

Putting cosmology to the test with computer simulations

Carlos S. Frenk
Institute for Computational Cosmology,
Durham





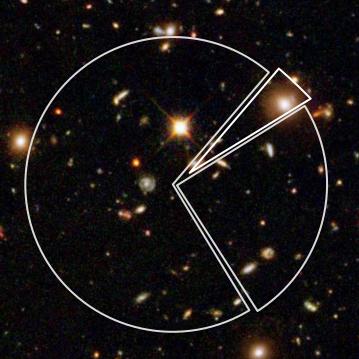
Flammarion 1888: tete des etoiles



What is the Universe made of?

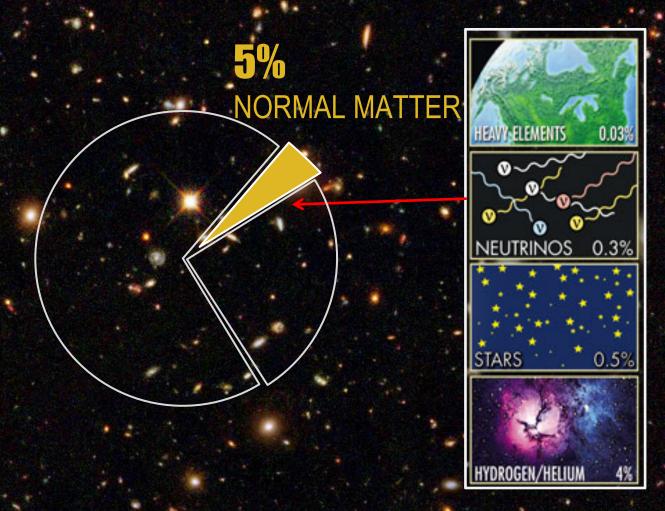
The (bizarre) contents of our Universe





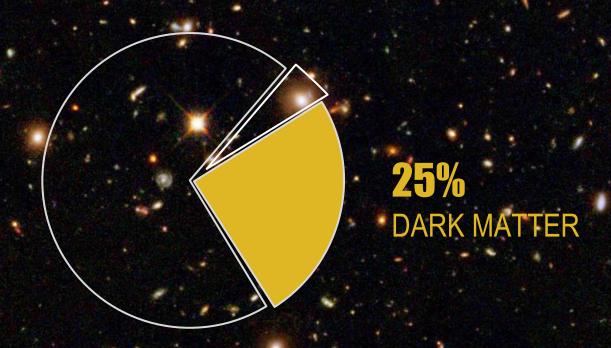
Normal matter = matter made of ordinary atoms





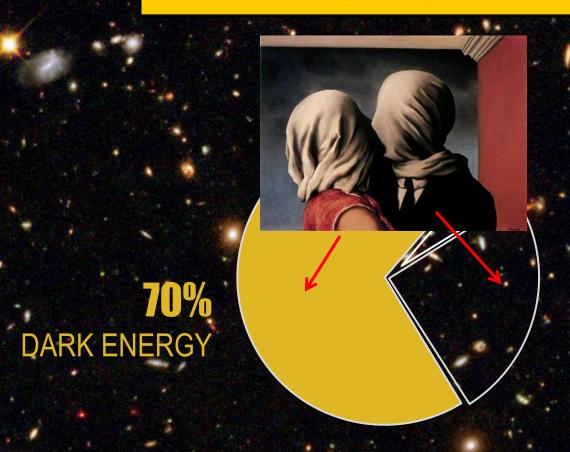
Normal matter = matter made of ordinary atoms





Dark matter = matter that does not emit light at any wavelength

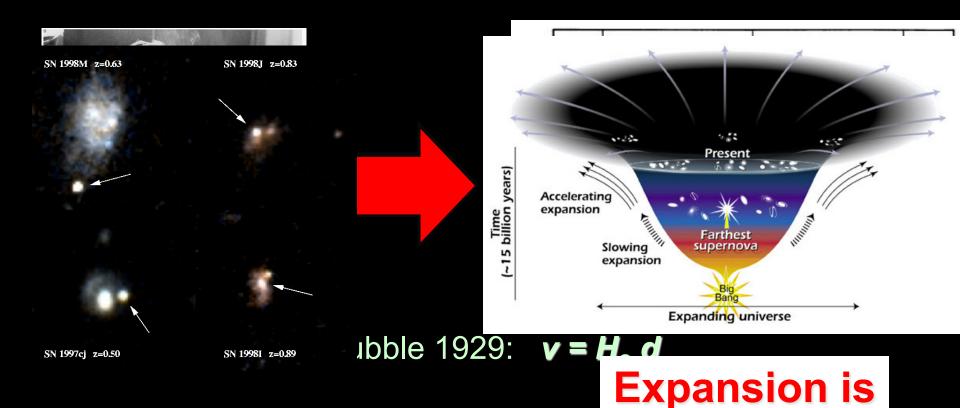




Dark energy = mysterious form of energy which opposes gravity and is causing the cosmic expansion to accelerate

The cosmic expansion

1998



2011 Nobel prize in physical energy

Computational Cosmolog

rating



What is the cosmic dark energy?

A form of energy that produces a repulsive force, causing the universal expansion to accelerate

It is likely to be energy associated with empty space – the vacuum

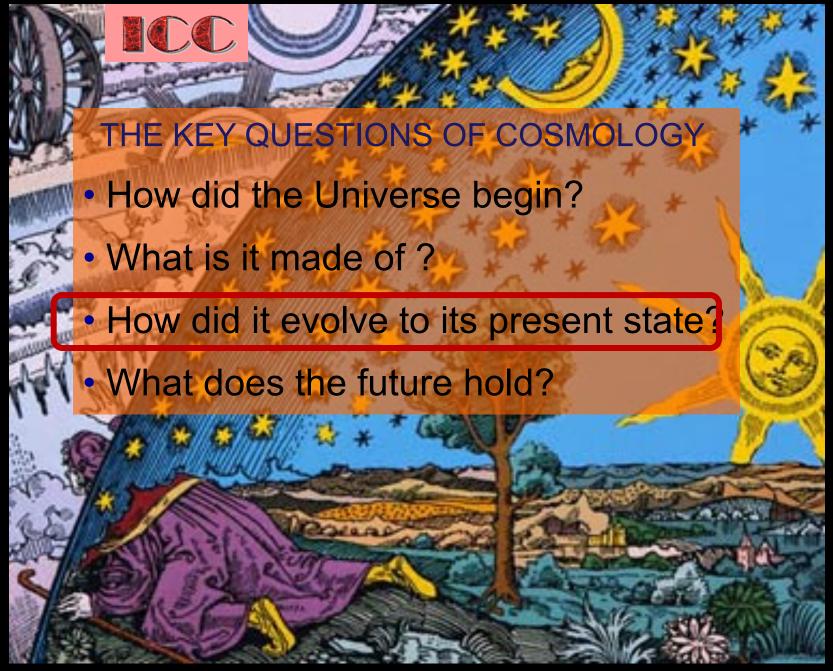


The gravitational constant



The simplest form of vacuum energy is the gravitational constant ∆ introduced by Einstein (for the wrong reasons)

There is no physical explanation for the measured value of Λ



Flammarion 1888: tete des etoiles



The standard model of cosmology



The ACDM model of cosmogony



- Proposed in 1980s, it is an ab initio, fully specified model of cosmic evolution and the formation of cosmic structure
- Has strong predictive power and can, in principle, be ruled out
- Has made a number of predictions that were subsequently verified empirically (e.g. CMB, LSS, galaxy formation)

Three Nobel Prizes in Physics since 2006, including 2019



Non-baryonic dark matter candidates

From the early 1980s:

Type	example	mass

hot	neutrino	few tens of eV
warm	sterile v	keV-MeV
cold	axion neutralino	10 ⁻⁵ eV - 100 GeV

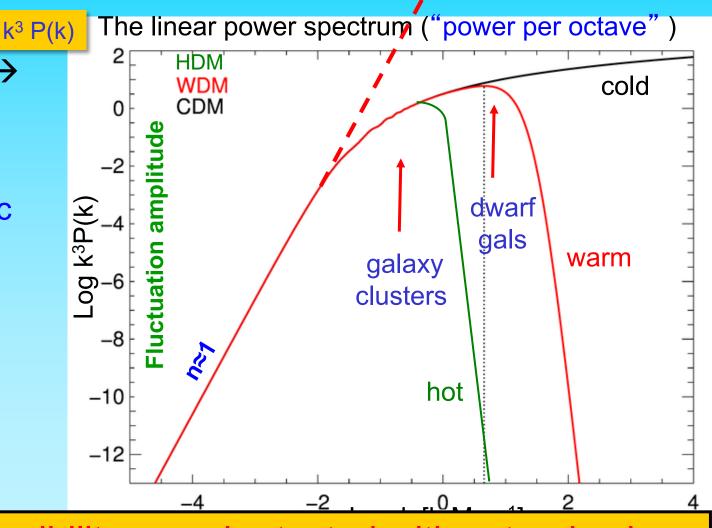


The dark matter power spectrum

Free streaming →

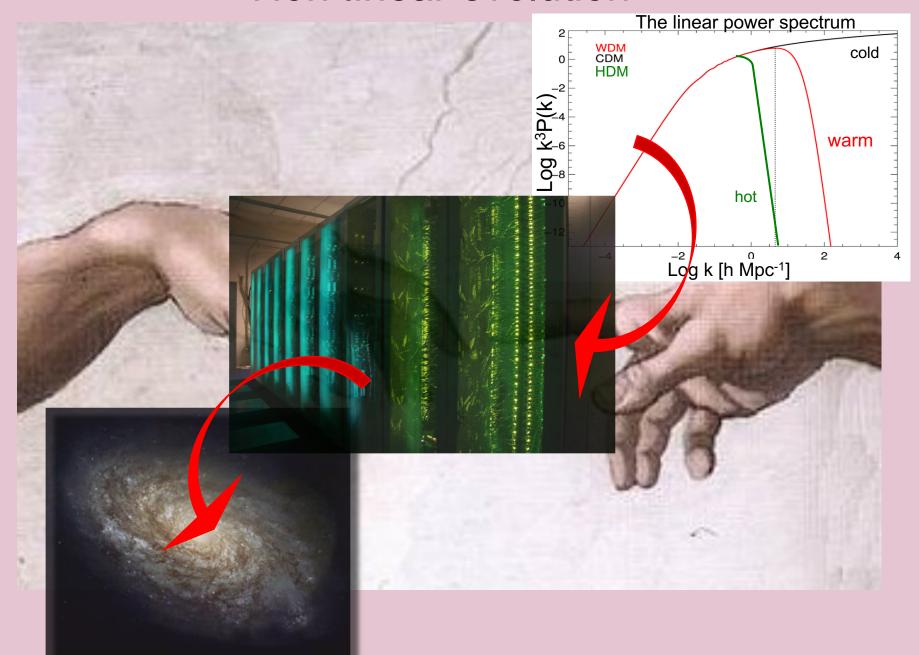
 $\lambda_{cut} \alpha m_x^{-1}$

for a thermal relic



These possibilites can be tested with astrophysics

Non-linear evolution





Non-linear evolution: simulations

Assumption about content of Universe -> Initial conditions

Relevant equations:

