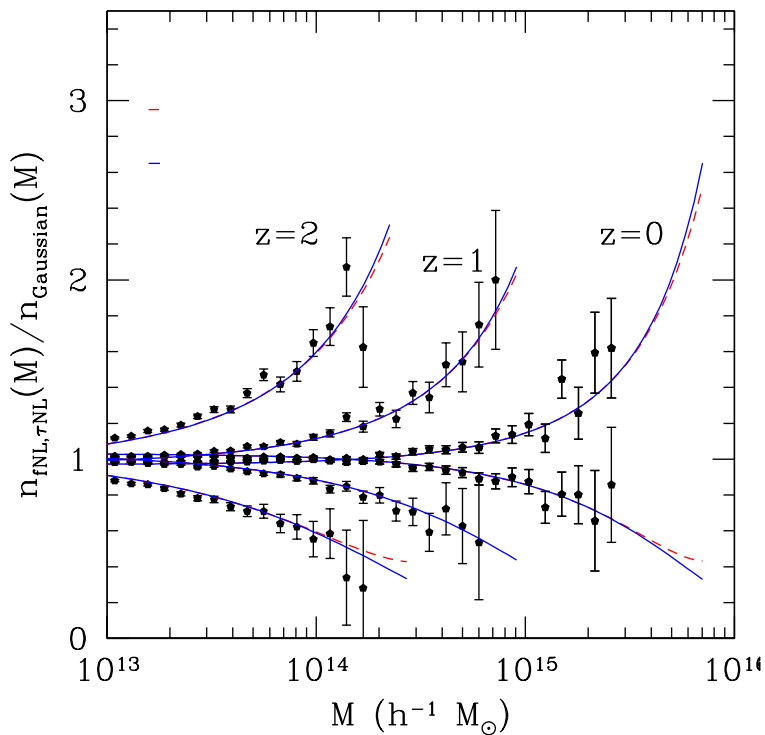
$f_{\text{NL}} = -500$
 $M = 4.3 \cdot 10^{15} M_{\odot}$

Mapping between
 M_{G} and $M \equiv M_{\text{NG}}$:



\Rightarrow NG mass function:

$$\frac{dN}{dM} = \int \frac{dP(M|M_{\text{G}})}{dM} \frac{dN}{dM_{\text{G}}} dM_{\text{G}}$$



NG/Gaussian **mass function** ratios:
for fixed M , more sensitivity
at higher redshift

Smith & LoVerde 2011; Pillepich, Porciani and Hahn 2009;
many others going back to 1990s

Unfortunately, cluster counts are **weakly**
sensitive to NG

e.g. $\sigma(f_{NL})=450$ measured from SPT (Williamson et al 2010)

Nevertheless:

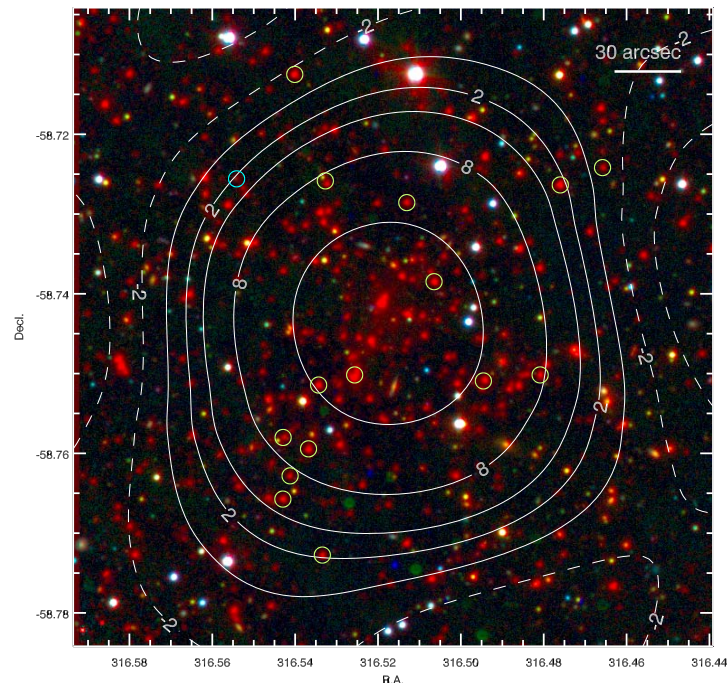
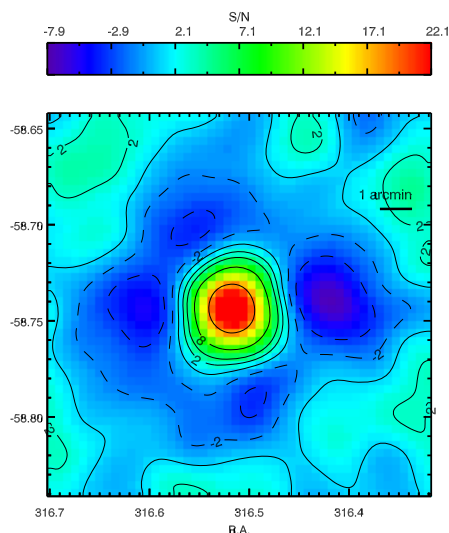
- cluster abundance is sensitive to ALL non-Gaussianity
- (large) amount of (local model) NG can boost the number of ‘pink elephant’ clusters

High-z, high-M - "pink elephant" - clusters of galaxies

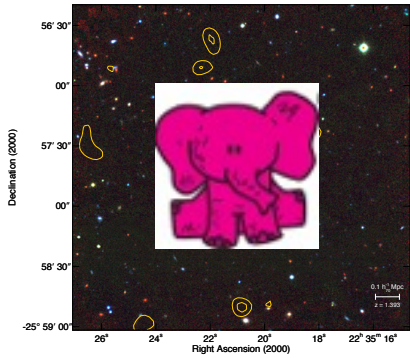
- SPT-CL J0546-5045: $z=1.067$, $M \approx (8.0 \pm 1.0) \cdot 10^{14} M_{\text{sun}}$
- XMMU J2235.3-2557: $z=1.39$, $M \approx (8.5 \pm 1.7) \cdot 10^{14} M_{\text{sun}}$
- SPT-CL J2106-8544: $z=1.132$, $M \approx (1.3 \pm 0.2) \cdot 10^{15} M_{\text{sun}}$

Some authors have claimed the existence of these clusters is in **conflict** with LCDM, but can be explained with (huge; $f_{\text{NL}} \sim 500$)
non-Gaussianity

Hoyle, Jimenez & Verde (2011);
Cayon, Gordon & Silk (2011);
Holz & Perlmutter 2011



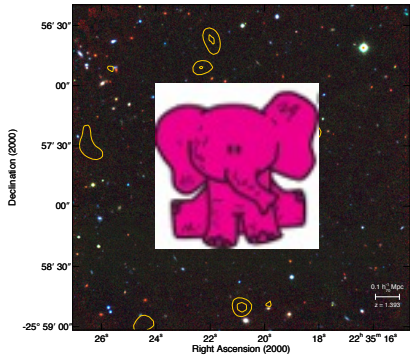
Are the pink elephants in conflict with LCDM?!



