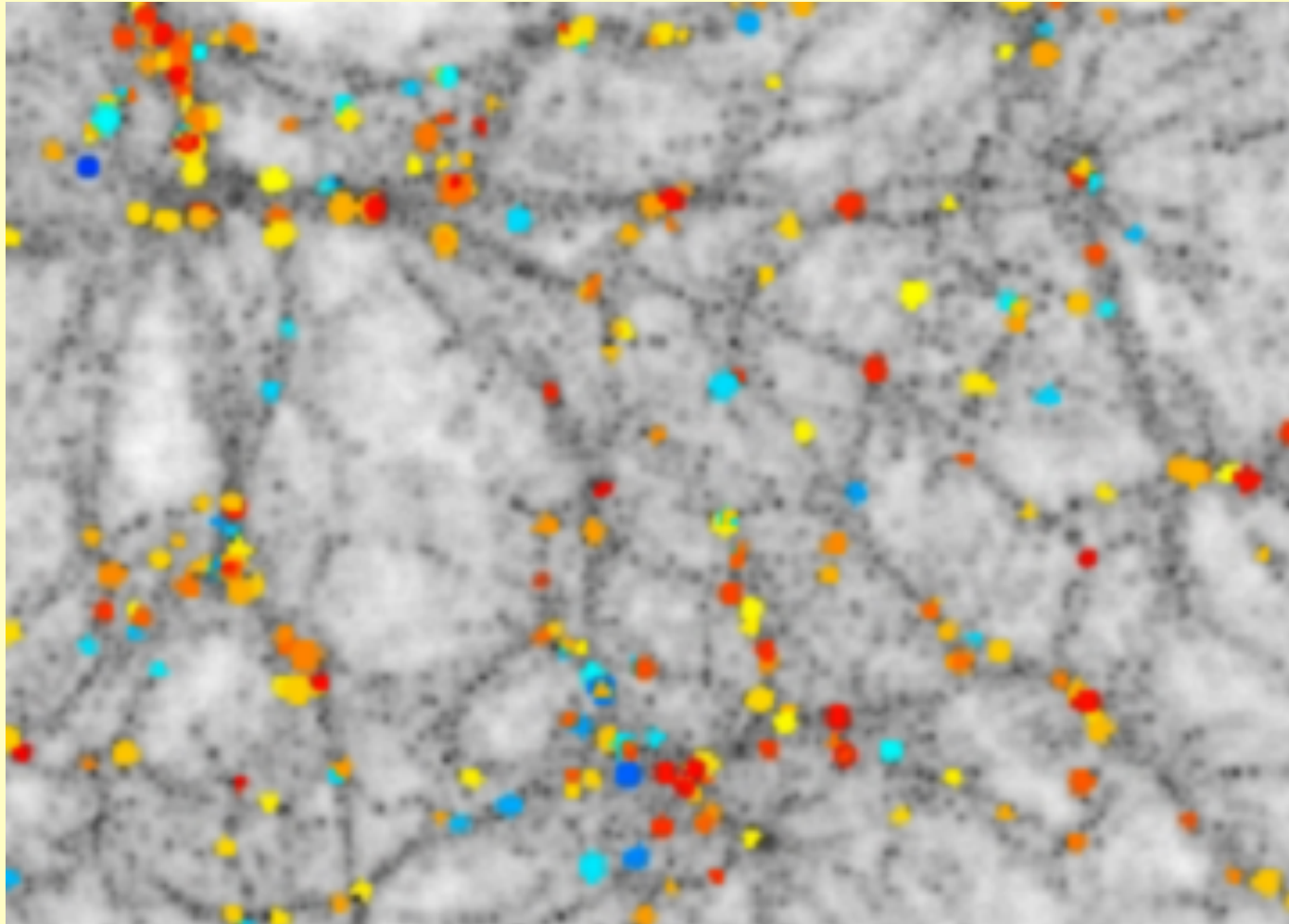


Multiple tracers of cosmological structure



John Peacock

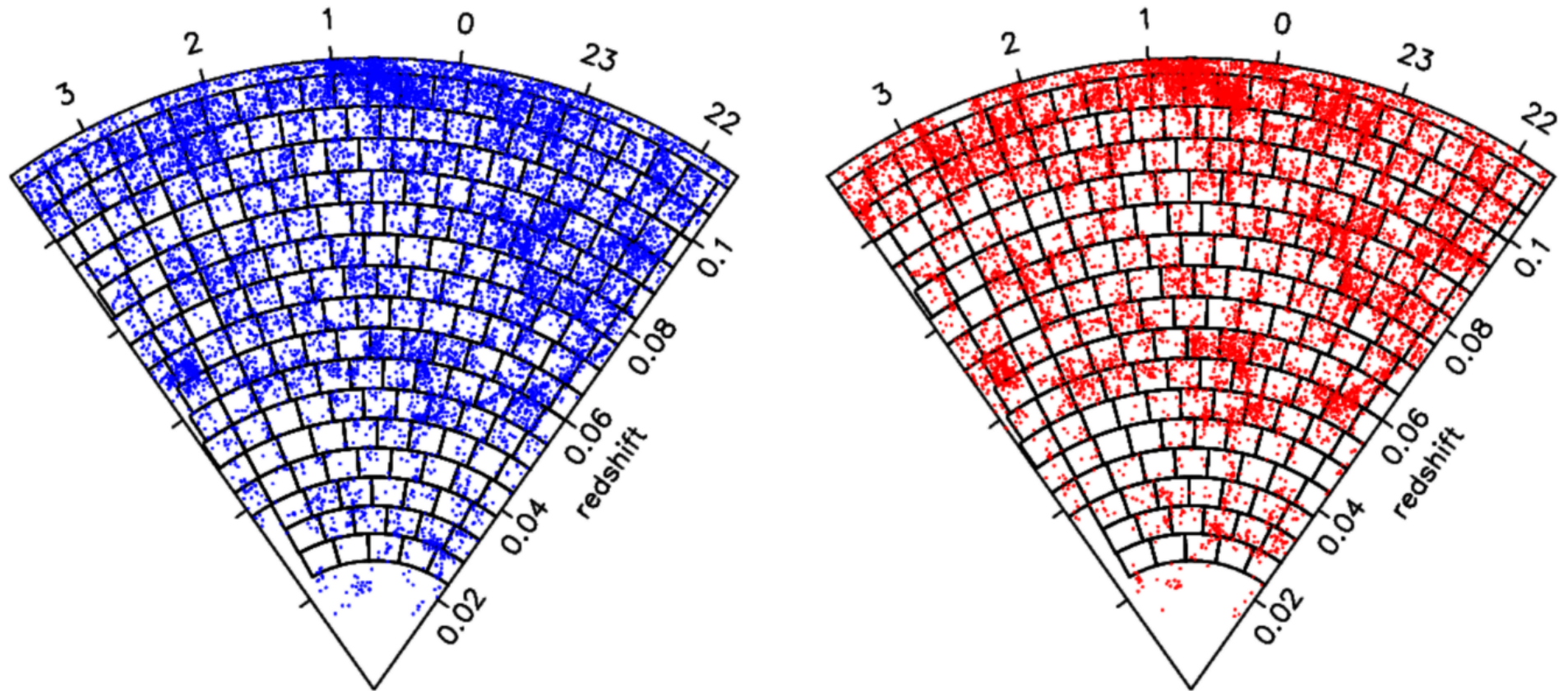
Ripples in the Cosmos: Durham

26 July 2013

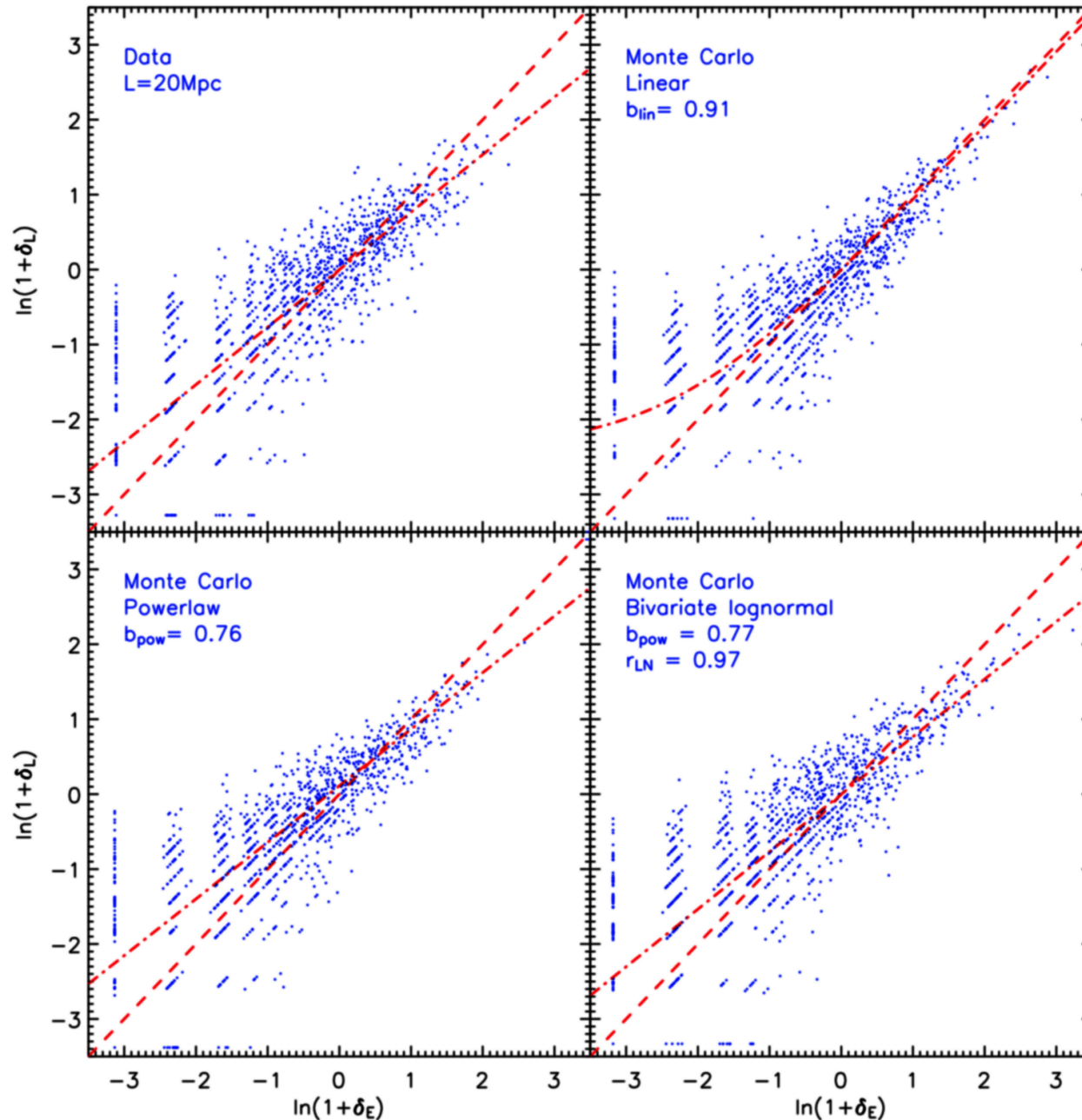
Topics

- Response of different tracers to DM density field
 - Bias, stochasticity and noise
- Information from cross-correlation
 - Exploiting CMB lensing
 - Tomographic photo-z calibration
- Evading cosmic variance
 - GAMA multi-tracer RSD

Stochastic bias in 2dFGRS



Wild et al. (2004): counts-in-cells for different populations



Stochastic bias in 2dFGRS

Extra scatter in blue vs red overdensity beyond Poisson:

$r = 1 - \text{few \% correlation in density}$

Stochastic bias and discreteness

$$\delta_t(\mathbf{x}) = f(\delta[\mathbf{x}]) + \epsilon(\mathbf{x})$$

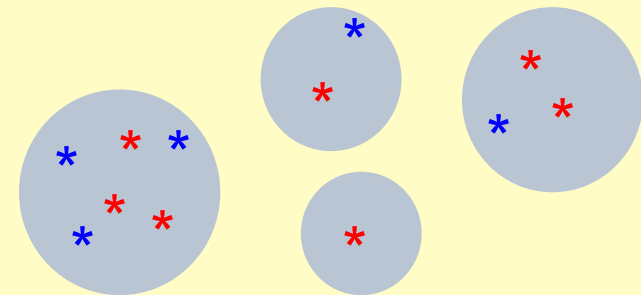
– e.g. $f(\delta) = b_0 + b_1\delta + b_2\delta^2/2$ or $1 + \delta_t \propto (1 + \delta)^b$

$$\Rightarrow P_t(k) = b^2(k) P(k) + P_\epsilon + P_{\text{shot}}; \quad P_{\text{shot}} = 1/n$$

Halo model interpretation:

ϵ field arises mainly from discrete nature of haloes (like shot noise); **non-local bias**. see Baldauf et al. 1305.2917

But overall density field is sum of haloes, so in principle no shot noise if haloes are weighted by mass (Seljak, Hamaus & Desjacques 0904.2963)