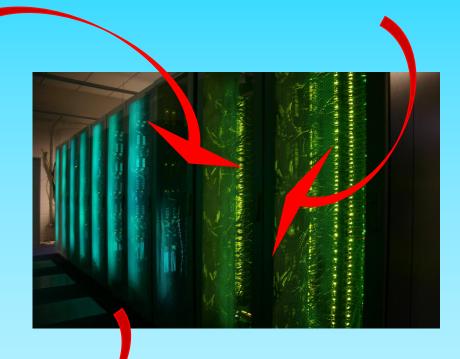
Collisionless Boltzmann;

Poisson; Friedmann eqns;

Radiative hydrodynamics

Subgrid astrophysics





How to make a virtual universe

LUBIMOV -7-

$m_v = 30 \text{ ev} \rightarrow \Omega_m = 1$

HAS THE NEUTRINO A NON-ZERO REST MASS? (Tritium β-Spectrum Measurement)

V. Lubimov, E. Novikov, V. Nozik, E. Tretyakov Institute for Theoretical and Experimental Physics, Moscow, U.S.S.R.

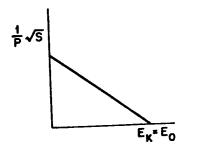
> V. Kosik Institute of Molecular Genetics, Moscow, U.S.S.R.

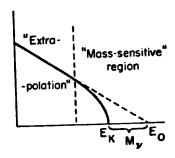
ABSTRACT

The high energy part of the β-spectrum of tritium in the molecule was measured with high precision by a toroidal β-spect meter. The results give evidence for a non-zero electron antineutrino mass.

Fifty years ago Pauli introduced the neutrino to explain the 3-spectrum shape. Pauli made the first estimate of the neutrino mass (E $_{3~max}$ muclei mass defect): it should be very small or maybe zero. Up to now the study of the β -spectrum shape is the

most sensitive, direct method of neutrino mass measurement. For allowed β -transitions, if M_{γ} = 0, then $S \simeq (E-E_{0})^{2}$. The Kurie plot is then a straight line with the only kinematic parameter being $E_k = E_0$ (total β -transition energy). If $M_0 \neq 0$, then $S = (E_o - E) \sqrt{(E_o - E)^2 - M_V^2}$. The Kurie plot is then distorted, especially near the endpoint.





1981

Fig. 2. Kurie plot for $M_{ij} \neq 0$. Fig. 1. Kurie plot for $M_{ij} = 0$.

The method for the neutrino mass measurement is to obtain E from the extrapolation and obtain Ek from the spectrum intercept. Then $\frac{1}{4}$ = $E_0 - E_k$. Qualitatively, $\frac{1}{4}$ = 0 if the 3-spectrum near the endpoint runs below the extrapolated curve.

things are more complicated. The apparatus resocongly affects the spectrum endpoint and rather e spectrum slope.

M, = 0 Background ealistic Kurie plot.

extrapolation. However, we are unable then once again the lack of counts near the indicate that $M_{\downarrow} = 0$. If $M_{\downarrow} \leq R$, the changes due to mass and the influence of R are indistinguishable. For M ermination the knowledge of R is compulsory. The background $exttt{de}^{ee}$ termines the statistical accuracy near the endpoint, i.e., in the region of the highest sensitivity to the ν mass. So: 1) R should be \sim M_J, 2) the smaller M_J is, the smaller the background (\sim M_Z) must be and the higher the statistics (\sim M_J⁻³) must be. For example, suppose that for M, = 100 eV we need resolution R, background Q, and statistics N. If M_{γ} = 30 eV, to achieve the same $\Delta M/M$ they should be R/3, Q/10, and N \times 30, respectively.

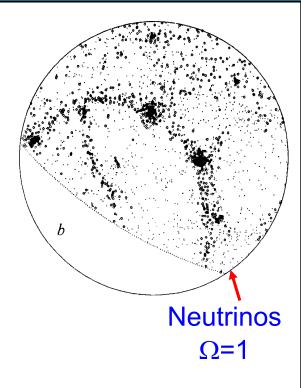
The shorter the $\beta\mbox{-spectrum,}$ the less it is spread due to R (as R $\sim \Delta p/p$ = const.). A classical example is ³H β-decay, which has 1) the smallest $E_{\mbox{\scriptsize 0}} \sim 18.6$ keV, 2) an allowed $\beta\text{-transition},$ simple nucleus, and simple theoretical interpretation, 3) highly reduced radioactivity. The first experiments with ${}^{3}\mathrm{H}^{}$ were by S. Curran et al. (1948) and G. Hanna, B. Pontecorvo (1949). Using $^3\mathrm{H}$ gas in a proportional counter, they obtained $M_{i,j} \leq 1$ keV. Further progress required magnetic spectrometer development. This allowed the resolution to be improved considerably, and L. Langer and R. Moffat (1952) obtained $M_{\odot} \le 250$ eV. The best value was obtained by K. Bergkvist (1972): R \sim 50 eV and M, \leq 55 eV.

The ITEP spectrometer is of a new type: ironless, with toroidal magnetic field (E. Tretyakov, 1973). The principle of the toroidal magnetic field focusing systems was proposed by V. Vladimirsky et al. (An example is a "Horn" of v-beams.) It turns out that a rectilinear conductor (current) has a focusing ability for particles emitted perpendicular to the rotation axis. This system has infinite periodical focusing structure. The ITEP spectrometer is based on this principle.

Paper presented by Oleg Egorov.