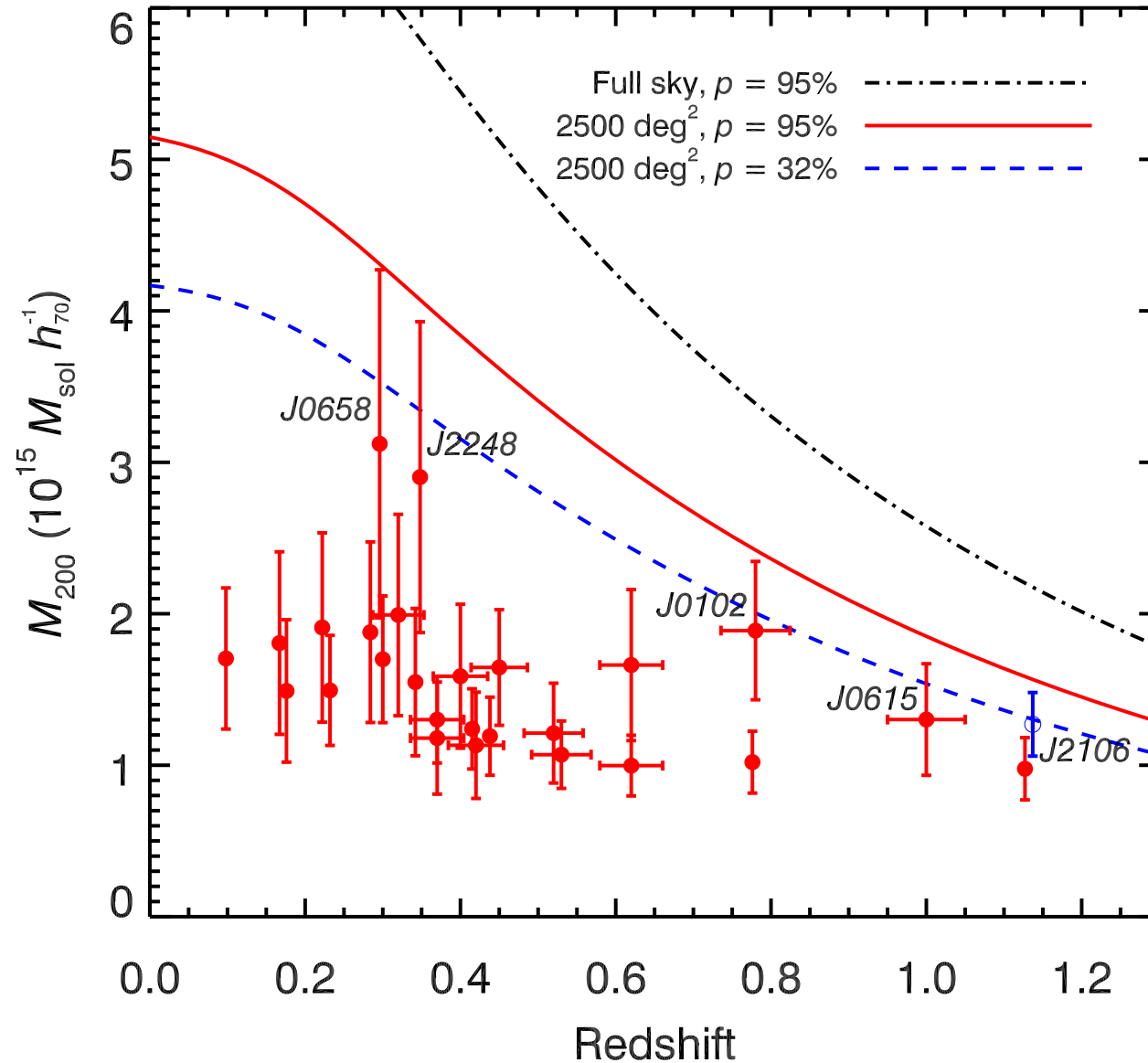
4 things to account for:

1. **Sample variance** - the Poisson noise in counting rare objects in a finite volume
2. **Parameter variance** - uncertainty due to fact that current data allow cosmological parameters to take a range of values
3. **Eddington bias** - mass measurement error will preferentially 'scatter' the cluster into higher mass
4. **Survey sky coverage** - needs to be fairly assessed

N.B. If a cluster rules out LCDM, it will rule out quintessence too!



No conflict - for now.



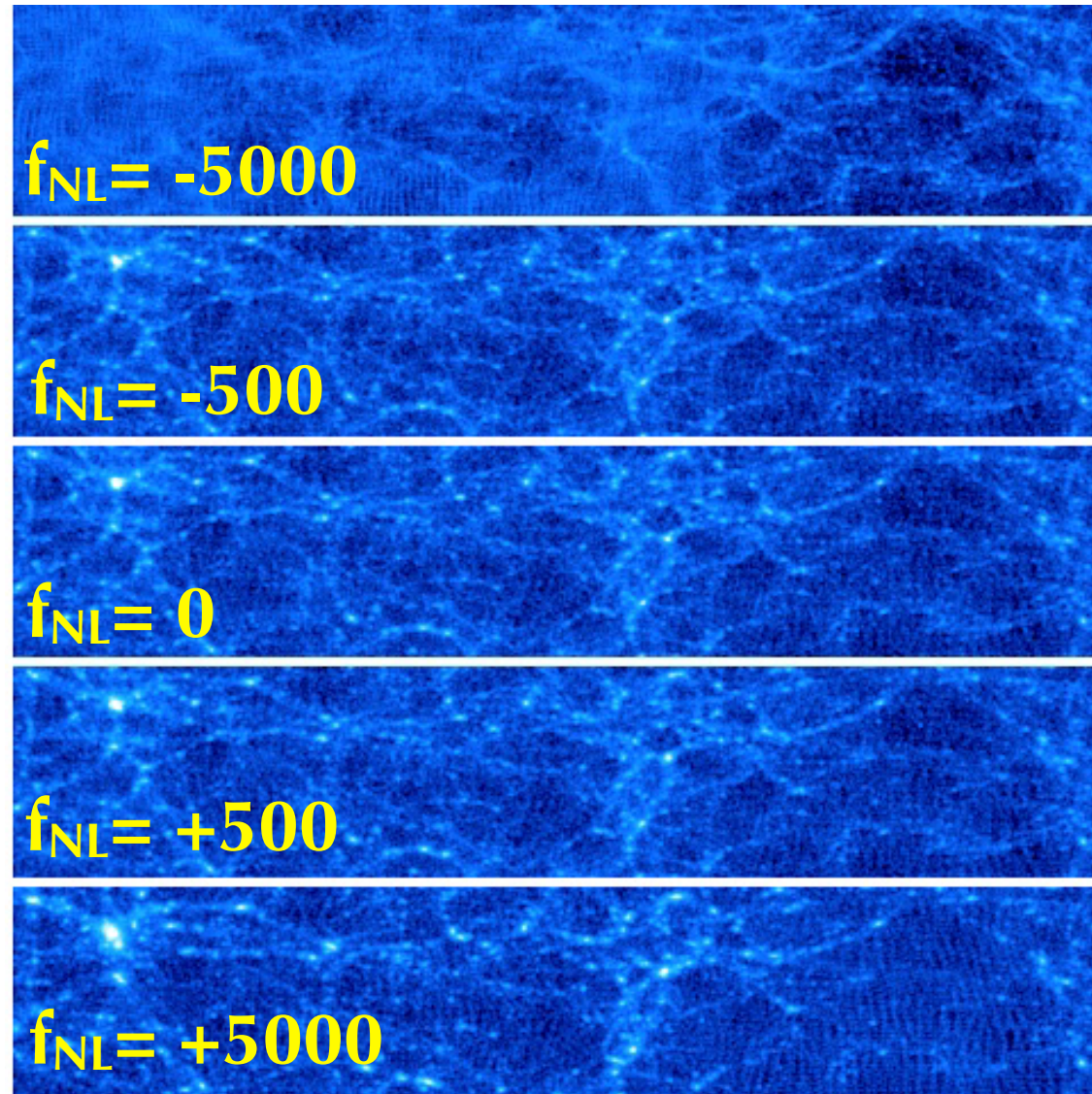
Foley et al 2011 arXiv:1101.1286 (SPT team);
 Mortonson, Hu & Huterer 2011

Next Frontier: Large-Scale Structure

	CMB	LSS
dimension	2D	3D
# modes	$\propto l_{\max}^2$	$\propto k_{\max}^3$
systematics & selection func.	relatively clean	relatively messy
temporal evol.	no	yes
can slice in	λ only	$\lambda, M, \text{bias} \dots$

Effects of primordial NG on the bias of virialized objects

Simulations with non-Gaussianity (f_{NL})



- Under-dense region evolution decrease with f_{NL}
- Over-dense region evolution increase with f_{NL}

80 Mpc/h

375 Mpc/h

- Same initial conditions, different f_{NL}
- Slice through a box in a simulation $N_{\text{part}}=512^3$, $L=800$ Mpc/h

Does galaxy/halo bias depend on NG?