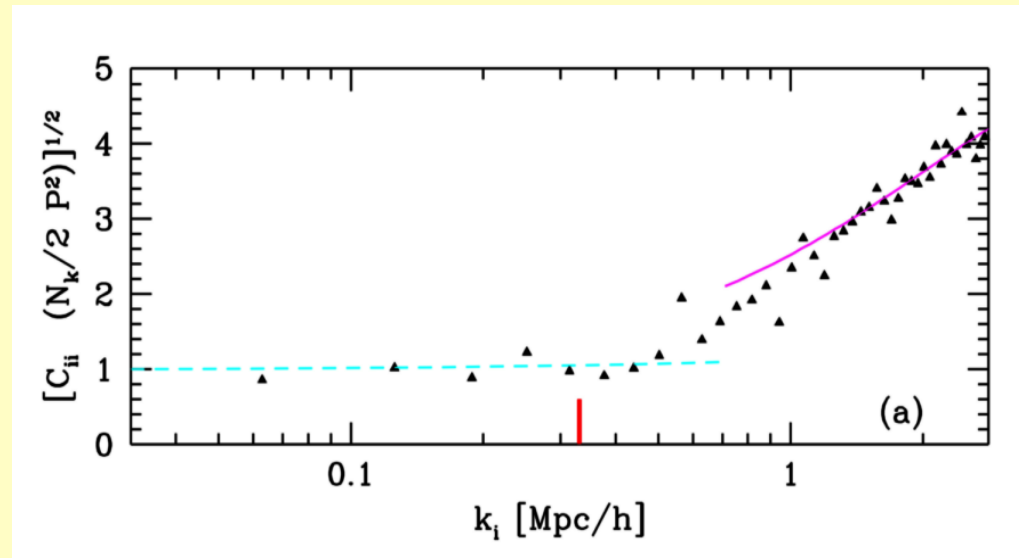
Cosmic variance

$$\text{Gaussian field: } \delta \ln P = N_{\text{modes}}^{-1/2} = \left(\frac{V}{4\pi^2} k^3 d \ln k \right)^{-1/2}$$

$$\begin{aligned} \text{Halo model: } P_{\text{nonlin}} &= P_{\text{lin}} + P_{1\text{-halo}} [+P_{\text{shot}}] \\ &\simeq P_{\text{lin}} + 1/n_{\text{halo}} [+1/n_{\text{galaxy}}] \end{aligned}$$

$$\begin{aligned} \Rightarrow \text{Var}(\ln P) &\simeq 1/N_{\text{modes}}(k) + 1/N_{\text{halo}} \\ &= \text{Var}(\ln P)_{\text{Gaussian}} \times \left(1 + \frac{1}{2} \Delta^2(k) d \ln k \right) \end{aligned}$$

Hence enhanced covariance (and lost signal) beyond NL scale: Scoccimarro et al. (1999); Meiksin & White (1999)



So ultimate limit to precision is number of linear modes?

Cross-correlations

- Natural quantity to consider if you have >1 tracer
- Cross-power affected by cosmic variance in same way as auto-power
- Informative about e.g. stochasticity

$$P_{12} \equiv \delta_1(k)\delta_2^*(k); \quad r(k) = P_{12}/[P_{11}(k)P_{22}(k)]^{1/2}$$

- Interesting tomographic information from samples with partial overlap

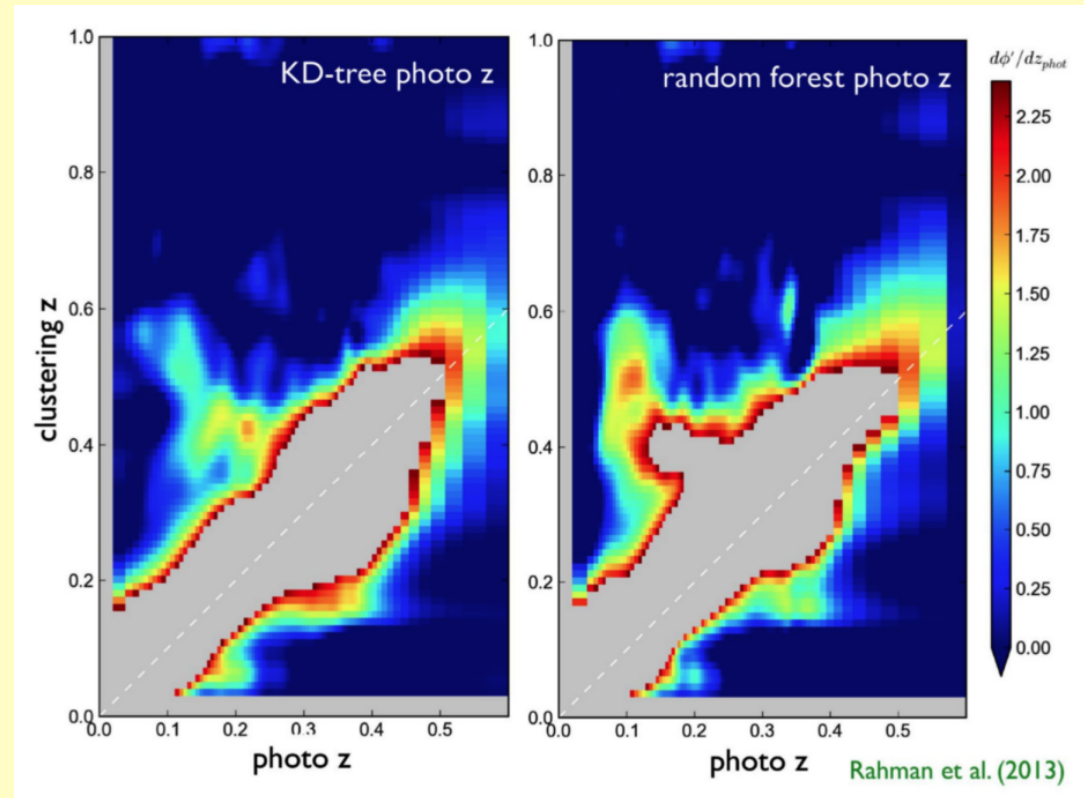
Cross-correlation with photo-z's

$$w_{pt}(\theta, z_p) = \int \frac{b_p(z)}{b_t(z)} \xi_t(r[z, z_p, \theta]) \phi(z|z_p) dz$$