Non-baryonic dark matter cosmologies



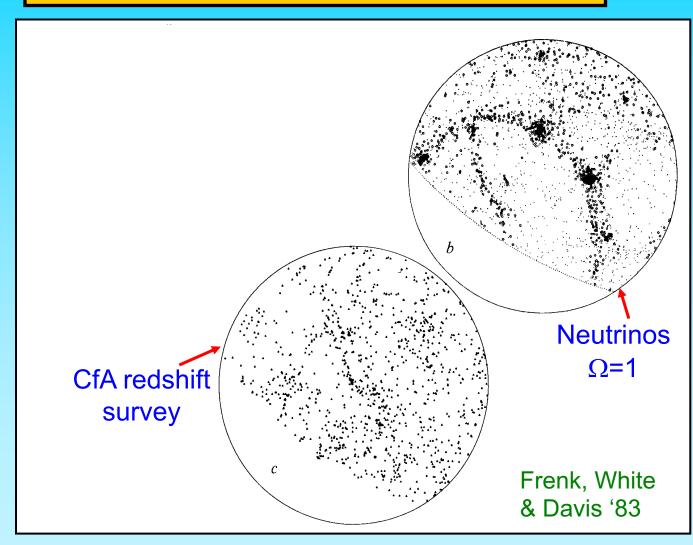
Frenk, White & Davis '83



Neutrino DM → wrong clustering

Neutrinos cannot make appreciable contribution to Ω \rightarrow m_v<< 30 ev

Non-baryonic dark matter cosmologies





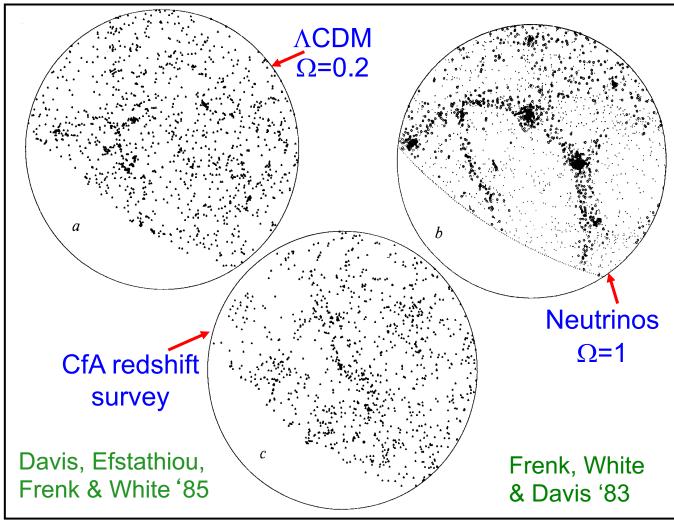
Neutrino DM → wrong clustering

Neutrinos cannot make appreciable contribution to Ω \rightarrow m, << 30 ev

Early CDM N-body simulations gave promising results

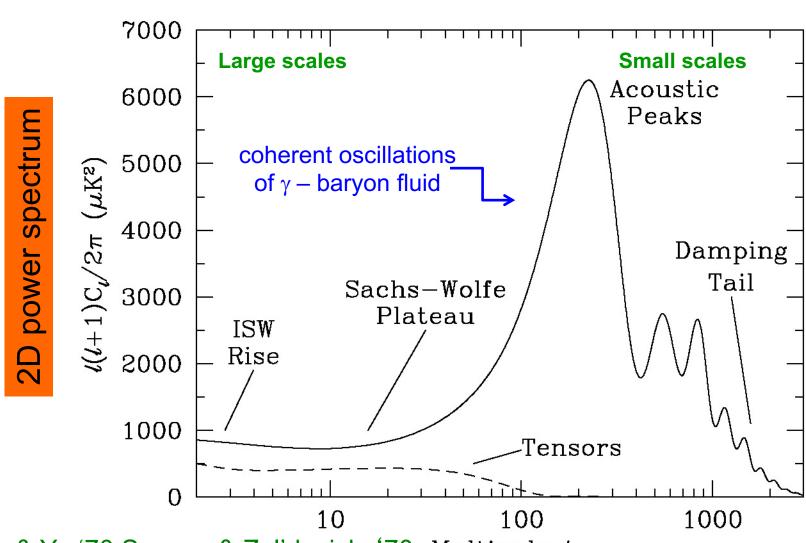
In CDM structure [forms hierarchically

Non-baryonic dark matter cosmologies





Temperature anisotropies in CMB



Peebles & Yu '70 Sunyev & Zel'dovich '70 Multipole l

For CDM: Peebles '82; Bond & Efstathiou '84



Jim Peebles
Nobel prize 2019

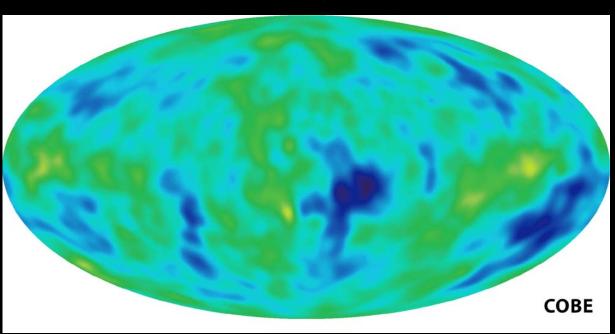




The CMB

1992





George Smoot - Nobel Prize 2006



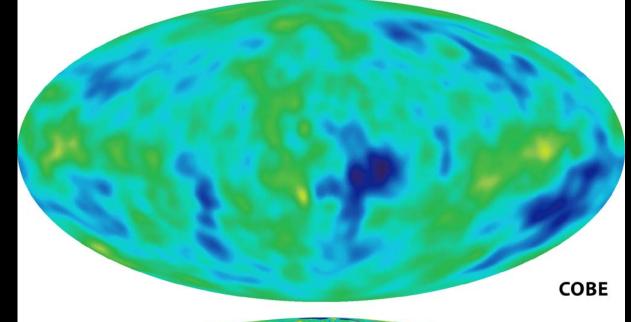




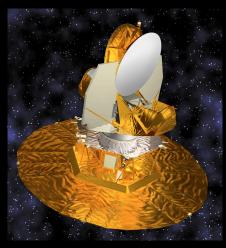
The CMB

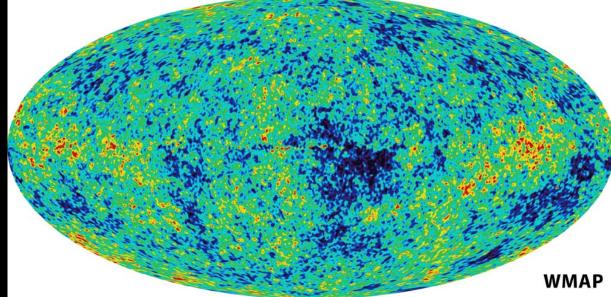
1992





2003

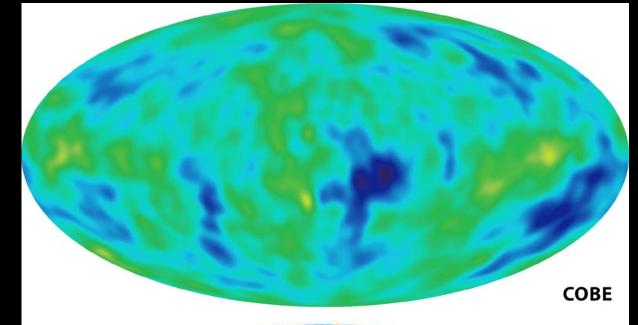


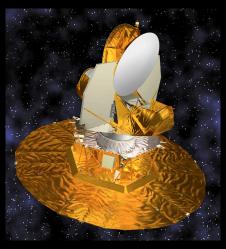


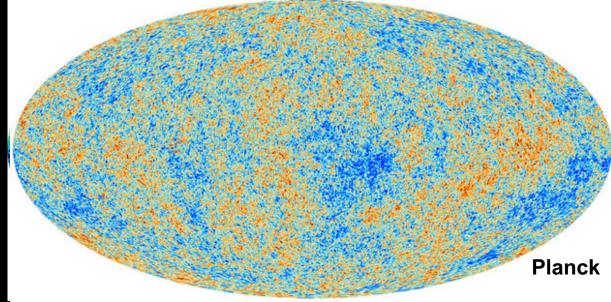


The CMB



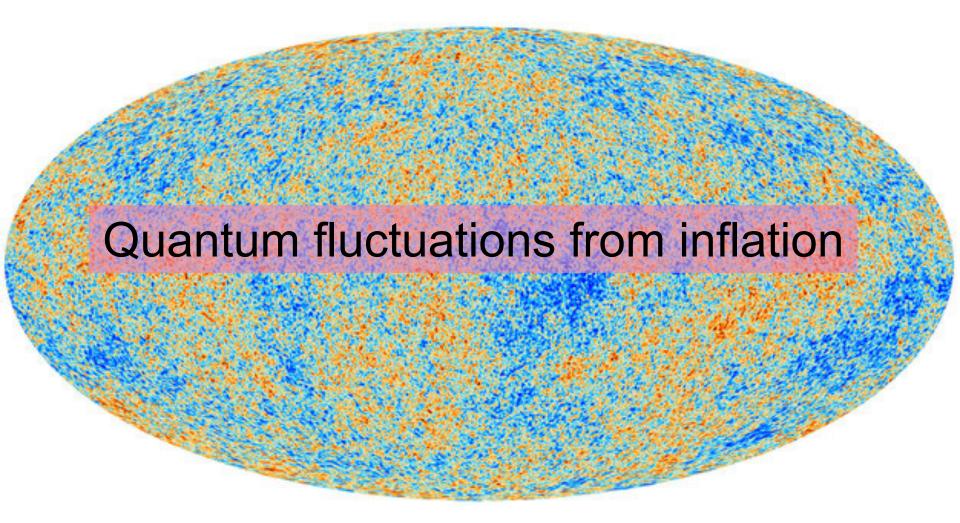






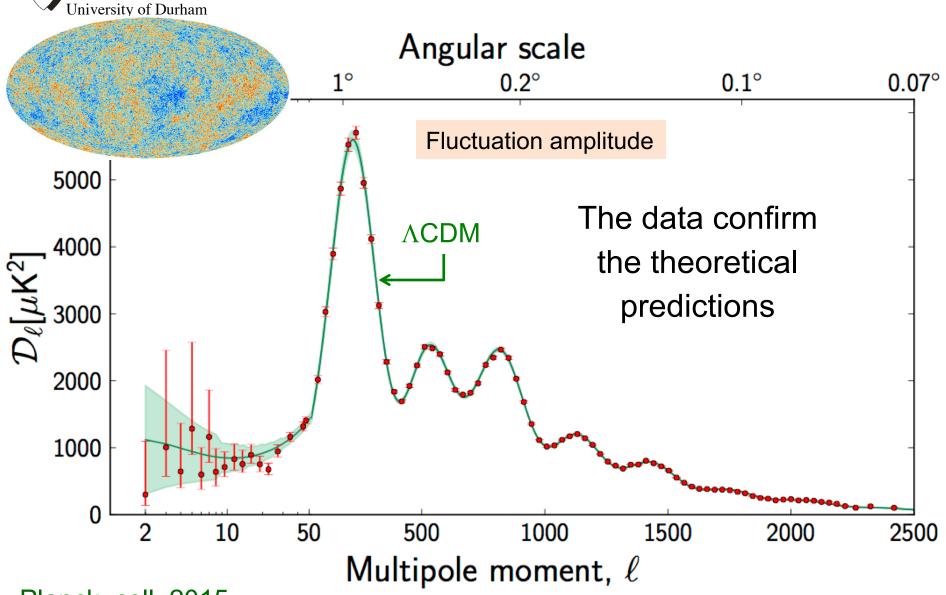


The initial conditions for galaxy formation





Planck: CMB temperature anisotropies



Planck coll. 2015





The six parameters of minimal \(\Lambda CDM \) model

	Planck+WP	
Parameter	Best fit	68% limits
Parameter $\Omega_{\rm b}h^2 \ . {\rm density} \ {\rm of} \ {\rm baryons} \ .$ $\Omega_{\rm c}h^2 \ . {\rm density} \ {\rm of} \ {\rm CDM} \ .$ $100\theta_{\rm MC} \ . \ . \ .$ $100\theta_{\rm MC} \ . \ . \ .$	0.022032	0.02205 atter 00028
$\Omega_{\mathrm{c}} h^2$, density of CDM	0.12038nic	0.1199 ± 0.0027
$100\theta_{\mathrm{MC}}$ of	non-04119	1.04131 ± 0.00063
τ. A400 detect.	0.0925	$0.089^{+0.012}_{-0.014}$
$n_{\rm S}$	0.9619	0.9603 ± 0.0073
$\ln(10^{10}A_{\rm s})\ldots\ldots$	3.0980	$3.089^{+0.024}_{-0.027}$