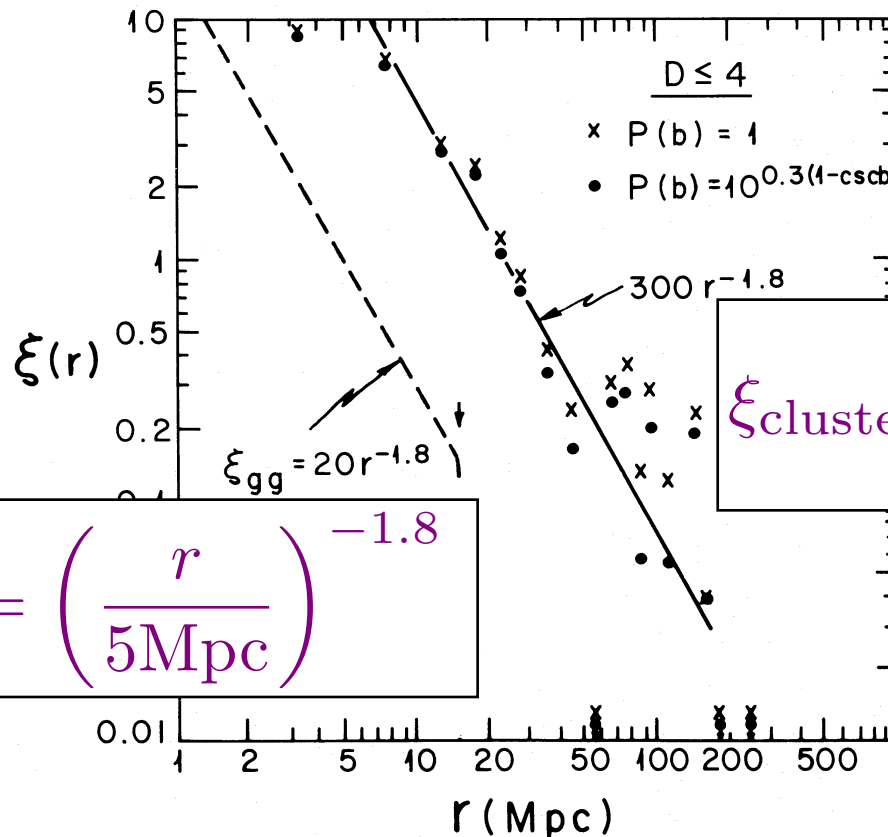
$$\text{bias} \equiv \frac{\text{clustering of galaxies}}{\text{clustering of dark matter}} = \frac{\left(\frac{\delta\rho}{\rho}\right)_{\text{halos}}}{\left(\frac{\delta\rho}{\rho}\right)_{\text{DM}}}$$

cosmologists measure

DM

usually nuisance parameter(s)