So measure $\phi(z|z_p)$ if bias evolution is known

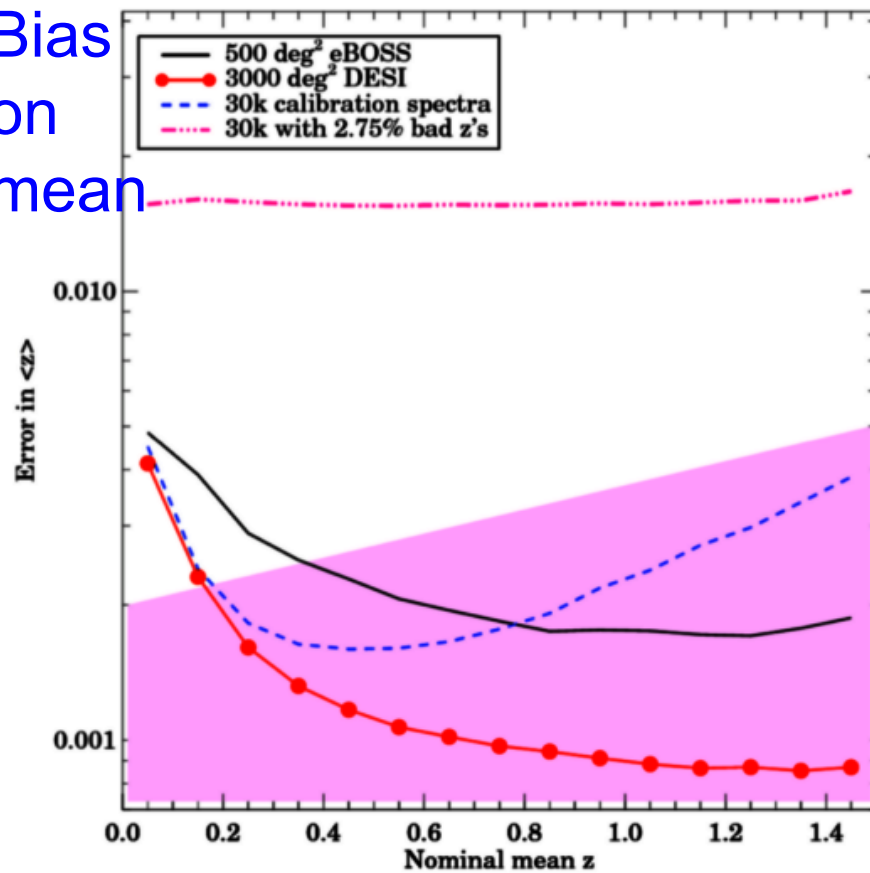
Newman (0804.1409); Schulz (0910.3683); McQuinn & White (1302.0857)

Use bright tracers to calibrate deep photo-z's since complete spectroscopy to requisite depth may be infeasible