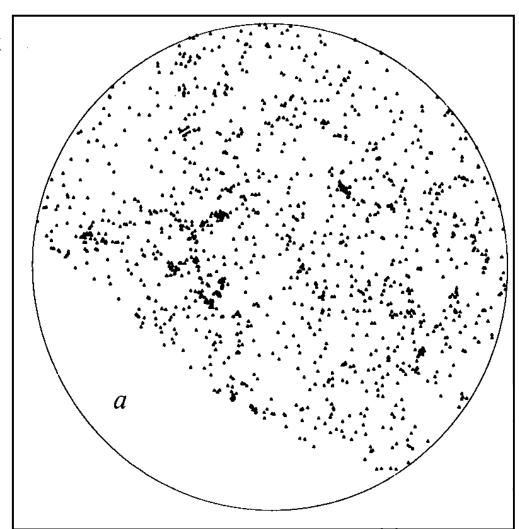
 $\Omega = 1$

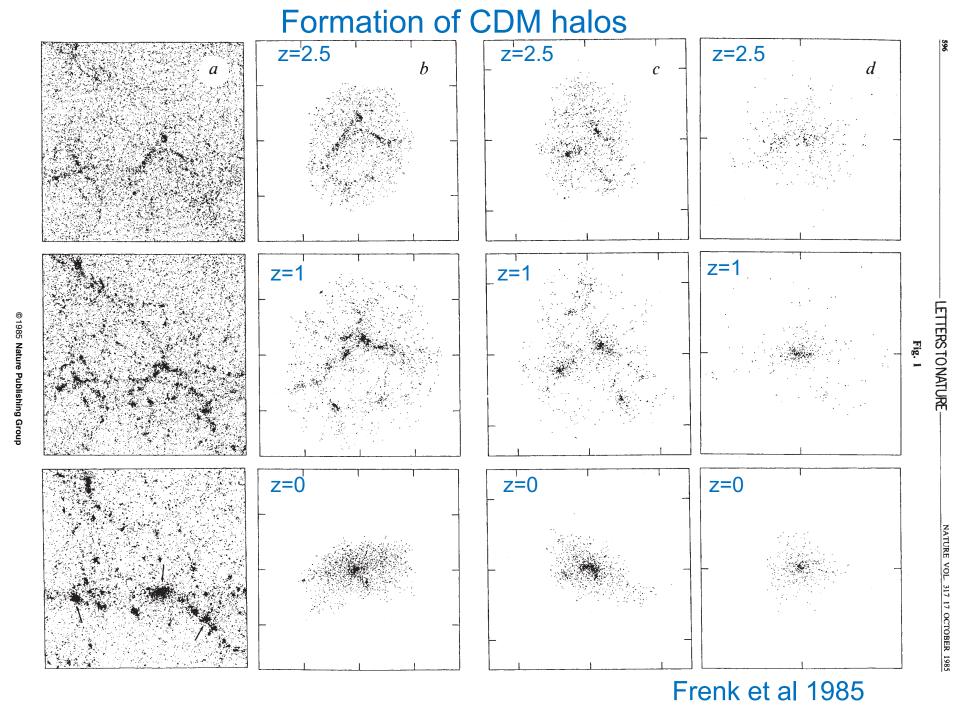
Planck collaboration '13



N-body simulations of largescale structure in ΛCDM

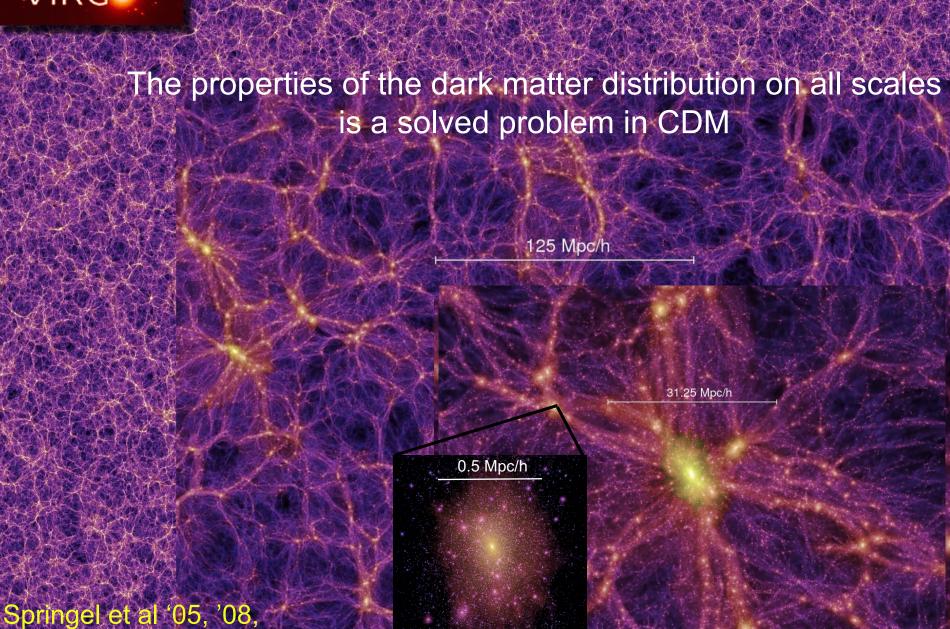
Davis, Efstathiou, Frenk & White '85





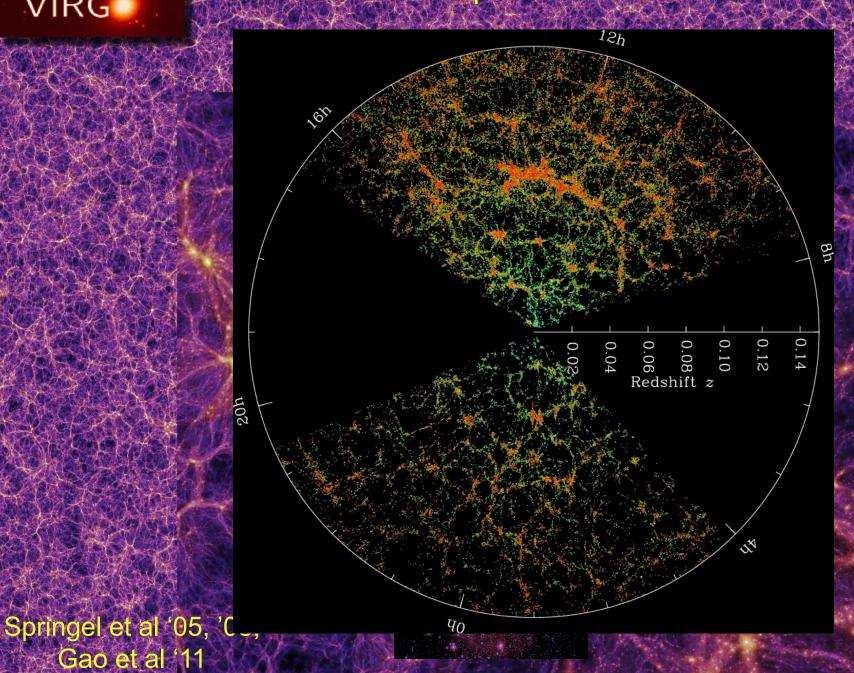
Gao et al '11

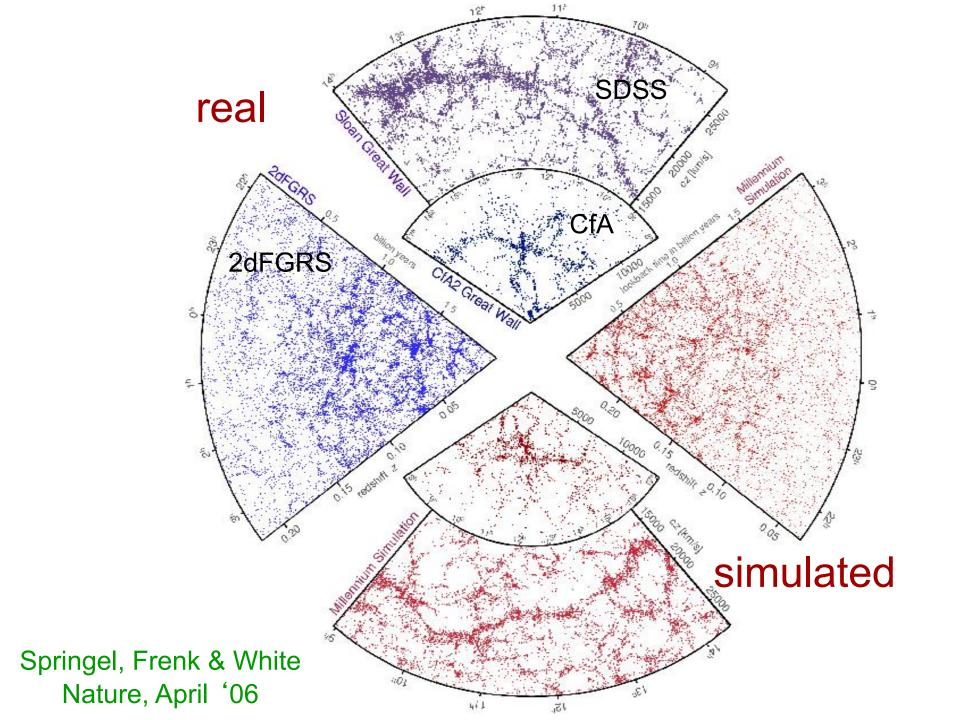
The Millennium/Aquarius/Phoenix simulation series





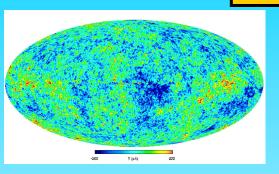
The Millennium/Aquarius/Phoenix simulation series





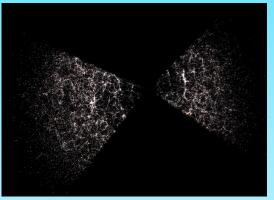


The cosmic power spectrum: from the CMB to the 2dFGRS



z~1000

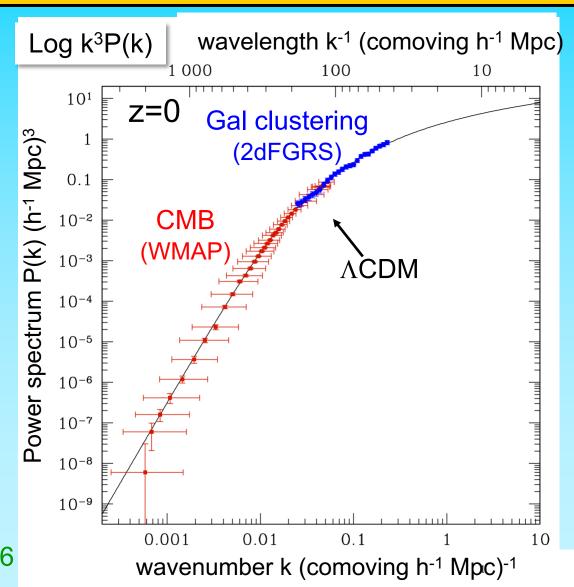
z~0



Z~U

→ ΛCDM provides an excellent description of mass power spectrum from 10-1000 Mpc

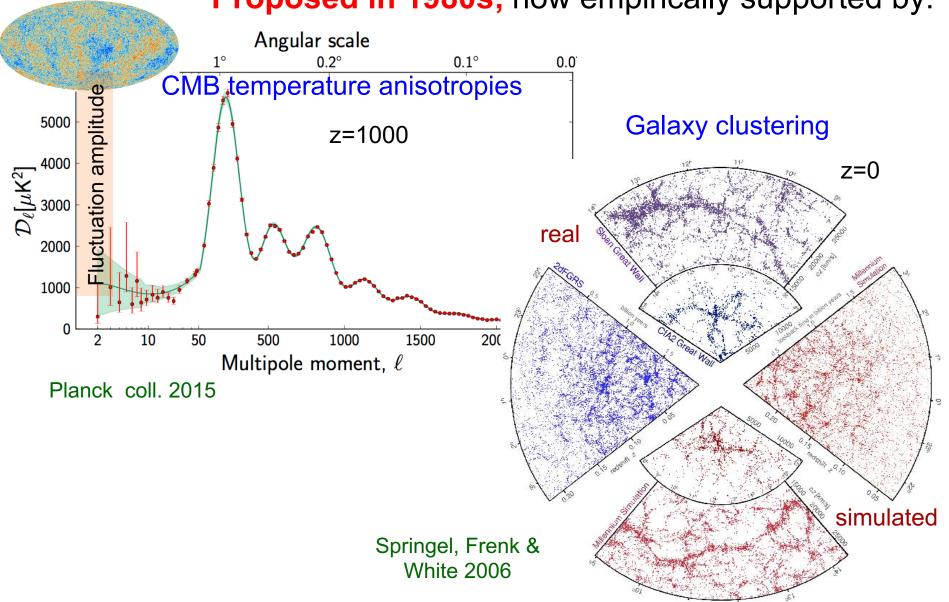
Sanchez et al 06





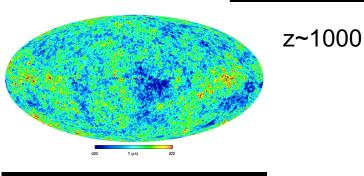
The ACDM model of cosmogony

Proposed in 1980s; now empirically supported by:





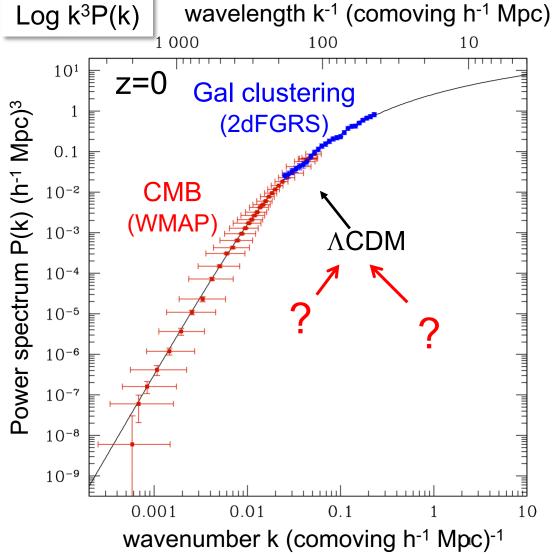
The cosmic power spectrum: from the CMB to the 2dFGRS





 \Rightarrow \land CDM provides an excellent description of mass power spectrum from 10-1000 Mpc

Sanchez et al 06





The cosmic power spectrum: from the CMB to the 2dFGRS

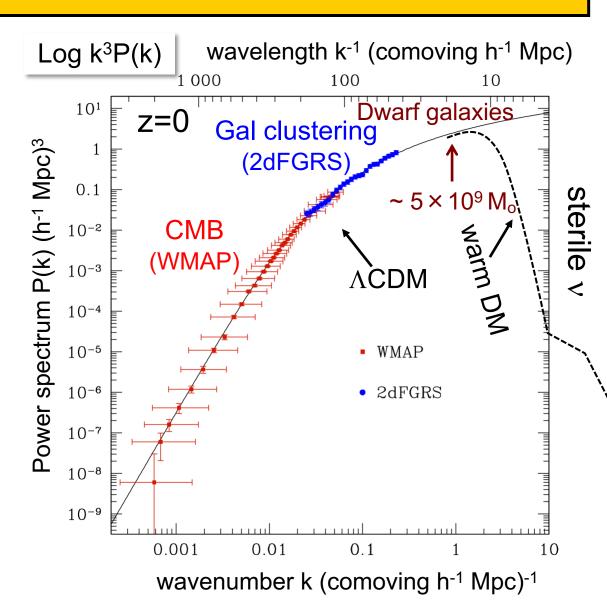
Free streaming →

 $\lambda_{cut} \alpha m_x^{-1}$

for thermal relic

 $m_{CDM} \sim 100 GeV$ susy; $M_{cut} \sim 10^{-6} M_o$

 $m_{WDM} \sim \text{few keV}$ sterile v; $M_{cut} \sim 10^9 M_o$





Both CDM & WDM compatible with CMB & galaxy clustering Claims that both types of DM have been discovered:

- ♦ CDM: γ-ray excess from Galactic Center
- ♦ WDM (sterile v): 3.5 X-ray keV line in galaxies and clusters



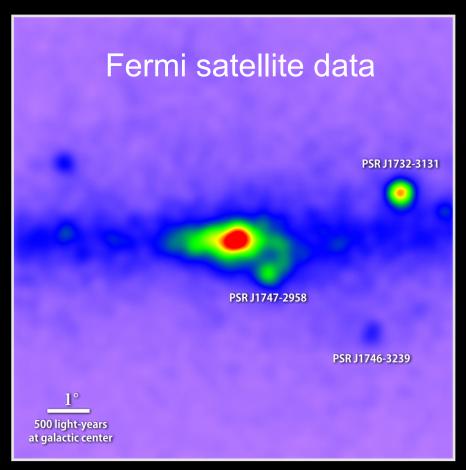
Cold dark matter In the Galactic Centre?

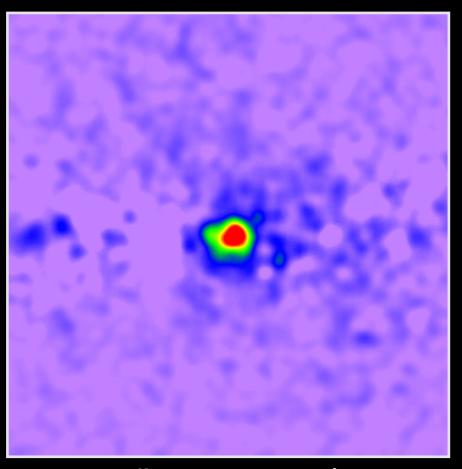
Cold dark matter

The Characterization of the Gamma-Ray Signal from the Central Milky Way:
A Compelling Case for Annihilating Dark Matter

Tansu Daylan,¹ Douglas P. Finkbeiner,^{1,2} Dan Hooper,^{3,4} Tim Linden,⁵ Stephen K. N. Portillo,² Nicholas L. Rodd,⁶ and Tracy R. Slatyer^{6,7}

Uncovering a gamma-ray excess at the galactic center





Unprocessed map of 1.0 to 3.16 GeV gamma rays

Known sources removed



Warm dark matter

Decay line at 3.51 keV in galaxies and clusters

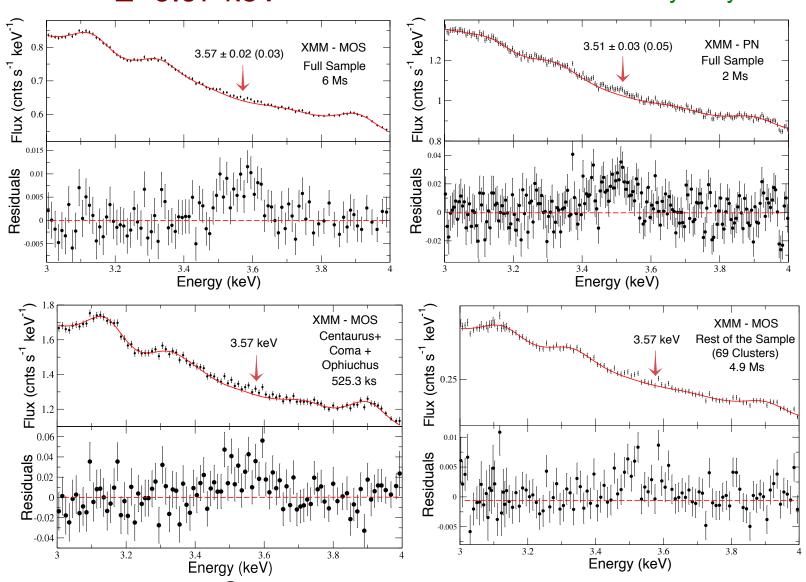
Institute for Computational Cosmology



Warm dark matter WDM decay line in 69 stacked clusters?

E=3.57 keV

Bulbul et al. '14 See also Boyarsky et al. '14





Both CDM & WDM compatible with CMB & galaxy clustering Claims that both types of DM have been discovered:

- ♦ CDM: γ-ray excess from Galactic Center
- ♦ WDM (sterile v): 3.5 X-ray keV line in galaxies and clusters

Very unlikely that both are right!



The identity of the dark matter is encoded in dwarf galaxies and in the halo of the MW

(strongly non-linear regime)



Cold Dark Matter

Warm Dark Matter

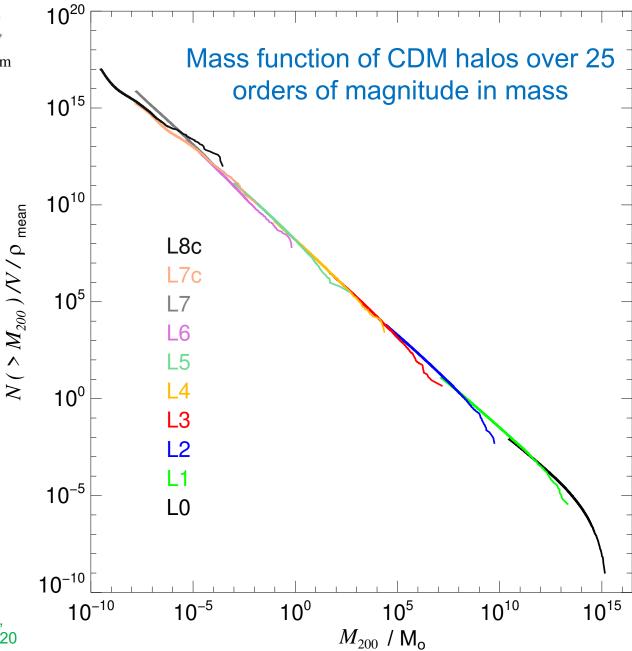
cold dark matter

warm dark matter



Lovell, Eke, Frenk, Gao, Jenkins, Wang, White, Theuns, Boyarski & Ruchayskiy '12



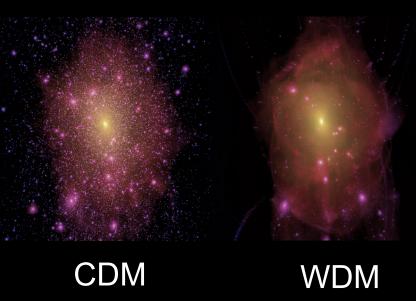


Wang, Bose, CSF, Gao, Jenkins, Springel, White '20

Institute for Computational Cosmology

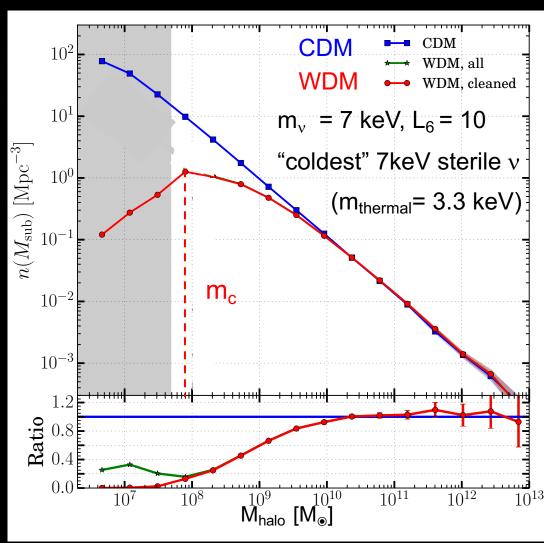


The subhalo mass function



3 x fewer WDM subhalos at $3x10^9\,M_o$

10 x fewer at 108 M_o





How can we distinguish the two?