theory predicts

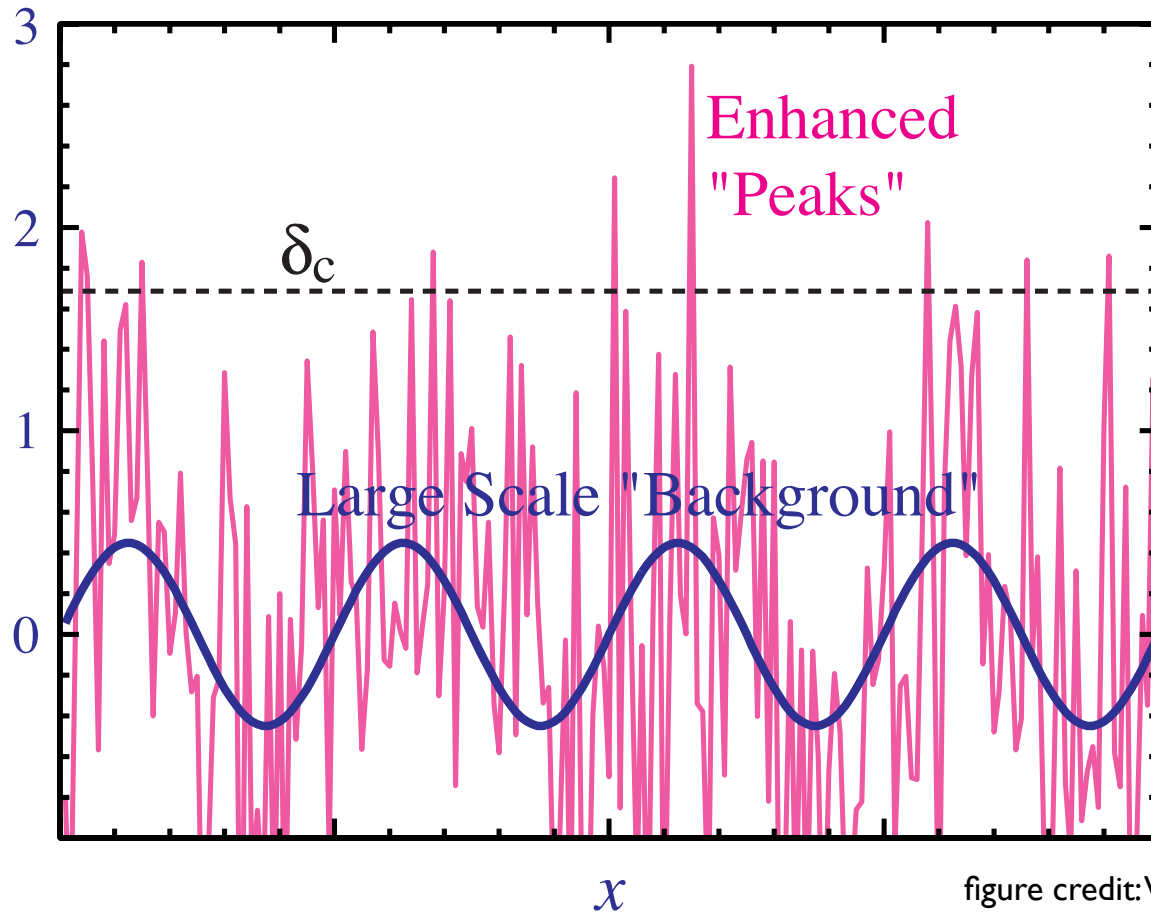


$$\xi_{\text{clusters}}(r) = \left(\frac{r}{25\text{Mpc}}\right)^{-1.8}$$

$$\xi_{\text{galaxies}}(r) = \left(\frac{r}{5\text{Mpc}}\right)^{-1.8}$$

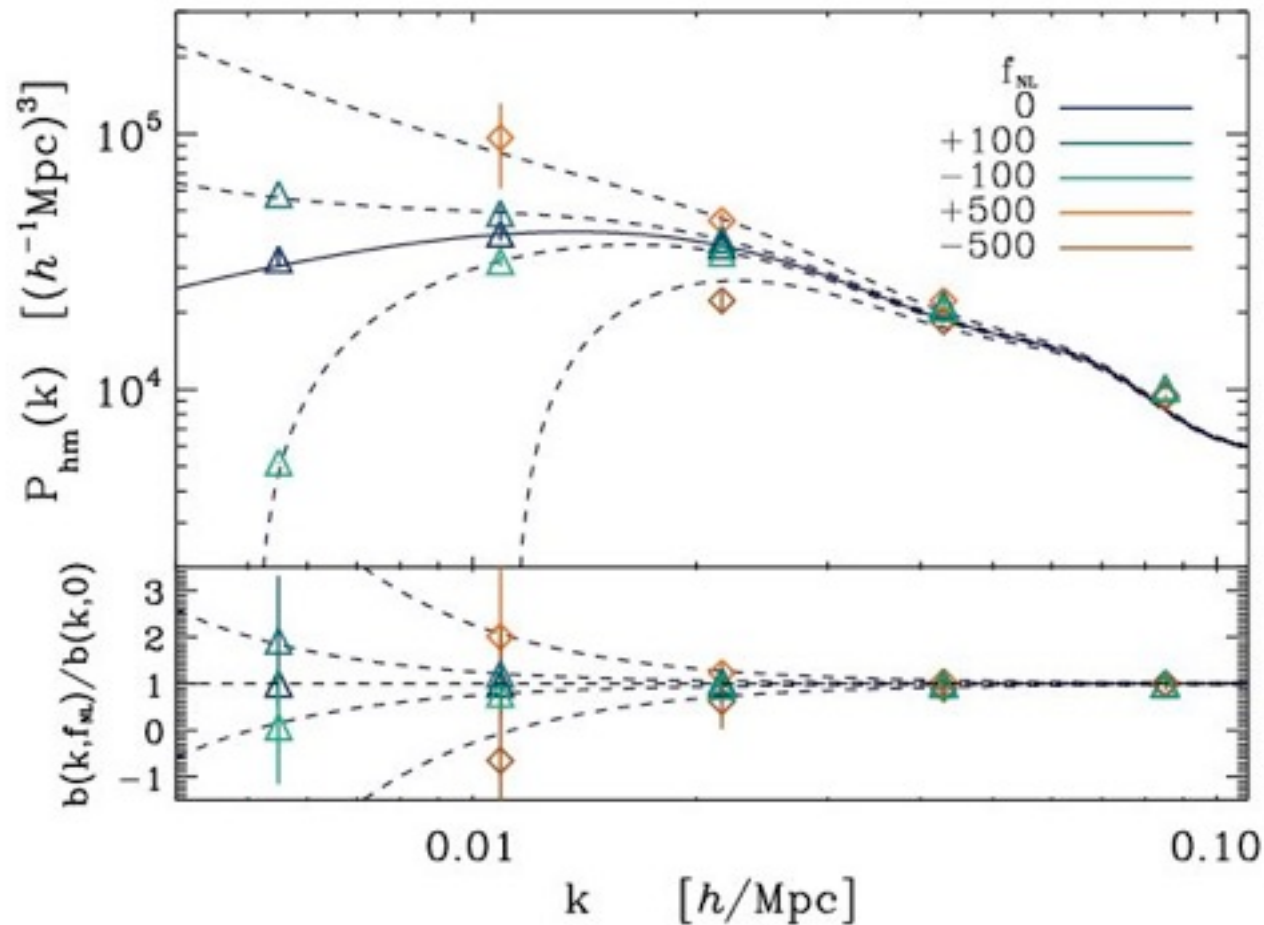
Bias of dark matter halos

$$P_h(k, z) = b^2(k, z) P_{\text{DM}}(k, z)$$



Simulations and theory both say: **large-scale bias is scale-independent**
(theorem if halo abundance is function of local density
and if the short and long modes are uncorrelated)

Scale dependence of NG halo bias



$$b(k) = b_G + f_{\text{NL}} \frac{\text{const}}{k^2}$$

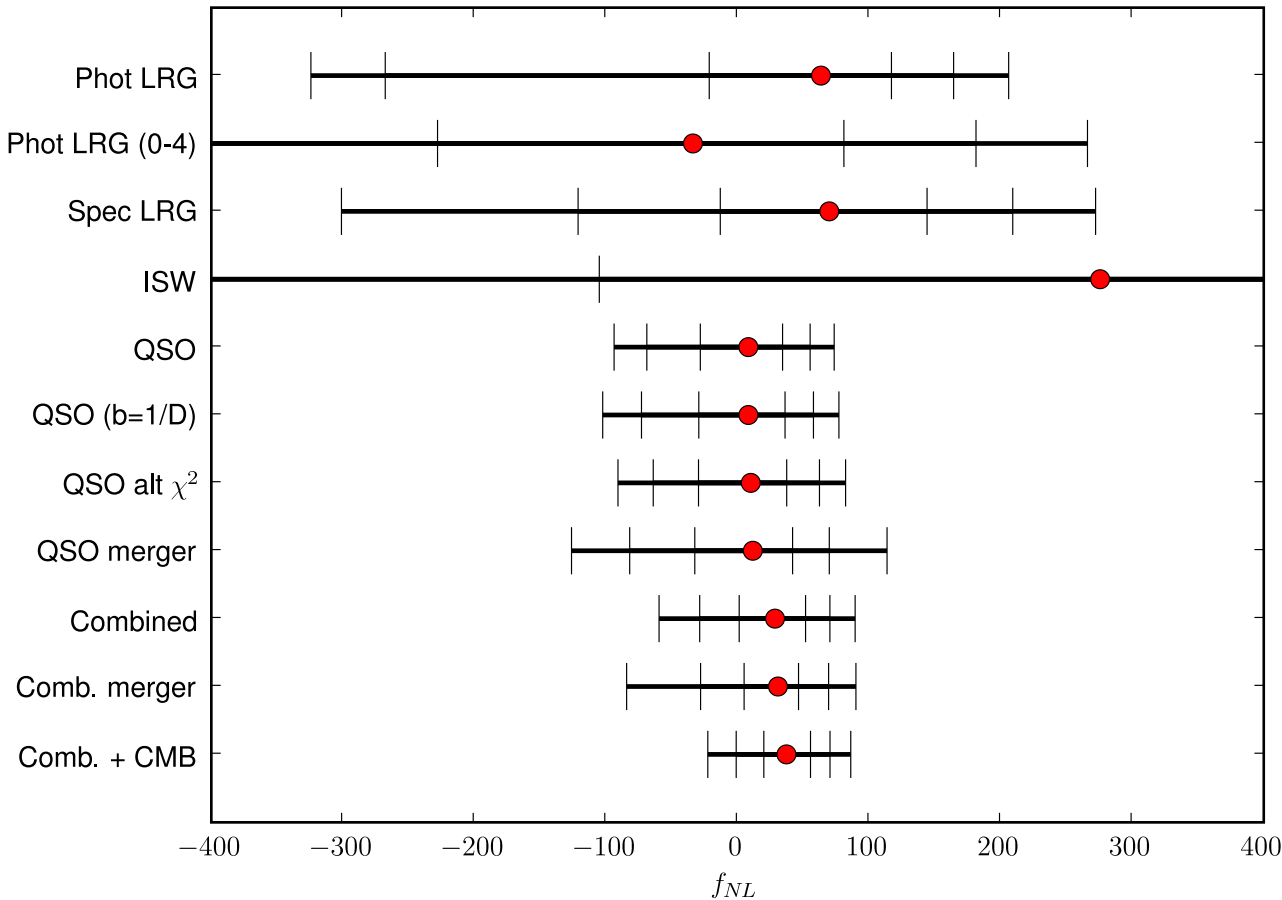
Verified using a variety of theoretical derivations and numerical simulations.

$$\Delta b(k) = f_{\text{NL}}(b_G - 1) \delta_c \frac{3 \Omega_M H_0^2}{T(k) D(a) k^2}$$

Implications:

- ▶ Unique $1/k^2$ scaling of bias; no free parameters
- ▶ Distinct from effect of all other cosmo parameters
- ▶ Straightforwardly measured (g-g, g-T,...)
- ▶ Derived theoretically several different ways
- ▶ Extensively tested with numerical simulations; good agreement found