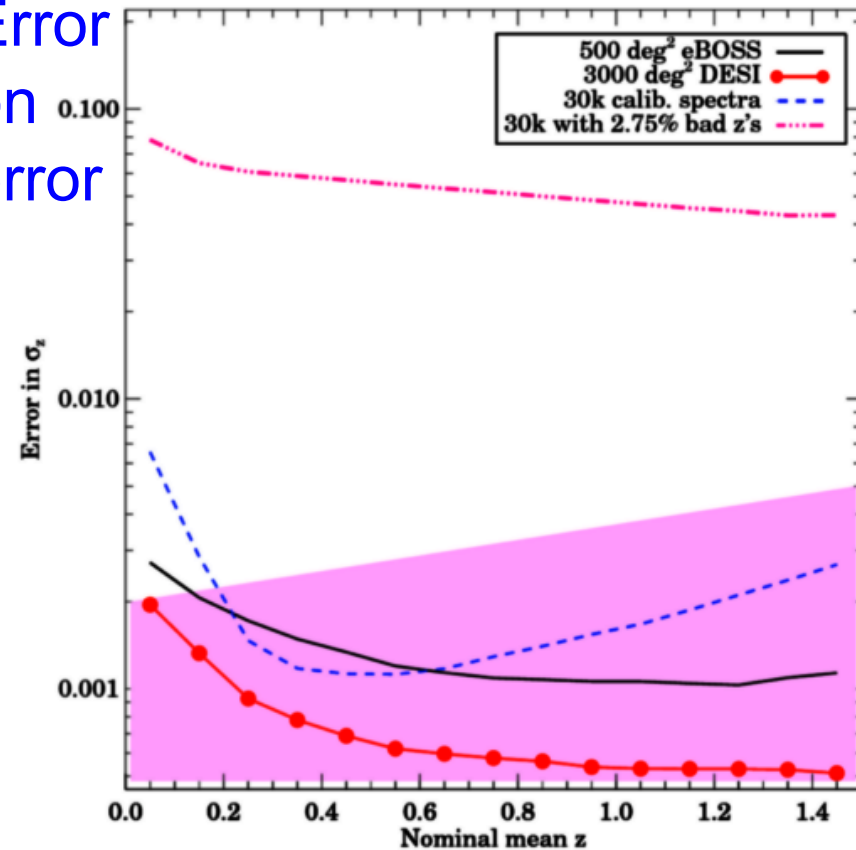


Prospects for LSST

Bias
on
mean

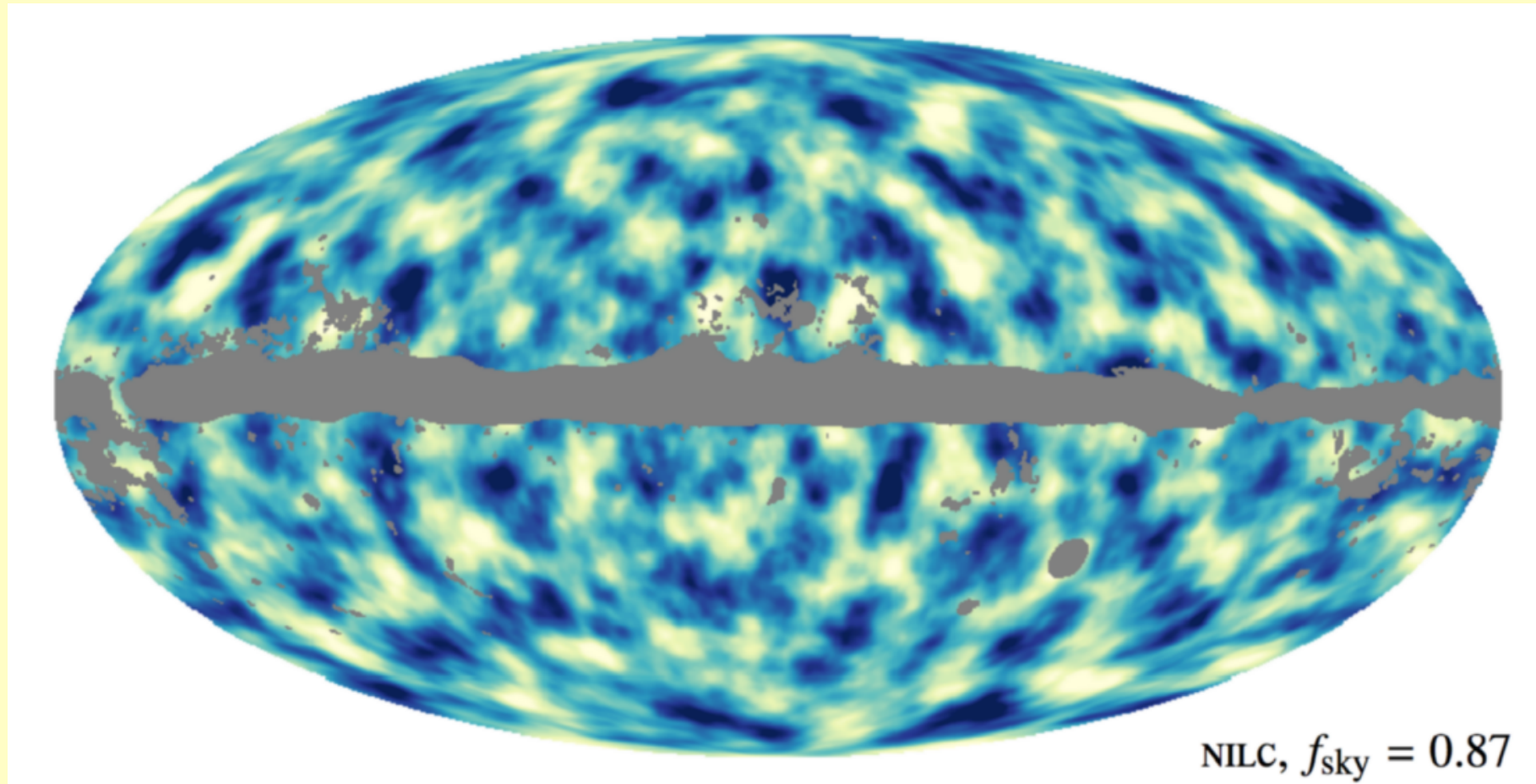


Error
on
error



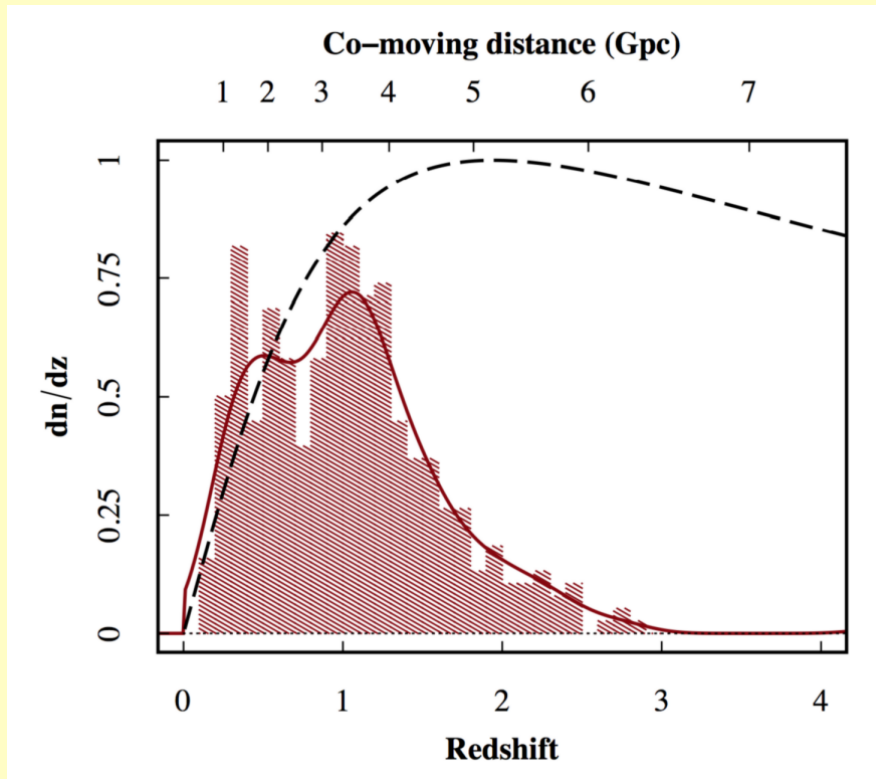
Newman et al. Snowmass 2013

Deprojecting CMB lensing



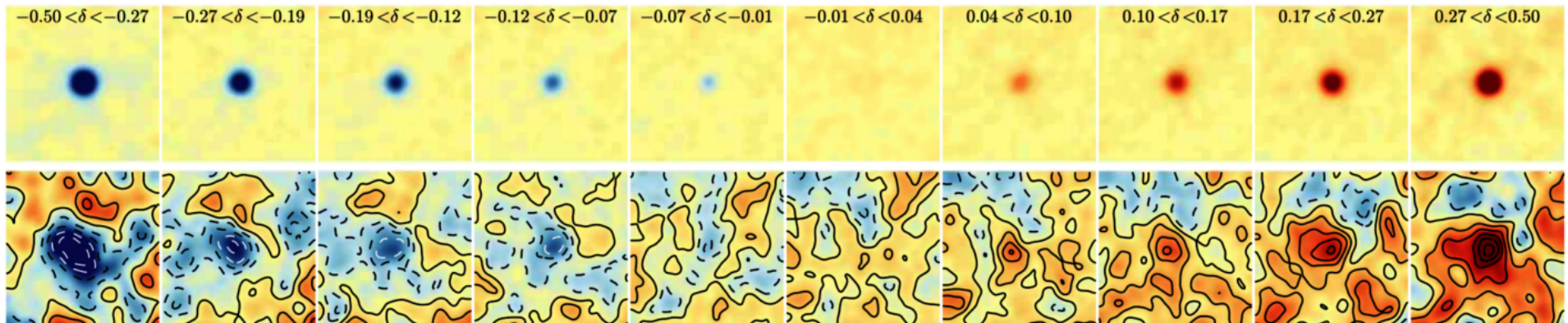
Lensing potential: Planck 2013 paper 17

Deprojecting CMB lensing



Geach et al. 1307.1706:
cross-correlate SPT lensing
with WISE-selected
quasars. Measure $b = 1.67$
 ± 0.24

Strong calibration test for
tomographic weak lensing



Beating cosmic variance

McDonald & Seljak 0810.0323. See also 1003.3238

In limit of linear bias, all populations share same fluctuations in phase and mode amplitude:

$$P(> |\delta|) = \exp(-|\delta|^2 / \langle |\delta|^2 \rangle); \text{ random phase}$$

Fourier Configuration space:

$$\delta_1 = b_1 \delta, \delta_2 = b_2 \delta \Rightarrow b_2/b_1 = |\delta_2|/|\delta_1|$$

Fourier Redshift space:

$$\delta_1 = (b_1 + f\mu^2)\delta, \delta_2 = (b_2 + f\mu^2)\delta \Rightarrow (b_2 + f\mu^2)/(b_1 + f\mu^2) = |\delta_2|/|\delta_1|$$

So can access RSD **without limit from mode amplitude and phase** if two b factors are different