Astrophysical tests of dark matter

Count the number of small-mass halos

- Number of dark matter halos (the halos mass fn.)
 (the ``missing satellites problem)
- 2. Annihilation/decay radiation

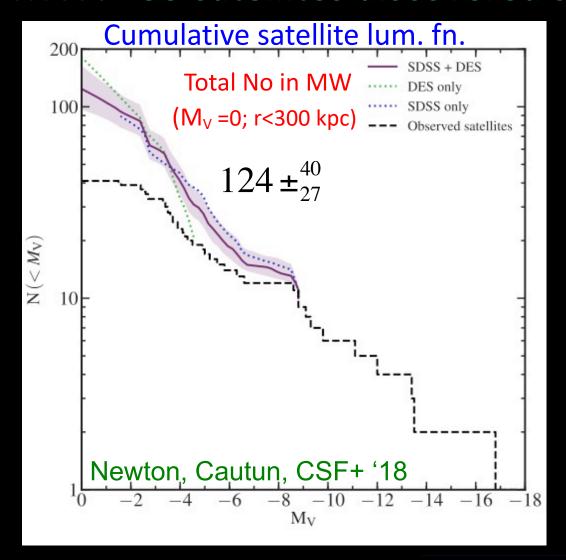
Let's begin by counting what we can see

Institute for Computational Cosmology



The satellites of the Milky Way

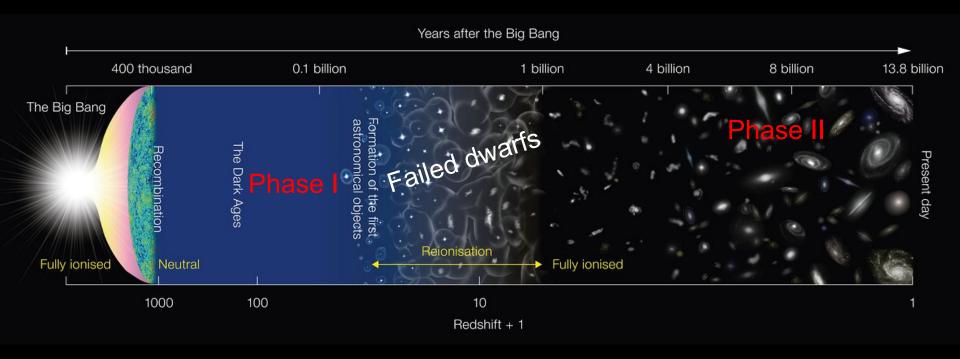
In the MW: ~55 satellites discovered so far







The two phases of galaxy formation



Phase I: During the "dark ages" H gas is neutral First stars reionize H and heat it up to 10⁴K

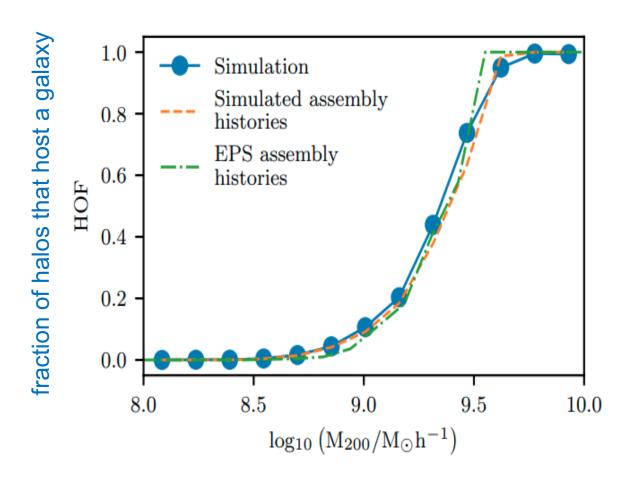
Phase II: H Gas is ionized ("T_{vir}" > 10⁴K form)

Institute for Computational Cosmology

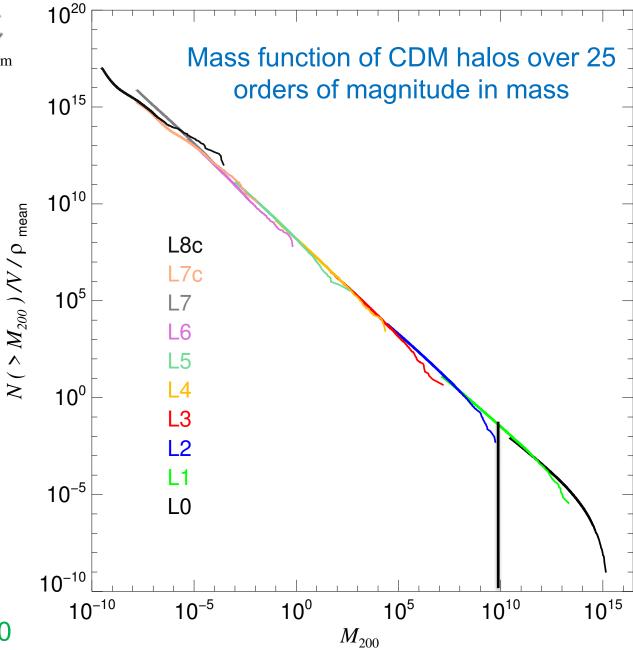


A galaxy formation primer

Halo Occupation Fraction (HOF): fraction of halos of a given mass that host a galaxy





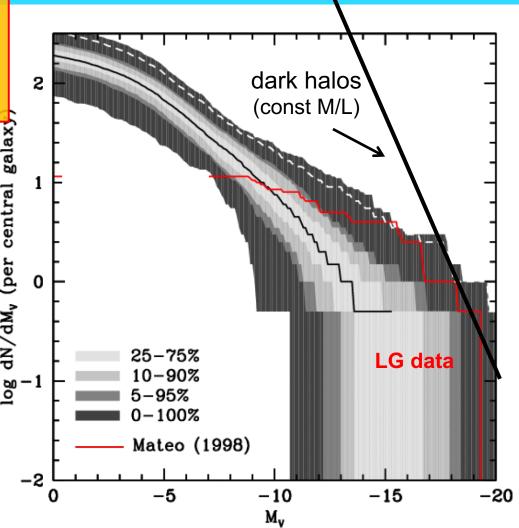




Luminosity Function of Local Group Satellites

Semi-analytic model of galaxy formation including effects of reionization and SN feedback

- Median model → correct abundance of sats brighter than M_V=-9 (V_{cir} > 12 km/s)
- Model predicts many, as yet undiscovered, faint satellites



Benson, Frenk, Lacey, Baugh & Cole '02 (see also Kauffman+ '93, Bullock+ '00, Somerville '02)

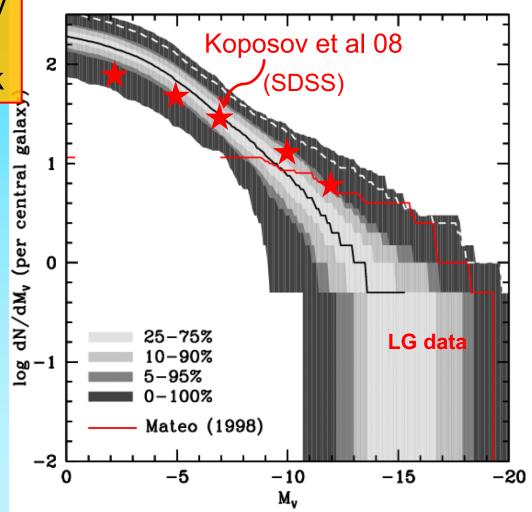
Institute for Computational Cosmology



Luminosity Function of Local Group Satellites

Semi-analytic model of galaxy formation ncluding effects of reionization and SN feedback

- Median model → correct abundance of sats brighter than M_V=-9 (V_{cir} > 12 km/s)
- Model predicts many, as yet undiscovered, faint satellites



Benson, Frenk, Lacey, Baugh & Cole '02 (see also Kauffman+ '93, Bullock+ '00, Somerville '02)

Institute for Computational Cosmology

"Evolution and assembly of galaxies and their environment"

THE EAGLE PROJECT

Virgo Consortium

Durham: Richard Bower, Michelle Furlong, Carlos Frenk, Matthieu Schaller, James

Trayford, Yelti Rosas-Guevara, Tom Theuns, Yan Qu, John Helly, Adrian Jenkins.

Leiden: Rob Crain, Joop Schaye.

Other: Claudio Dalla Vecchia, Ian McCarthy, Craig Booth...