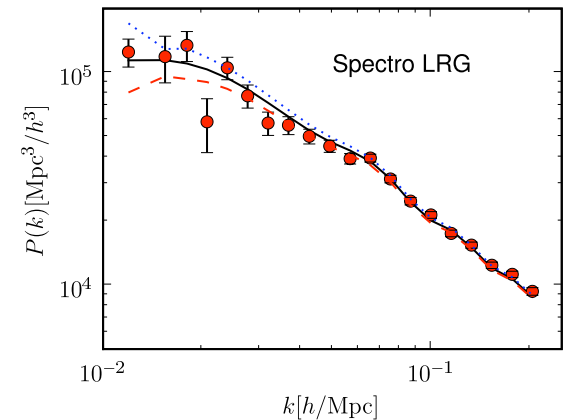
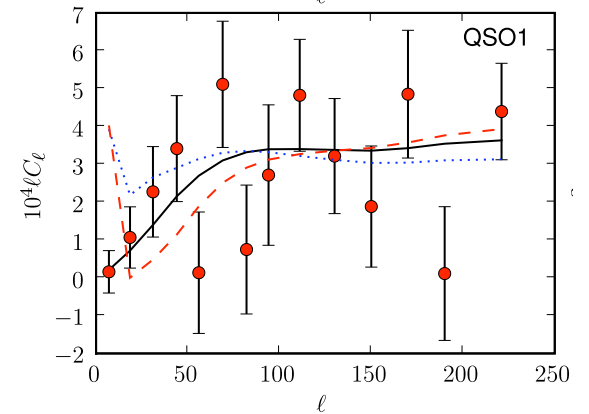
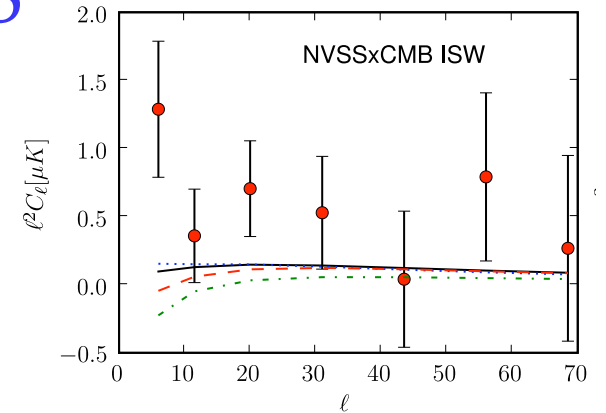
Constraints from **current** data: SDSS



$f_{NL} = 8 \pm 30$ (68%, QSO)

Slosar et al. 2008

$f_{NL} = 23 \pm 23$ (68%, all)



Future data forecasts for LSS: $\sigma(f_{NL}) \approx \mathcal{O}(\text{few})$

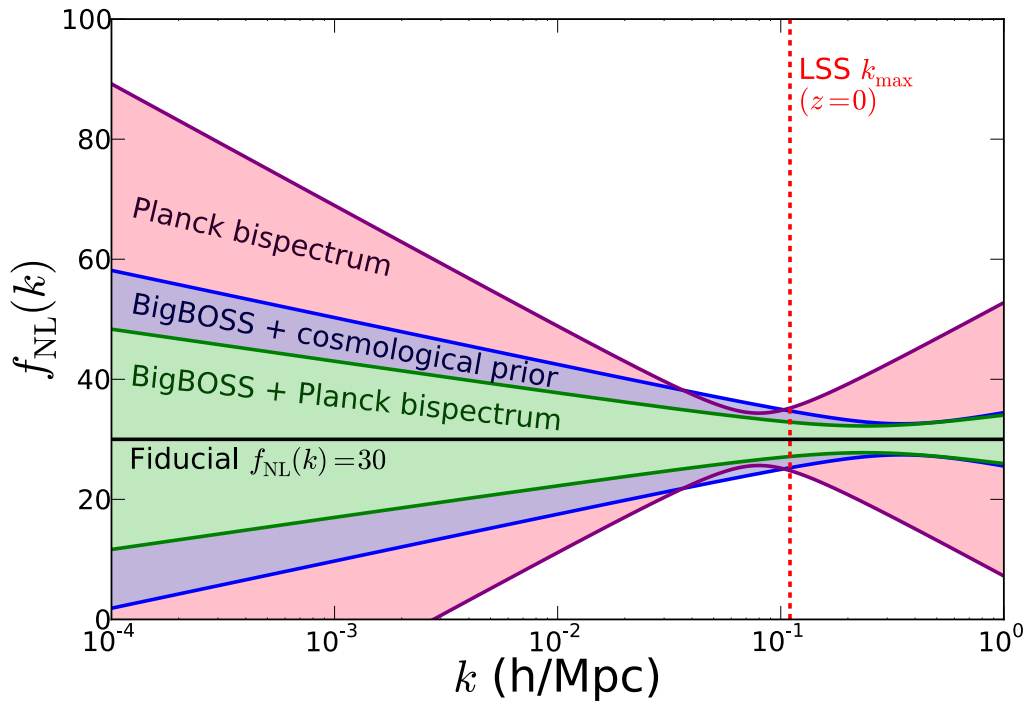
(at least?) as good as, and highly complementary, to Planck CMB

**More general NG models:
beyond f_{NL}**

More generic NG: $f_{\text{NL}}(k)$ forecasts

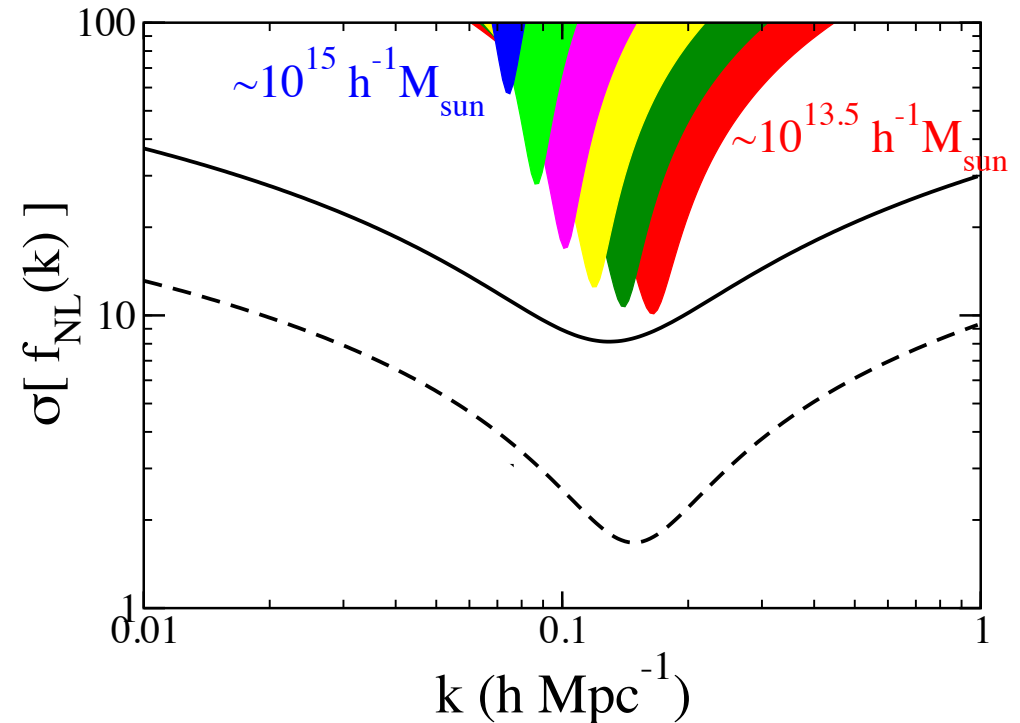
$$f_{\text{NL}}(k) = f_{\text{NL}}^* \left(\frac{k}{k_*} \right)^{n_{f_{\text{NL}}}}$$

CMB and LSS are very complementary



Becker, Huterer & Kadota 2012, Shandera, Dalal & Huterer 2012

Halos of mass M probe
NG on scale $k \sim M^{-1/3}$



In general, LSS can probe:

$$\Delta b_{\text{NG}} \propto \begin{cases} k^{-2} \text{ (local)} \\ k^{-1} \text{ (folded)} \\ k^0 \text{ (equilateral)} \\ k^{-\alpha} \text{ (generic); } 0 \leq \alpha \leq 3 \end{cases}$$

Forecasts for $f_{\text{NL}}(\mathbf{k})$

$$f_{\text{NL}}(\mathbf{k}) = f_{\text{NL}}^* \left(\frac{k}{k_*} \right)^{n_{f_{\text{NL}}}}$$

Projected errors $\sigma(f_{\text{NL}}^*)$ and $\sigma(n_{f_{\text{NL}}})$, and the corresponding pivots