All the information is in power

$$P_1 = A\bar{P}(1 + \alpha) + \Delta P_1$$

$$P_2 = B\bar{P}(1 + \alpha) + \Delta P_2,$$

Since $\hat{A} = P_1/\bar{P}$ etc., the covariance matrix is

$$\sigma_{AA}^2 = A^2\sigma^2 + \epsilon_1^2/\bar{P}^2$$

$$\sigma_{BB}^2 = B^2\sigma^2 + \epsilon_2^2/\bar{P}^2$$

$$\sigma_{AB}^2 = AB\sigma^2,$$

where $\sigma^2 \equiv \langle \alpha^2 \rangle$, $\epsilon^2 \equiv \langle (\Delta P)^2 \rangle$. From this, we get

$$\text{var}(A/B) = (A/B)^2 [\epsilon_1^2/A^2\bar{P}^2 + \epsilon_2^2/B^2\bar{P}^2].$$

Cross-correlation (i.e. tracers have same phase) gives no extra information

Getting errors on parameters

Covariance matrix for parameters λ and λ' comes from inverse of **Hessian matrix**:

$$\sigma_{\lambda\lambda'}^2 = (H^{-1})_{\lambda\lambda'}; \quad H_{\lambda\lambda'} \equiv -\frac{\partial^2 \ln \mathcal{L}}{\partial \lambda \partial \lambda'}; \quad \ln \mathcal{L} = D^\dagger \cdot C \cdot D/2$$

Fisher matrix is expectation of Hessian (for error forecasting). For Gaussian likelihood:

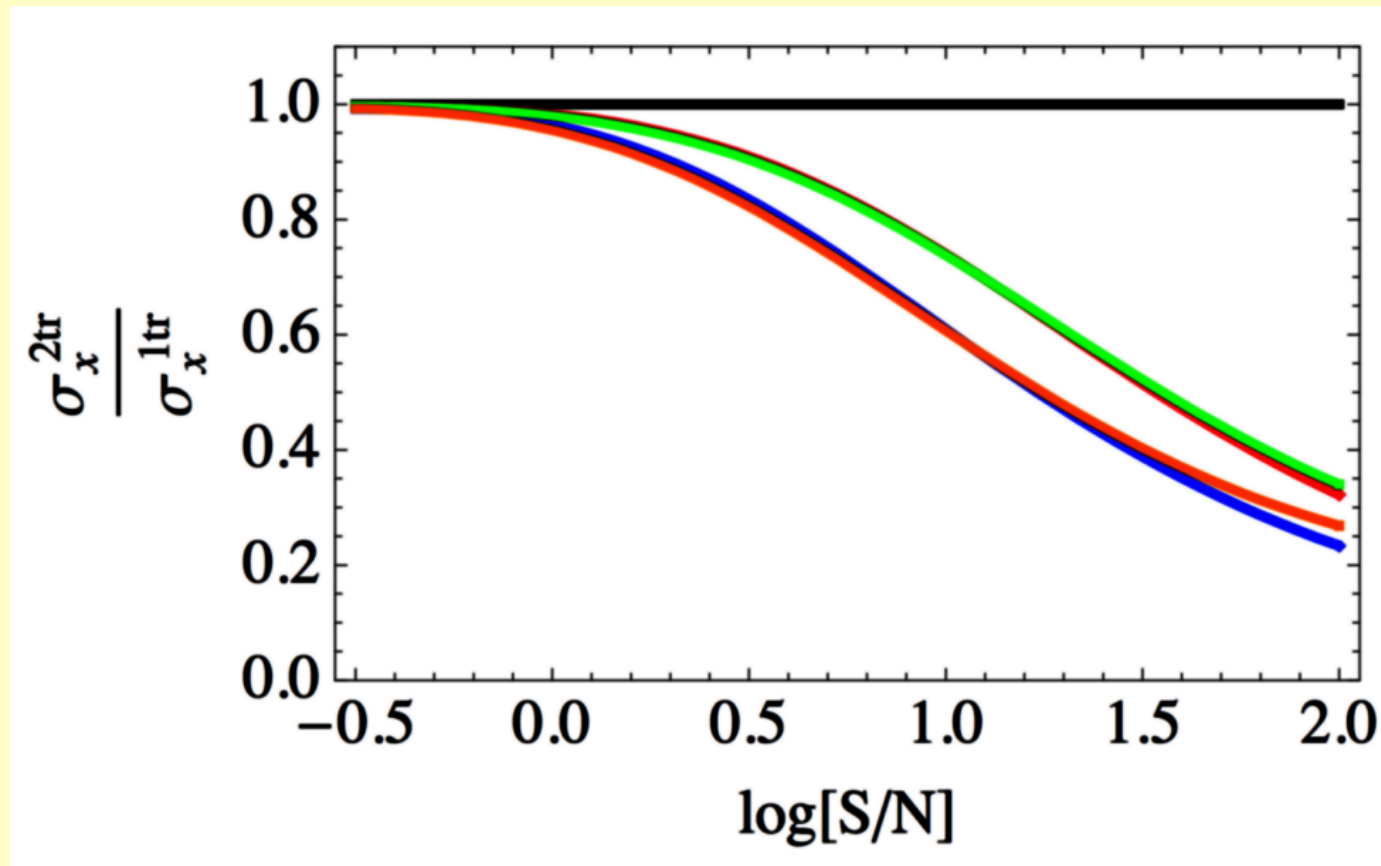
$$F_{\lambda\lambda'} = \frac{1}{2} \text{Tr} (C_{,\lambda} C^{-1} C_{,\lambda'} C^{-1})$$