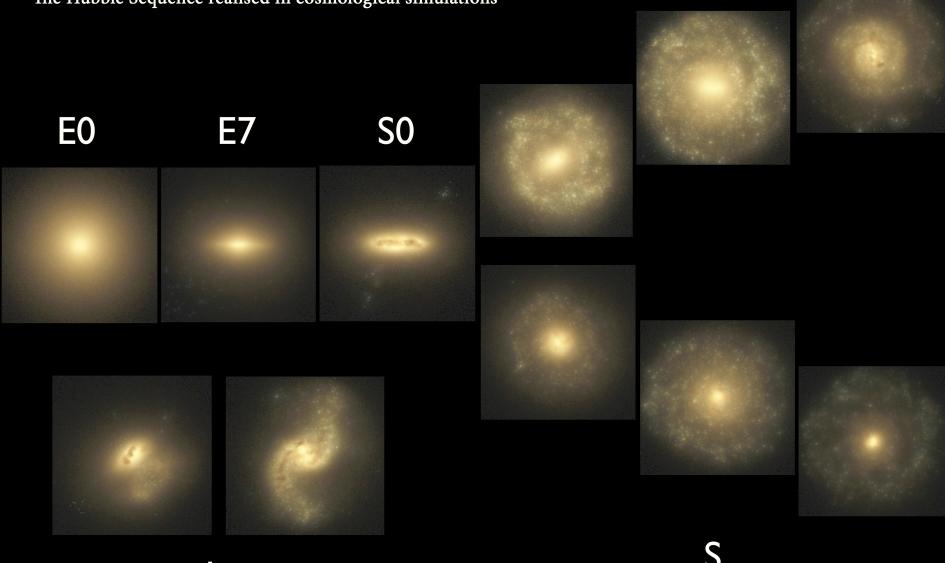




The Eagle Simulations

EVOLUTION AND ASSEMBLY OF GALAXIES AND THEIR ENVIRONMENTS

The Hubble Sequence realised in cosmological simulations



Irr

SB

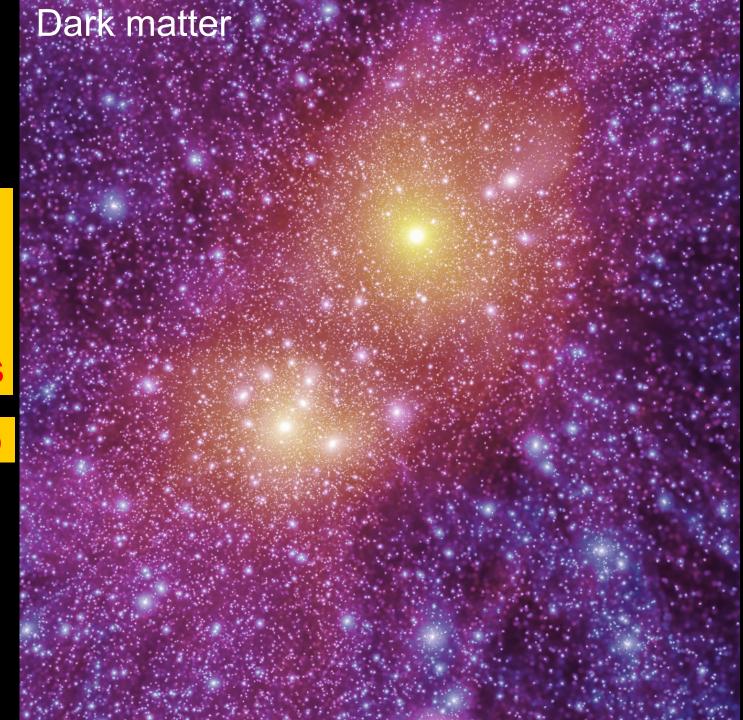
VIRG

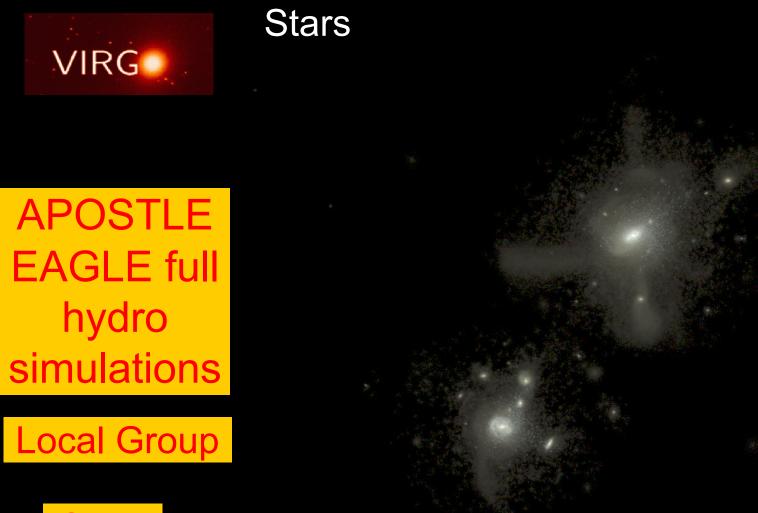
APOSTLE
EAGLE full
hydro
simulations

Local Group

CDM

Sawala, CSF et al '16





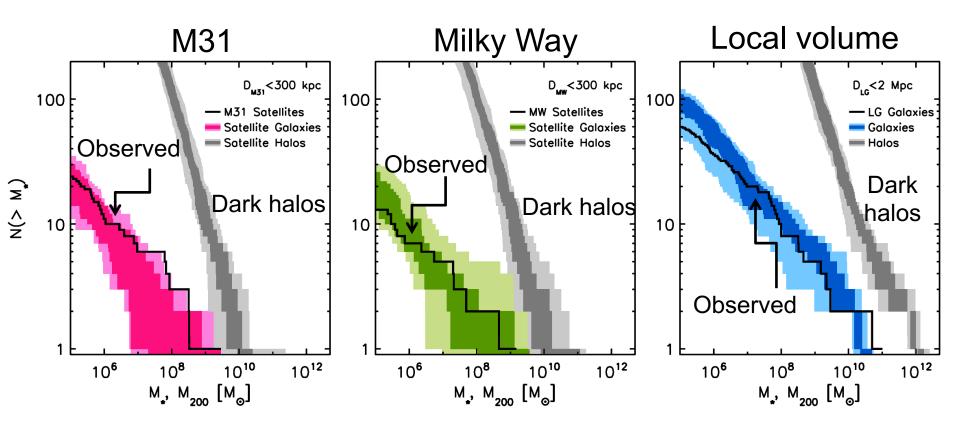
Stars

Far fewer satellite galaxies than CDM halos

Sawala, CSF et al '16



EAGLE Local Group simulation





When galaxy formation is taken into account

IIIIO GOOGUIII



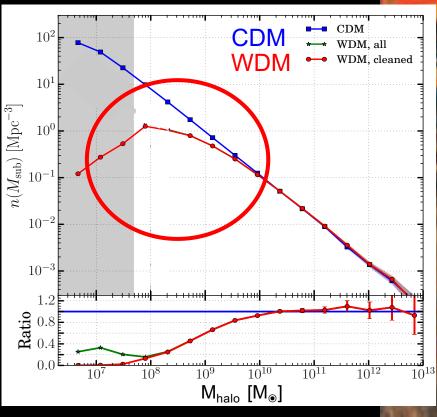
CDM predicts the observed abundance of satellites

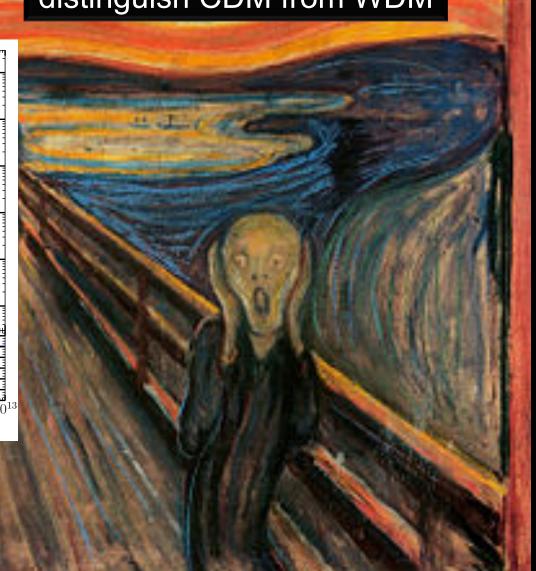


There is no such thing as a "missing satellite problem" in CDM!



But it doesn't help distinguish CDM from WDM







Can we distinguish CDM/WDM?

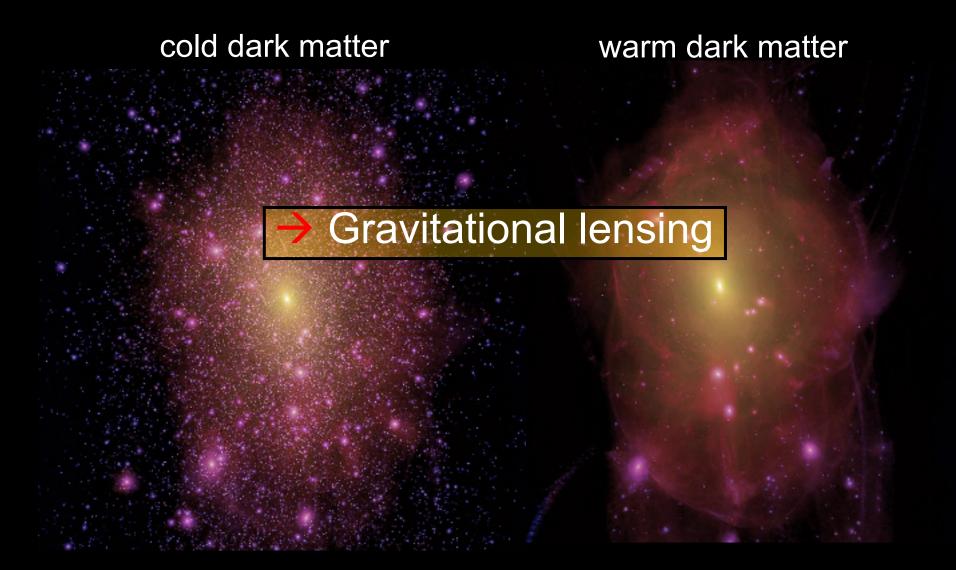
cold dark matter

warm dark matter

Rather than counting faint galaxies, count the number of dark halos ("failed dwarfs")

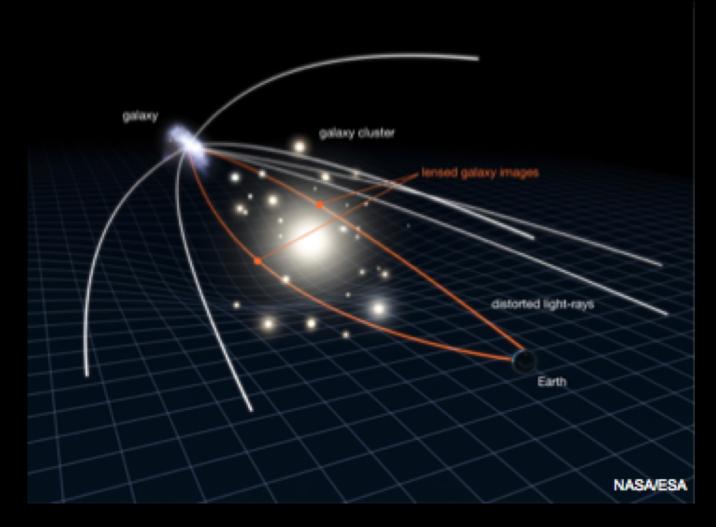


Can we count dark haloes?





Gravitational lensing: Einstein rings



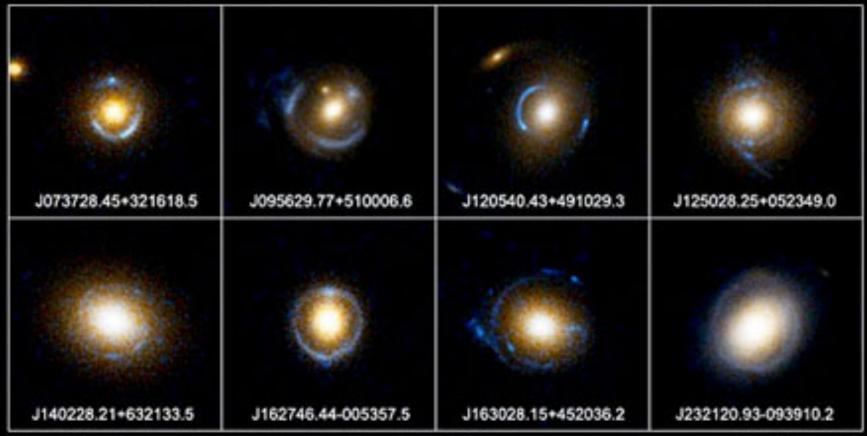
When the source and the lens are well aligned -> strong arc or an Einstein ring