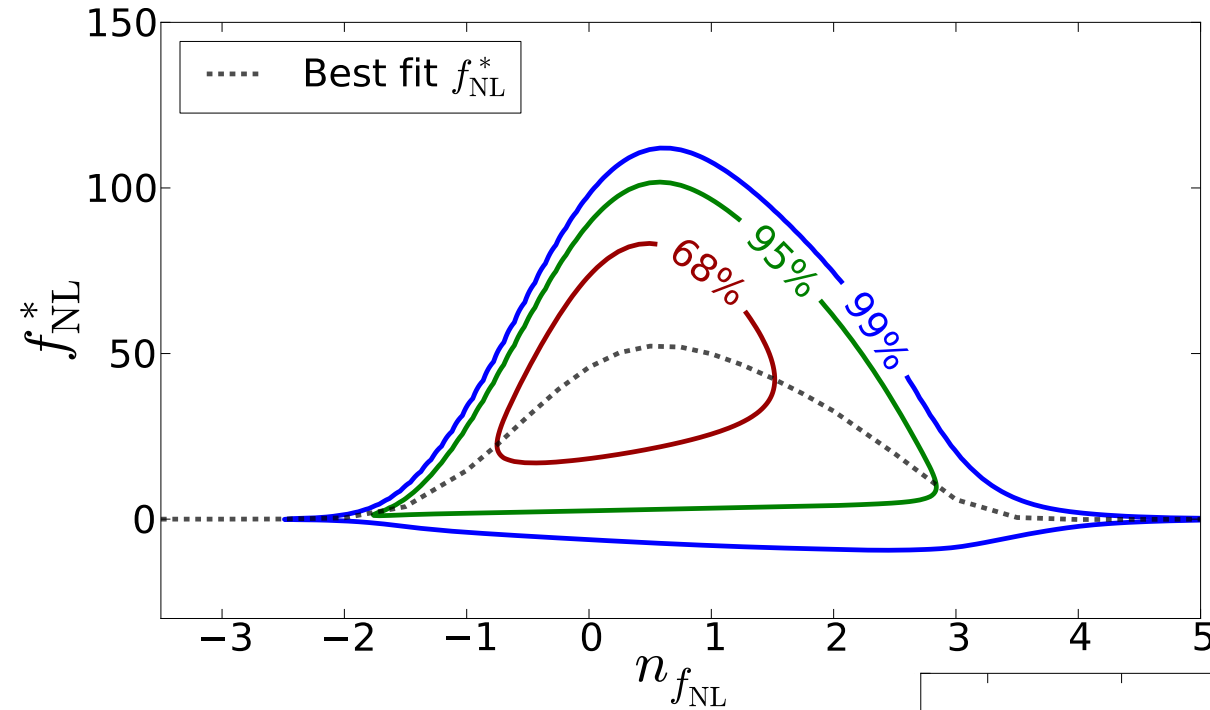
Variable	BigBOSS	BigBOSS+Planck $C_{\ell S}$	Planck bispec	BigBOSS+all Planck
$\sigma(f_{\text{NL}}^*)$	3.0	2.6	4.4	2.2
$\sigma(n_{f_{\text{NL}}})$	0.12	0.11	0.29	0.078
FoM ^(NG)	2.7	3.4	0.78	5.8
k_{piv}	0.33	0.35	0.080	0.24

area in f_{NL}^* - $n_{f_{\text{NL}}}$ plane

**NB: The LSS forecasts are very uncertain,
much more so than the CMB**

First constraints on the running of NG

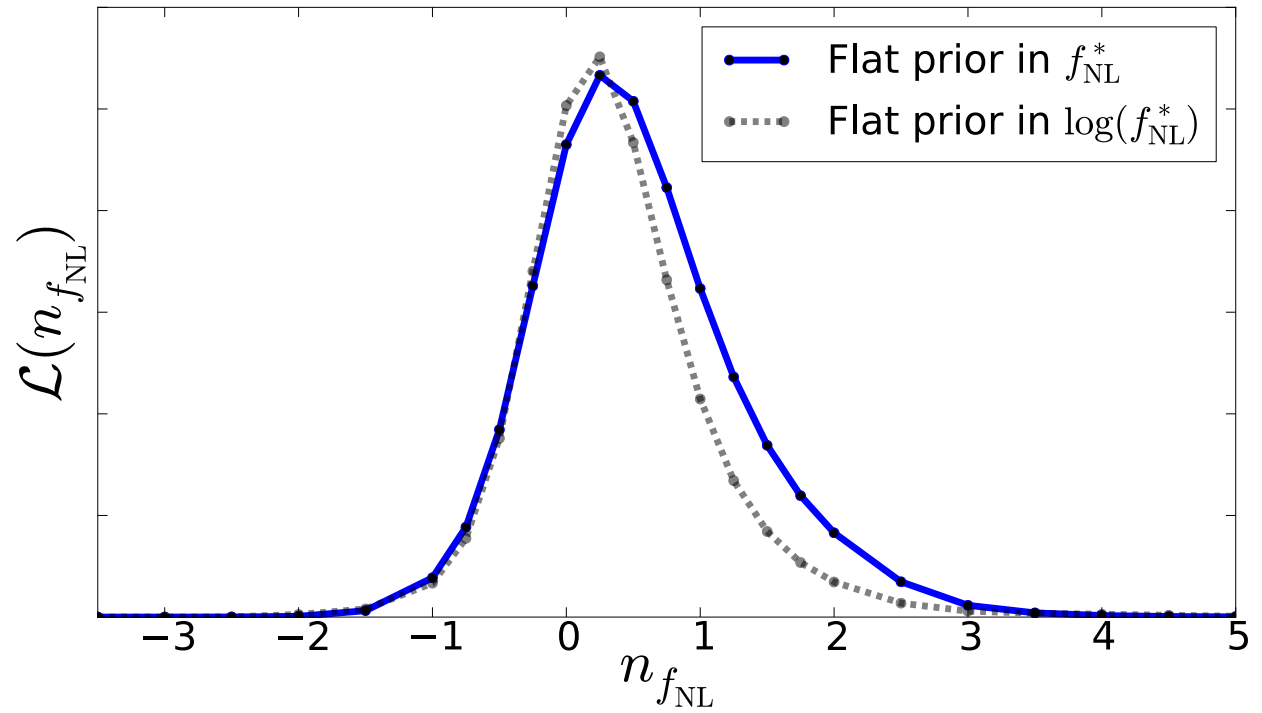
WMAP7 data, modified KSW estimator



$$f_{\text{NL}}(k) = f_{\text{NL}}^* \left(\frac{k}{k_*} \right)^{n_{f_{\text{NL}}}}$$