What goes in D (the data)? It doesn't matter

Thus easier in practice to use covariance of density, $C_{ij} = \langle \delta_1 \delta_2^* \rangle$, rather than covariance of power – just need power spectrum itself.

Gain as f(SNR)

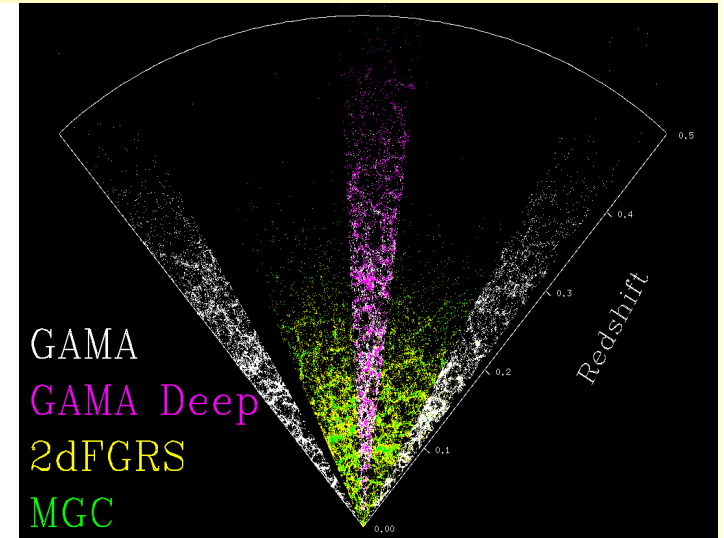
Extent of 2-tracer gain is limited by relative size of shot noise: $\text{SNR} = P_g / P_{\text{shot}}$



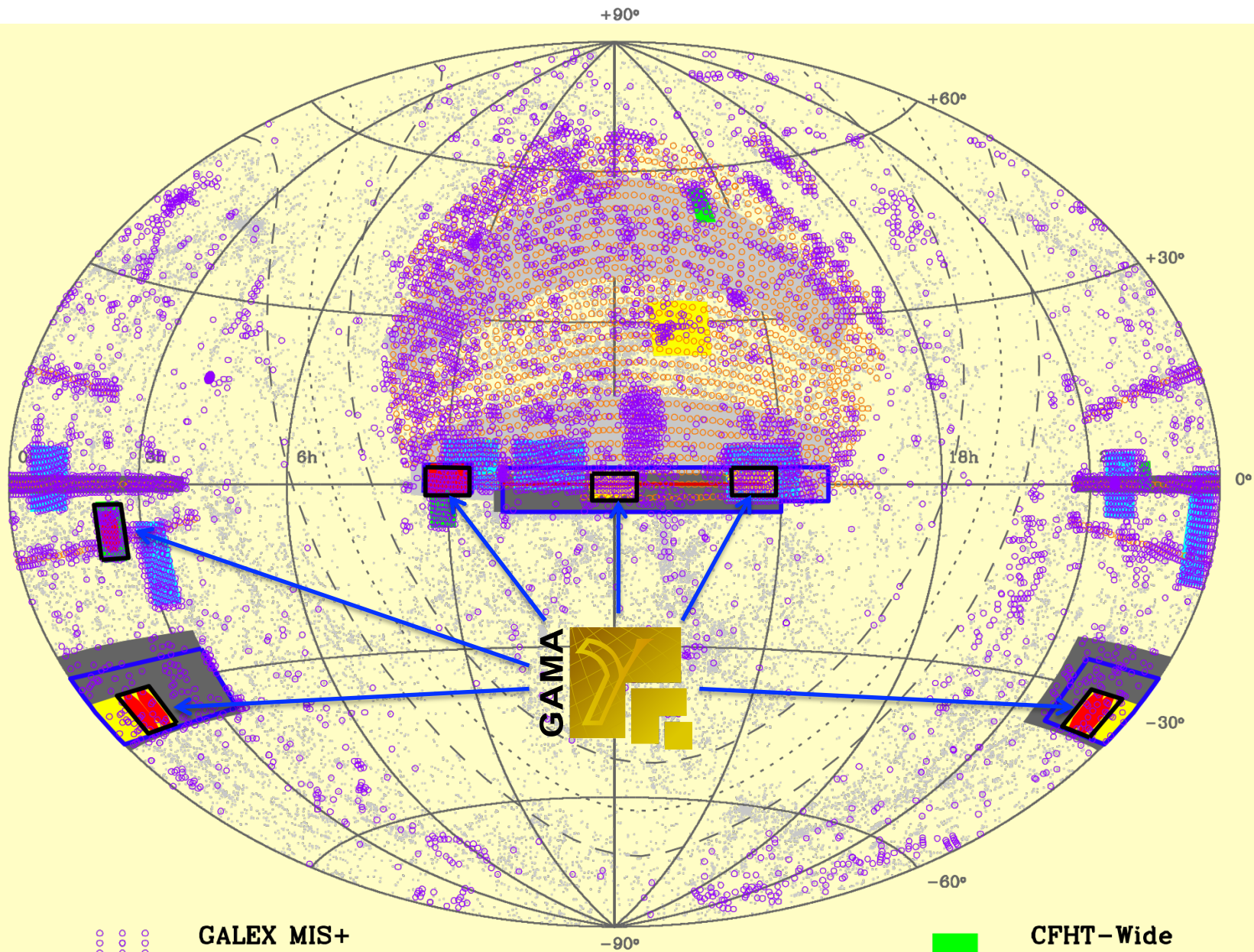
Gil-Marin et al. 1003.3238: $b_1=1$, b_2 varying












First application: GAMA

Galaxy And Mass Assembly – GAMA

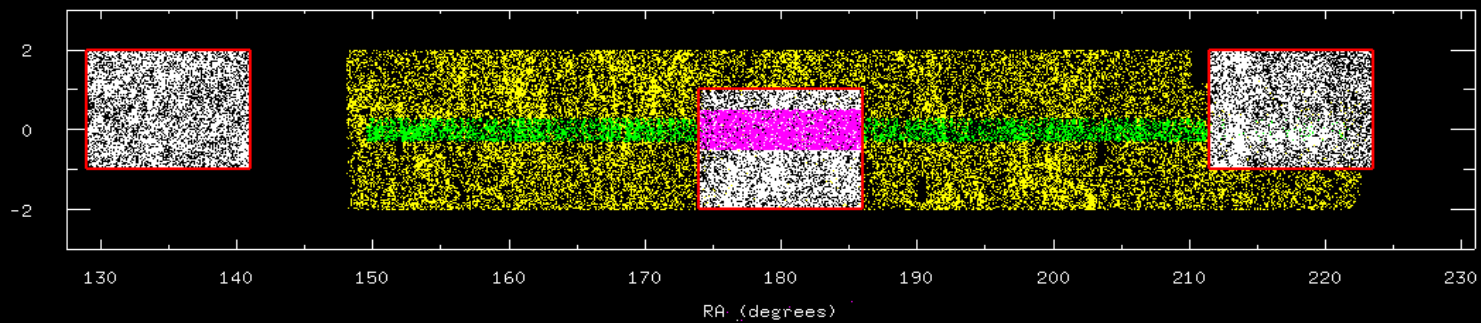


- 300 deg² in 6 fields
- to $r < 19.4 / 19.8$ (GAMA deep) – cf. SDSS 17.8
- Aim for >200,000 redshifts
- First 3 observing seasons (GAMA I; 150 deg²):
 - 63 AAT nights 08/09/10 – 75% clear
 - 140k new z's; 96% success
 - Over 160k including 2dFGRS/SDSS
 - GAMA-II 2011-14; should reach 400k z's



- | | | | | | |
|---|-------------------------------|---|-------------------------------|---|-------------------|
|  | GALEX MIS+ |  | ASKAP-DINGO |  | CFHT-Wide |
|  | GAMA |  | SDSS-Main (spec. only) |  | 2dFGRS |
|  | HERSCHEL-ATLAS |  | WiggleZ |  | UKIDSS-LAS |
|  | Millennium Galaxy Cat. |  | VST-KIDS/VISTA VIKING | | |

Dec (degrees)

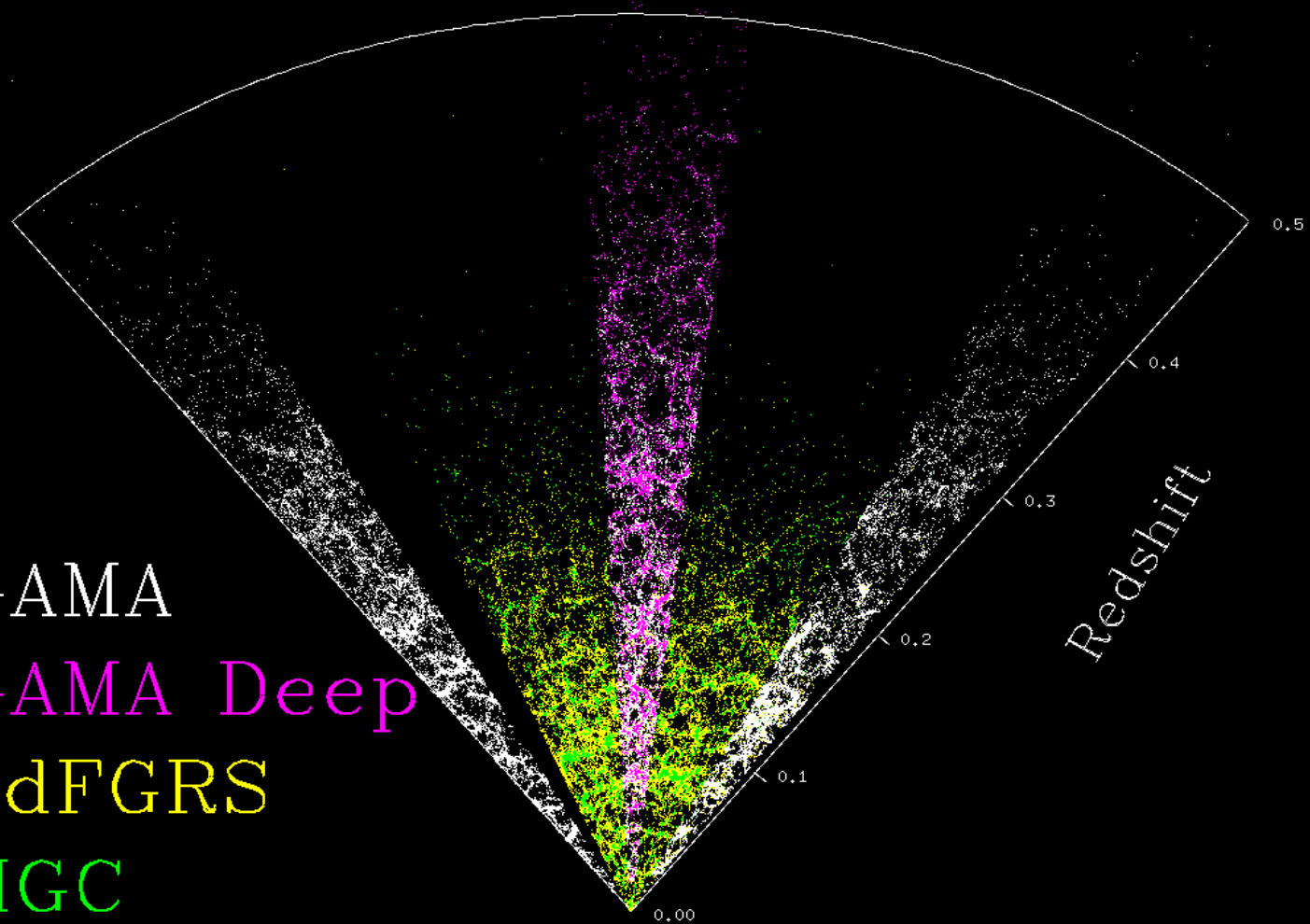


GAMA

GAMA Deep

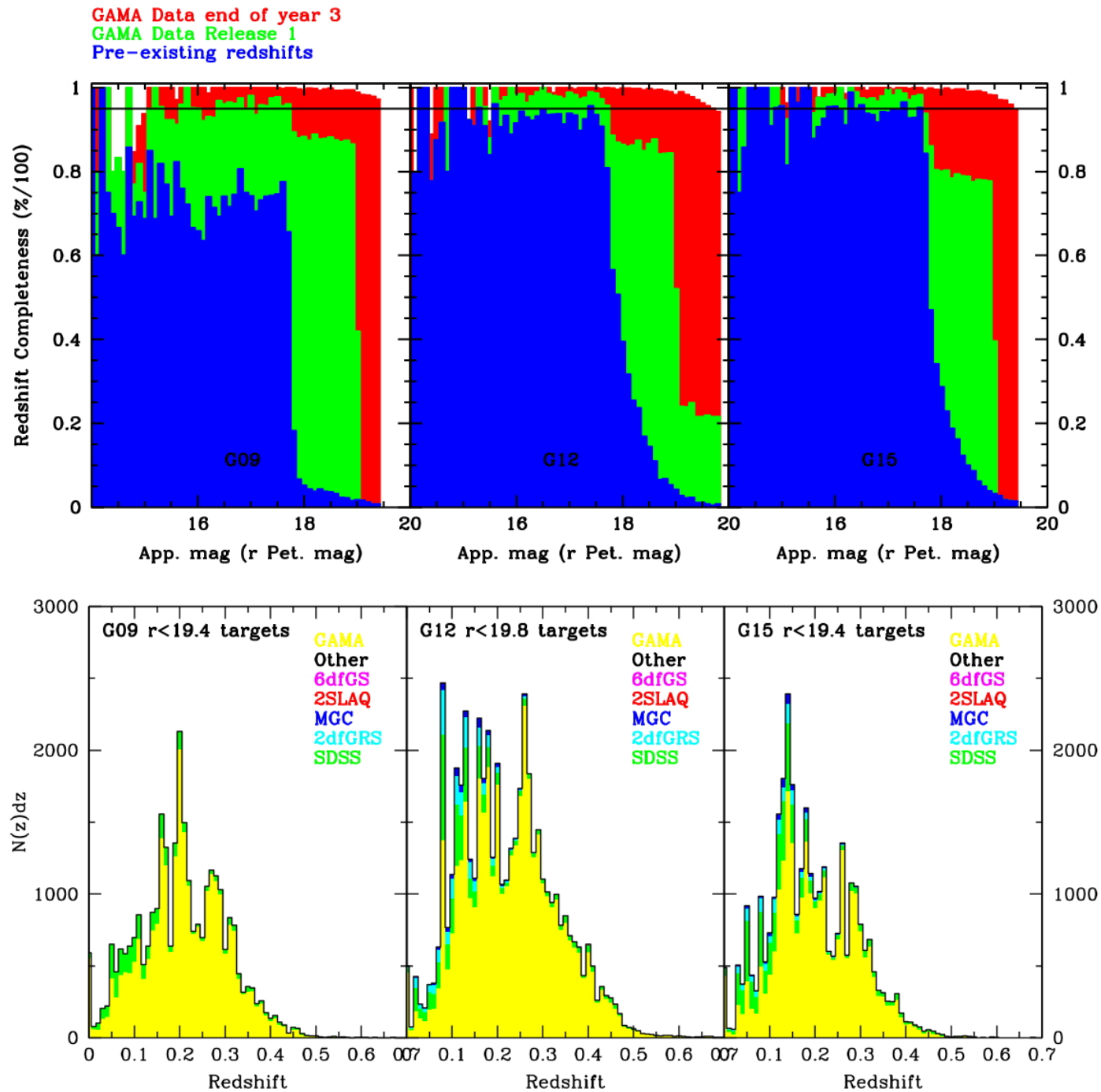
2dFGRS

MGC

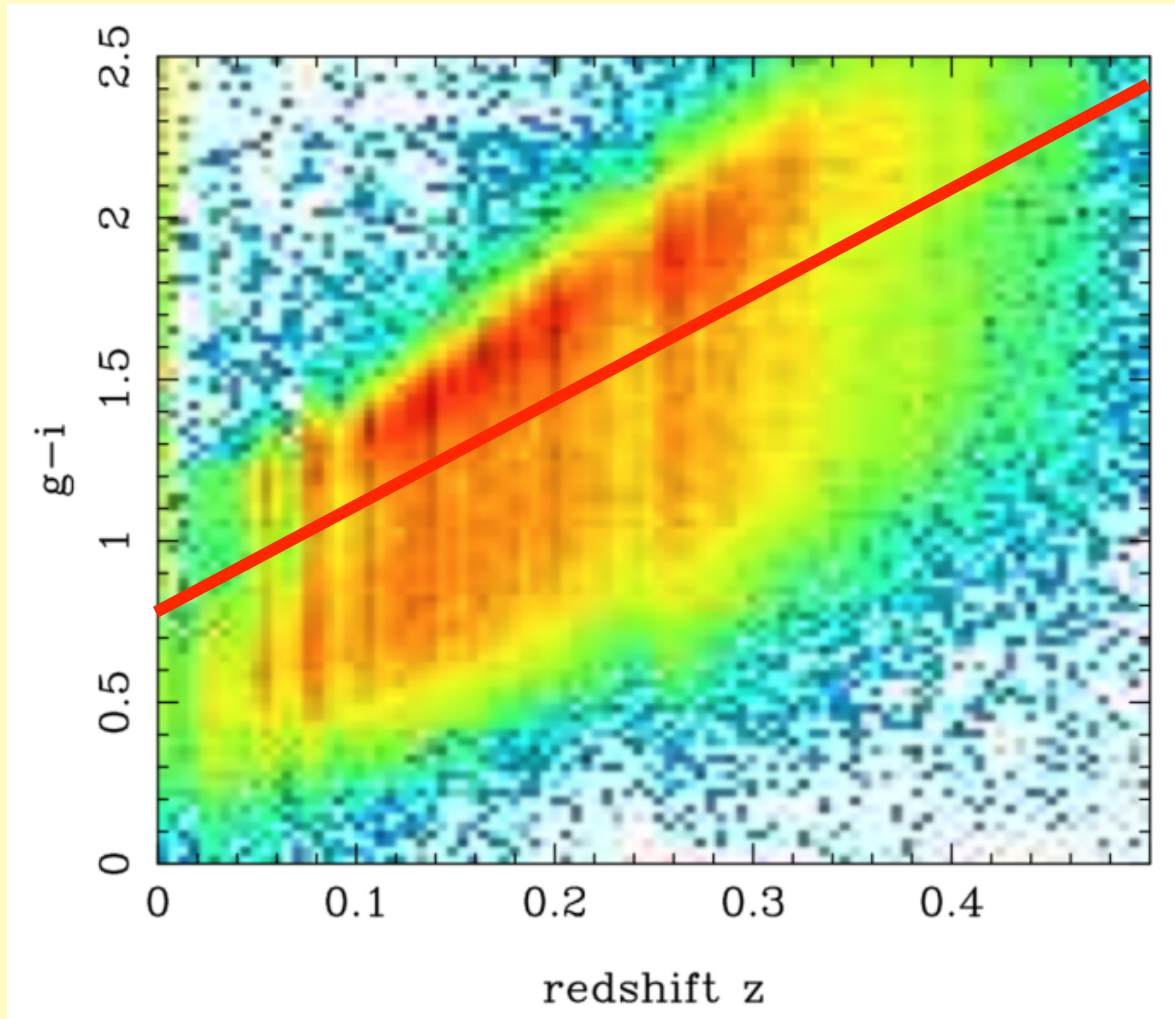


A very complete z-survey: (>96%)

(sampling, not success rate)

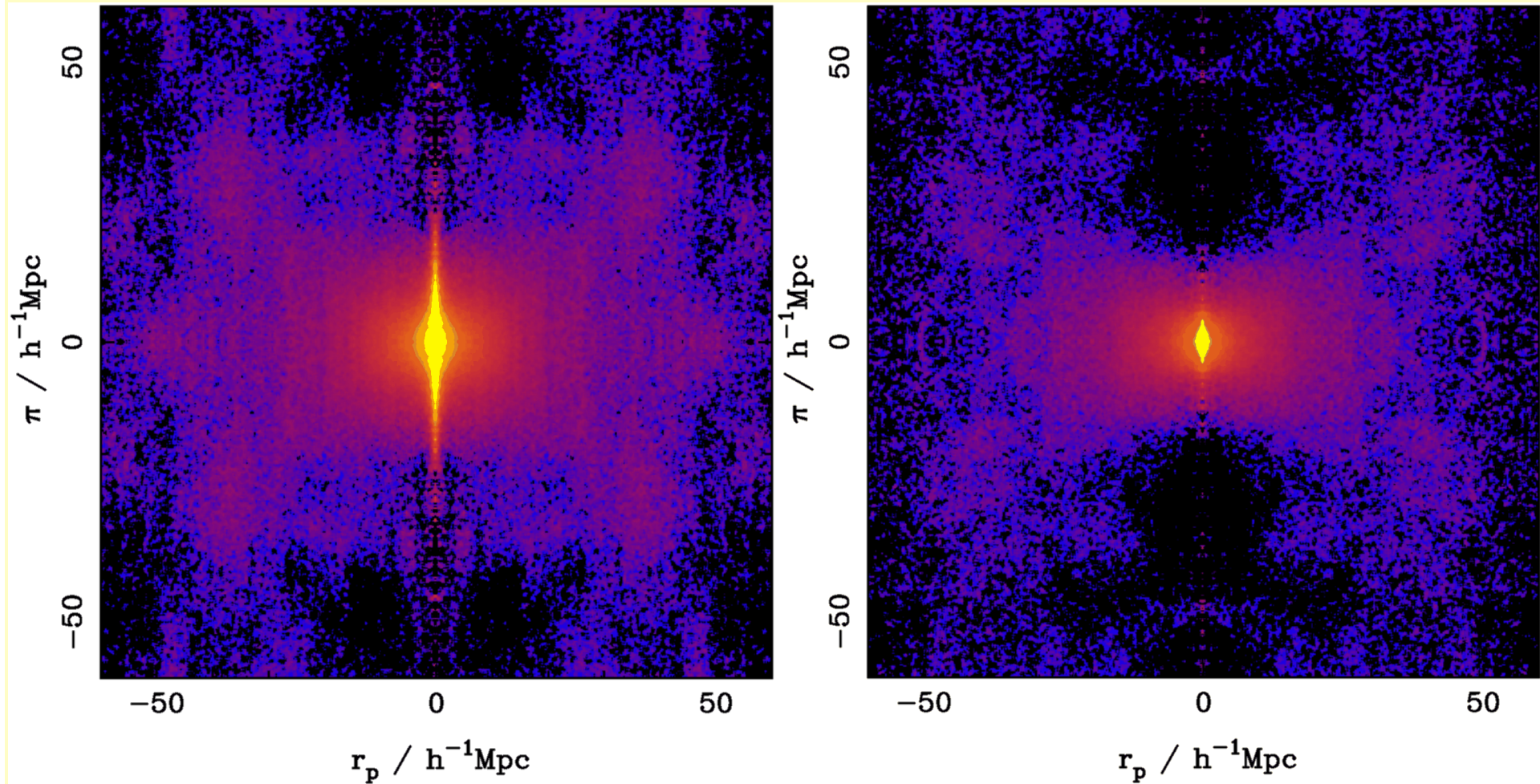


GAMA colour split



Split at $g-i = 0.8 + 3.2z$

RSD as $f(\text{colour})$ in Fourier space (not)