$$n_{f_{\text{NL}}} = 0.3^{+1.9}_{-1.2}$$

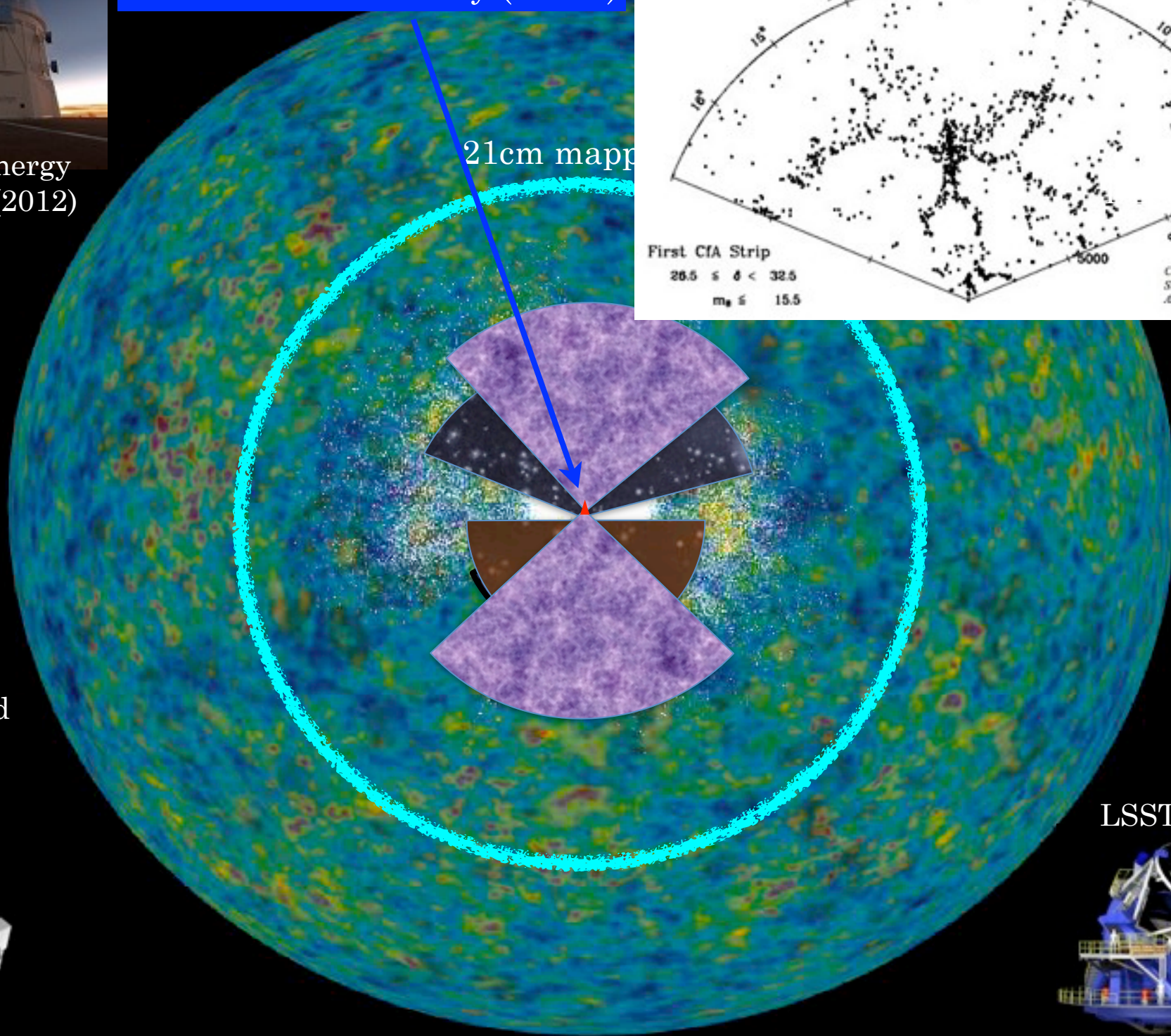
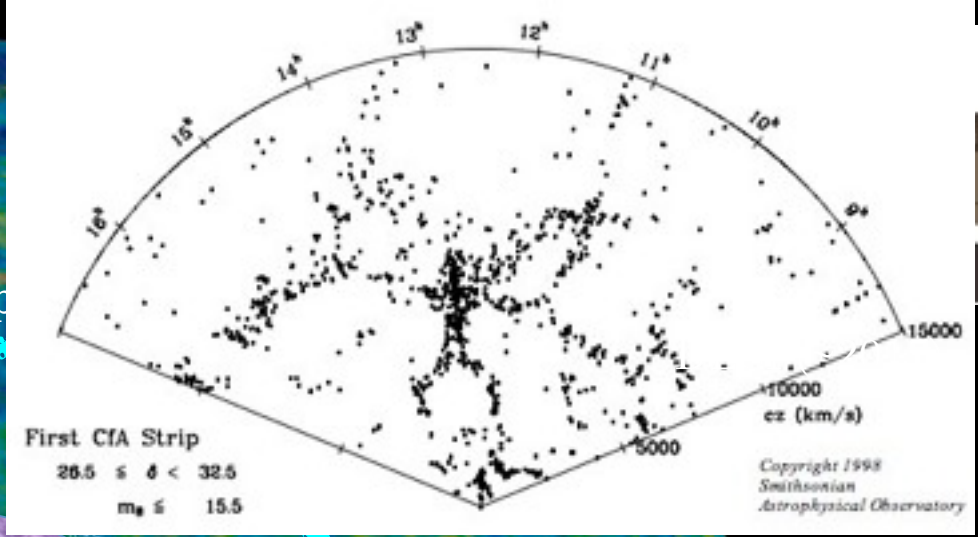
at 95% CL



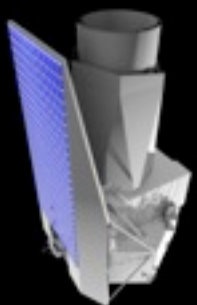
▲ Harvard-Cfa survey (1980s)



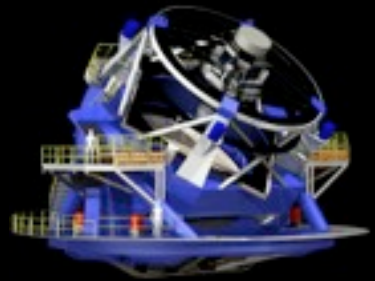
Dark Energy Survey (2012)



Euclid and WFIRST (~202X)



LSST (~2018)



Challenges for NG program

... and approximate **current** status

- Motivate simple and more complicated NG models
✓ (single-field, multiple fields, self-interactions)
- Utilize a variety of observables in LSS and CMB to get at NG ✓
- Develop fast, near-optimal estimators to extract NG from the CMB ✓
- Develop theory to relate NG models to LSS observables ✓ ✗ (messy; still need to check with sims)
- Develop theory to use LSS info from quasi-linear scales ($k \simeq 0.1 h^{-1} \text{ Mpc}$) ✓ ✗
- Control the systematic errors, esp large-scale LSS ✓ ✗
- Use LSS bispectrum to get at primordial NG ✗

Advances in Astronomy special issue on “Testing the Gaussianity and Statistical Isotropy of the Universe”

<http://www.hindawi.com/journals/aa/2010/si.gsiu/>

15 review articles (all also on arXiv)

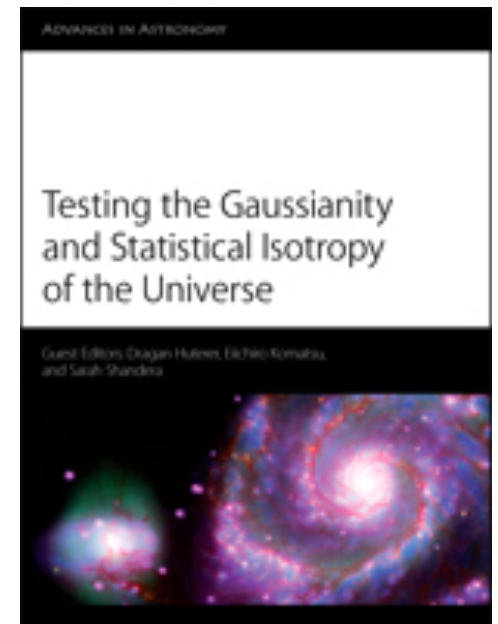
Testing the Gaussianity and Statistical Isotropy of the Universe

Guest Editors: Dragan Huterer, Eiichiro Komatsu, and Sarah Shandera

Non-Gaussianity from Large-Scale Structure Surveys, Licia Verde
Volume 2010 (2010), Article ID 768675, 15 pages

*Non-Gaussianity and Statistical Anisotropy from Vector Field
Populated Inflationary Models*, Emanuela Dimastrogiovanni, Nicola
Bartolo, Sabino Matarrese, and Antonio Riotto
Volume 2010 (2010), Article ID 752670, 21 pages

Cosmic Strings and Their Induced Non-Gaussianities in the Cosmic Microwave Background,



EXTRA SLIDES

Scale-dependent nongaussianity?

Generalized local ansatz

Becker, Huterer & Kadota 2011, 2012
theory motivation: Byrnes et al, etc

- Motivated by multi-field inflationary models
- In general, even if you are considering standard single-field inflation, interactions may lead to scale-dependence of f_{NL}

(Usual) local model...