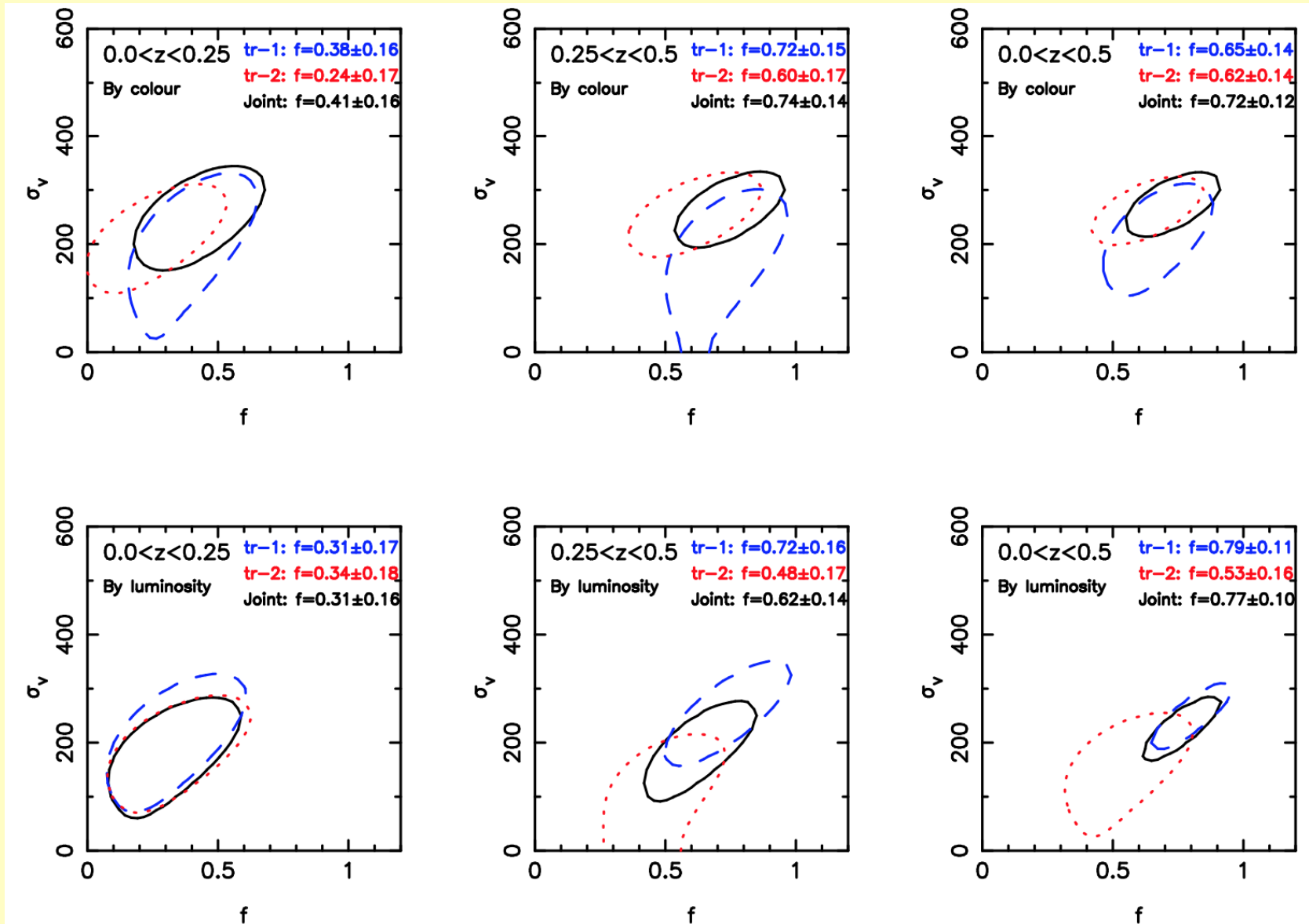


Red

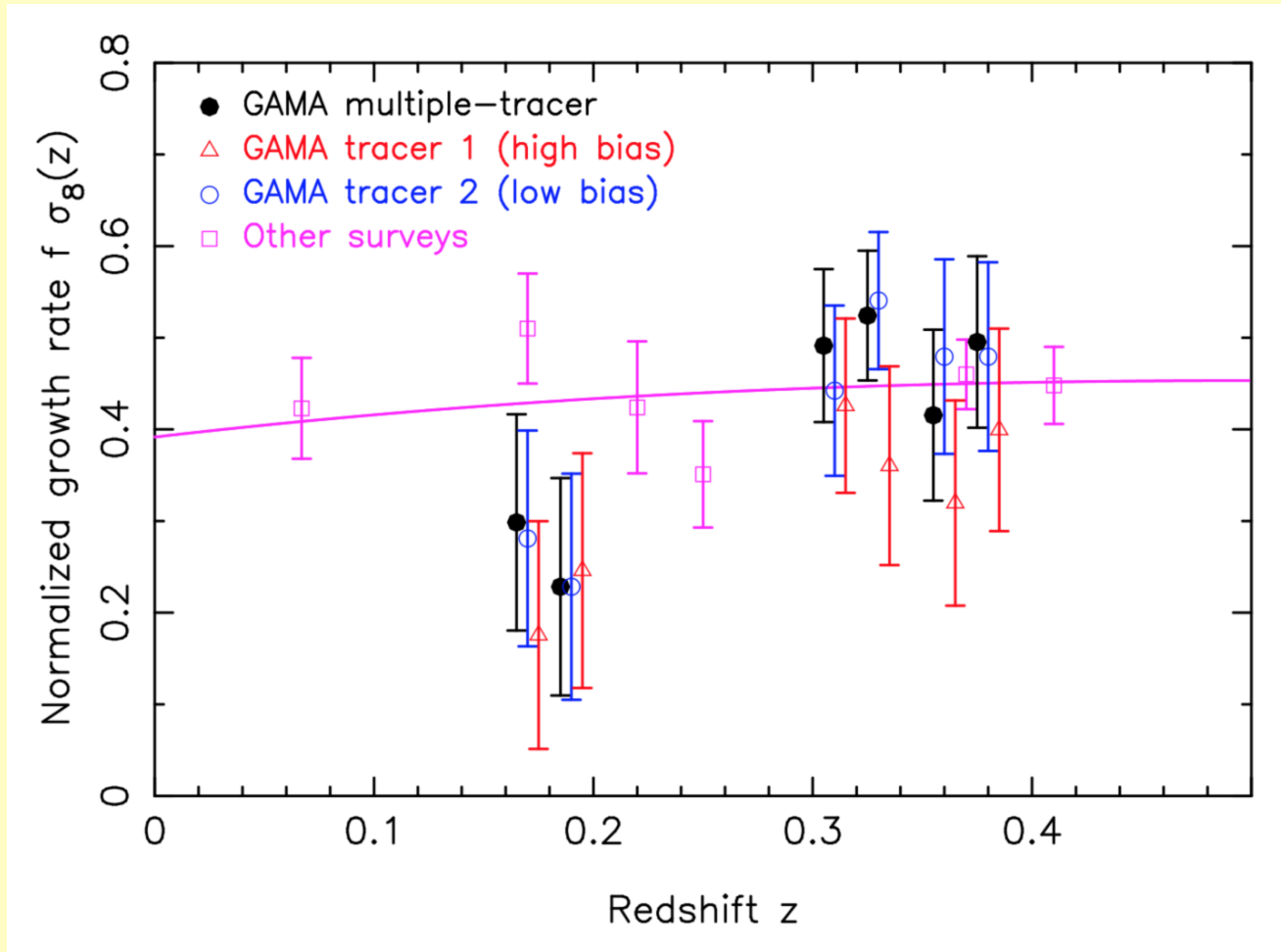
Blue

GAMA multi-tracer analysis

Blake et al. (2013): $\sim 10\%$ improvement on single-tracer f_g error

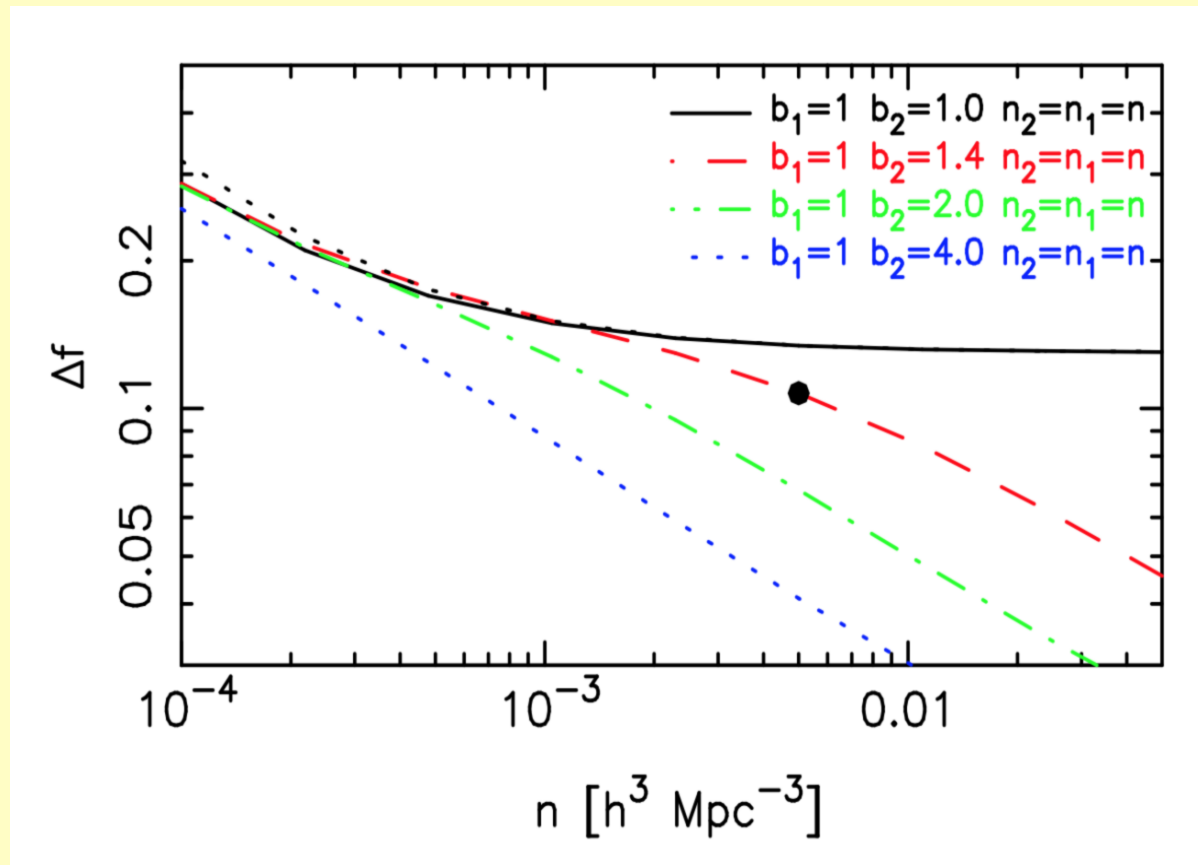


Growth-rate results



cf. BOSS DR9 0.415 ± 0.034 at $z=0.57$

Outlook



Gain requires (a) large bias difference; (b) low shot noise. Normally hard to reconcile – but Seljak et al. estimate shot noise (for one tracer) can be reduced $\times \sim 10$: should allow factor 2 improvement in multi-tracer result.

Summary

- Different tracers have complex relation to mass field
 - But most systematics can be understood via halo model
 - Discreteness noise can be tamed by halo weighting
- Cosmic variance is not the ultimate limit to precision
 - Practical implementation of McDonald-Seljak method
 - Proof of concept applied to GAMA
 - Scope for factor > 2 improvement in BOSS results if suitable subsets can be identified (harder with sparser samples)