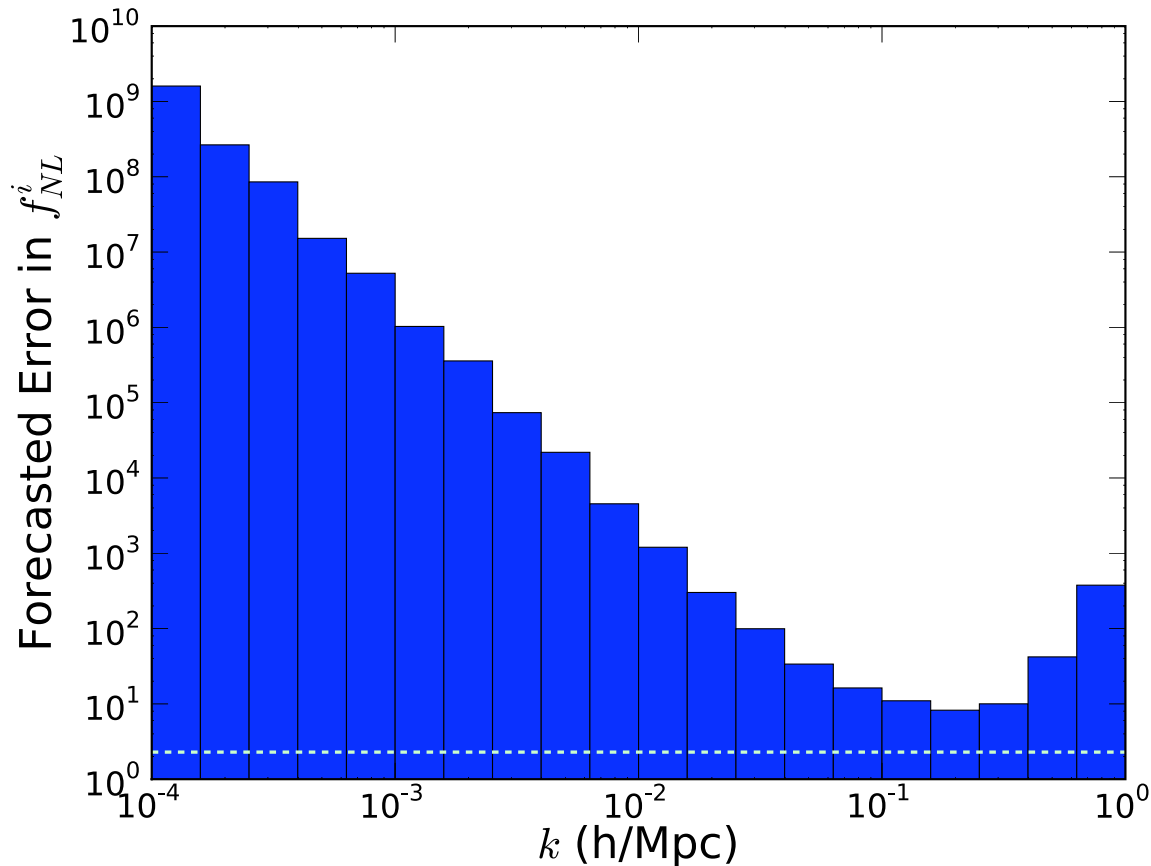
$$\Phi(x) = \phi_G(x) + f_{\text{NL}} [\phi_G^2(x) - \langle \phi_G^2 \rangle]$$

...we generalize to a scale dependent (non-local) model

$$\Phi(x) = \phi_G(x) + f_{\text{NL}}(x) * [\phi_G^2(x) - \langle \phi_G^2 \rangle]$$

$$\Phi(k) = \phi_G(k) + f_{\text{NL}}(k) \int \frac{d^3 k'}{(2\pi)^3} \phi_G(k') \phi_G(k - k')$$

A complete basis for $f_{NL}(k)$: piecewise-constant bins



Measurement forecasts
from
DES-type survey

Given this basis, projecting forecasts onto any
parametrized $f_{NL}(k)$ model is now trivial

Warning, however: theoretical predictions are uncertain and
(always) have to be checked with simulations first

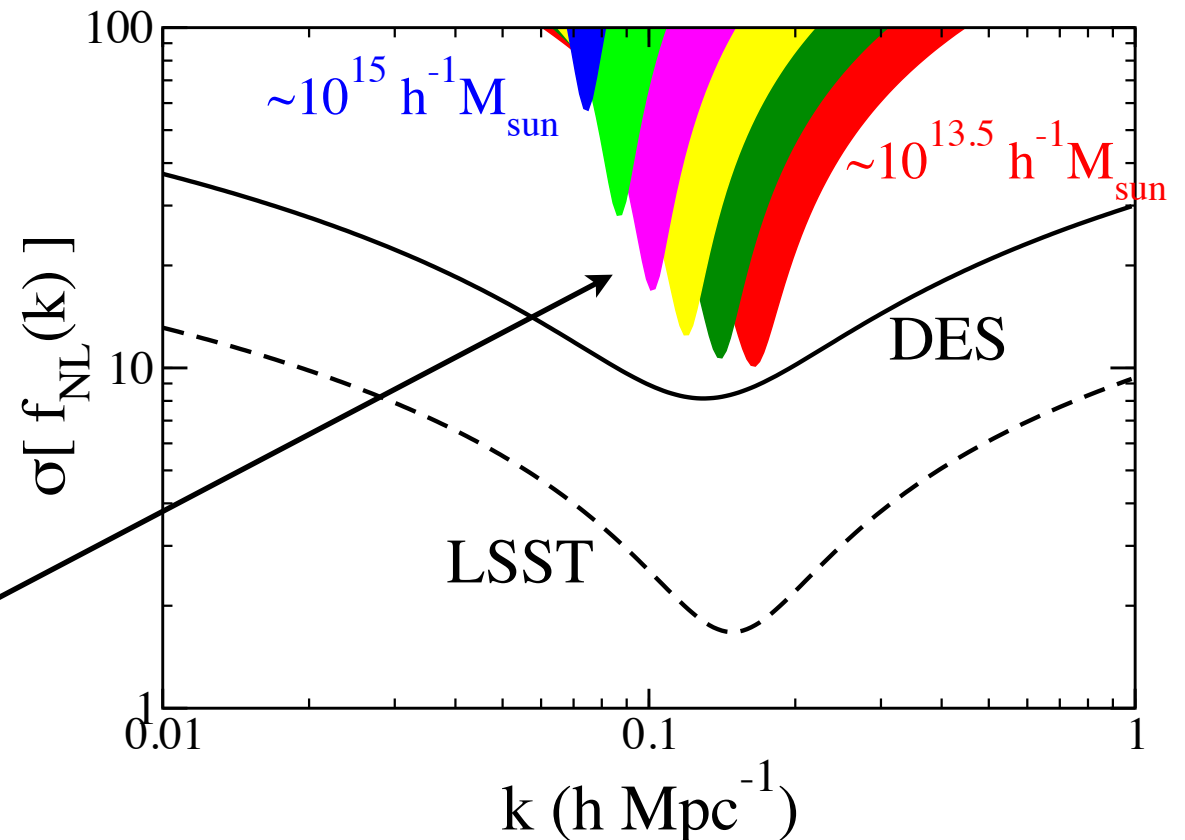
Future: using LSS to probe scale-dependent NG

- ▶ Scale-dep NG models are **motivated by particle theory** (single-field inflation with self-interaction; mixed curvaton-inflaton models)
- ▶ Effects on LSS are significant, but **theory predictions** are uncertain
⇒ ongoing theoretical and simulation work
- ▶ Understanding of **astrophysics** (of DM halos, etc) required in order to probe fundamental physics

Scale-dependent NG ansatz:

$$f_{\text{NL}}(k) = f_{\text{NL}}^* \left(\frac{k}{k_*} \right)^{n_{f_{\text{NL}}}}$$

Halos of mass M probe
NG on scale $k \sim M^{-1/3}$



CMB, LSS, and CMB+LSS forecasts

$$f_{\text{NL}}(k) = f_{\text{NL}}^* \left(\frac{k}{k_*} \right)^{n_{f_{\text{NL}}}}$$