SLAC sample of strong lenses

Einstein Ring Gravitational Lenses

Hubble Space Telescope . ACS

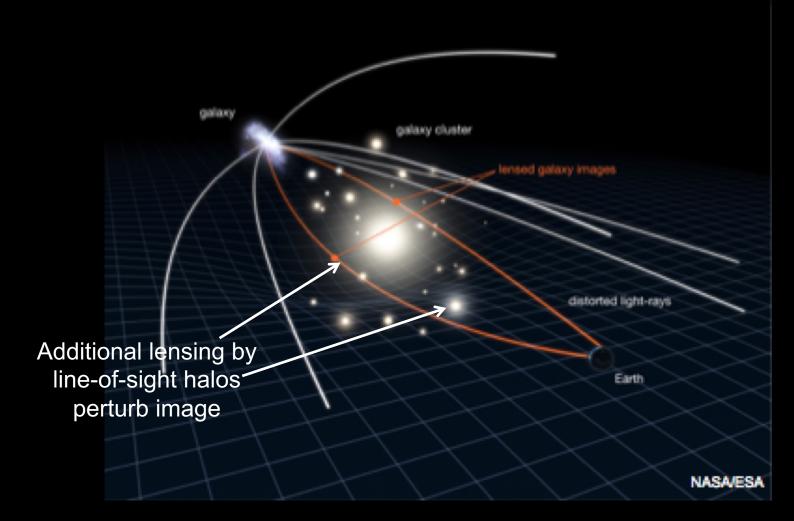


NASA, ESA, A. Bolton (Harvard-Smithsonian CfA), and the SLACS Team

STScI-PRC05-32



Gravitational lensing: Einstein rings

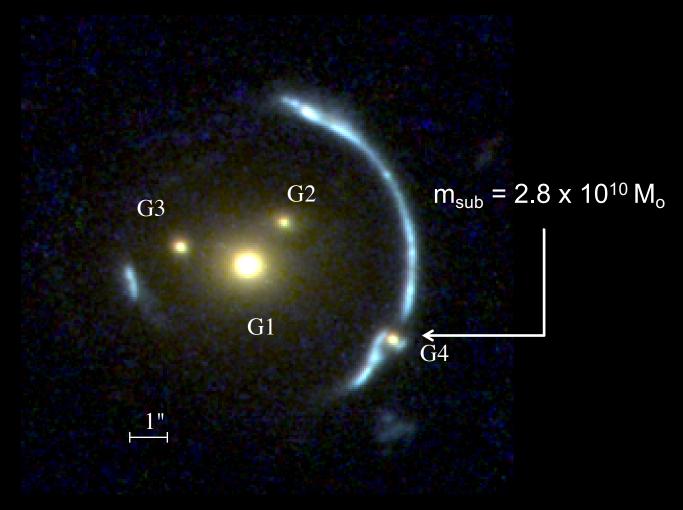


When the source and the lens are well aligned -> strong arc or an Einstein ring



Gravitational lensing: Einstein rings

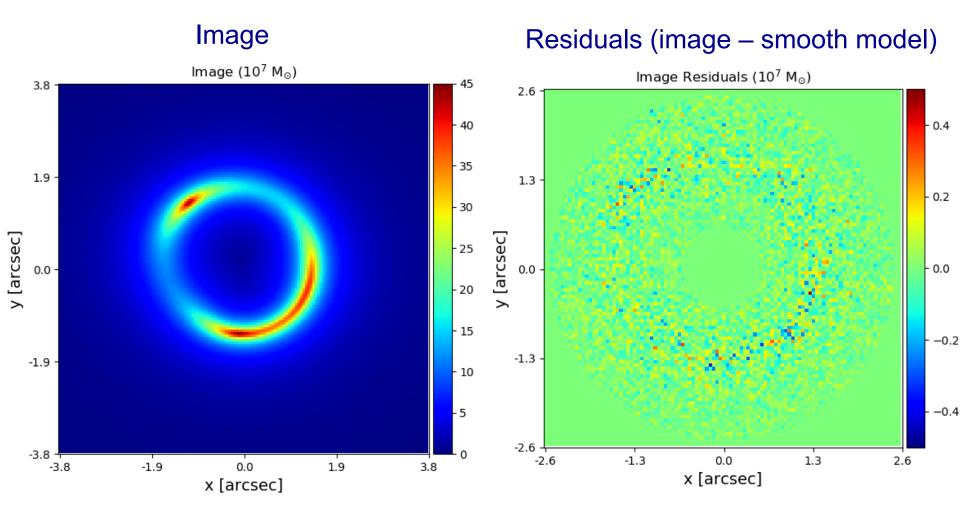
Halos projected onto an Einstein ring distort the image





Strong lensing: detecting small halos

HST "data": z_{source} =1; z_{lens} =0.2 10⁷ M_o halo – NOT so easy to spot

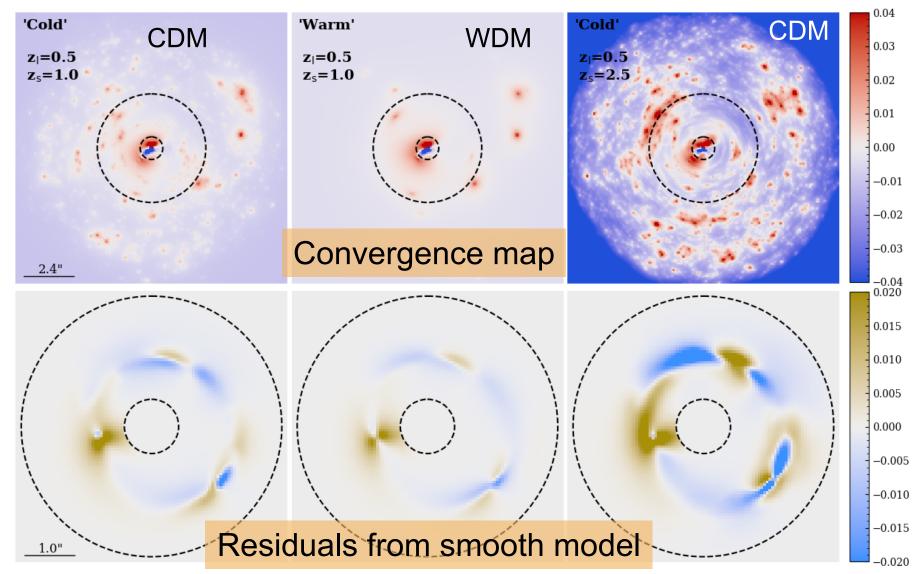


He, Li, CSF et al '19

Institute for Computational Cosmology



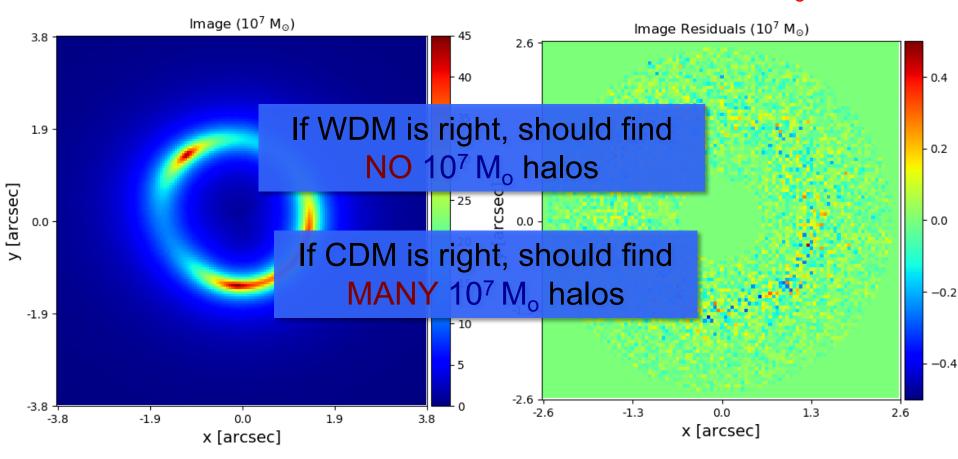
Strong lensing: detecting small halos





Detecting halos w. strong lensing

Can detect halos as small as $10^7 - 10^8 \,\mathrm{M}_{\odot}$

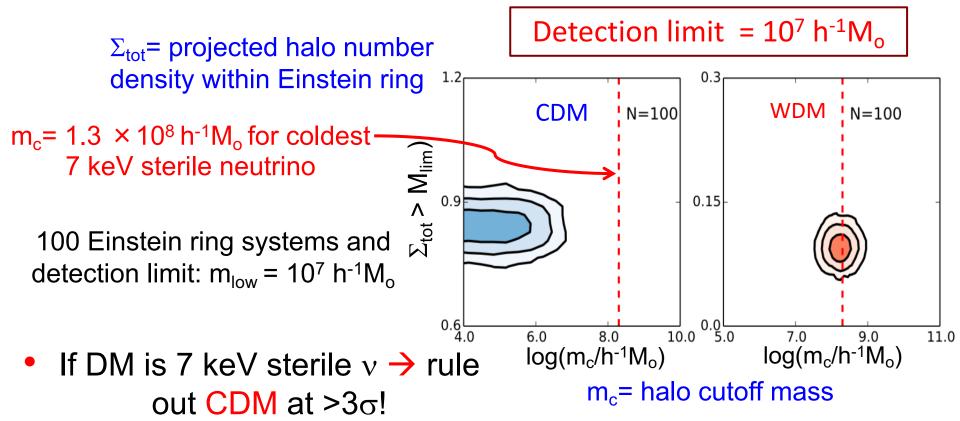


He, Li, CSF et al '19

Institute for Computational Cosmology



Detecting substructures with strong lensing



If DM is CDM → rule out 7 keV
 sterile v at many σ



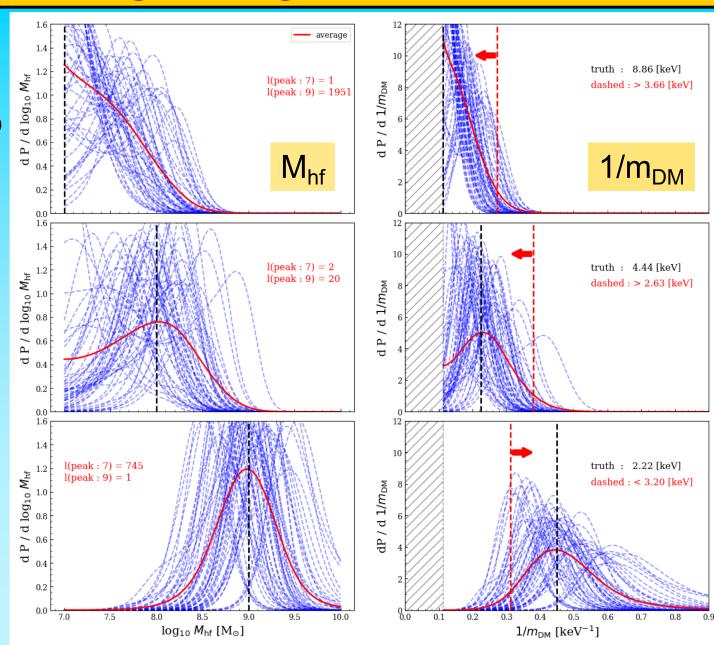
Strong lensing: statistical detection

Power spectrum of residuals map

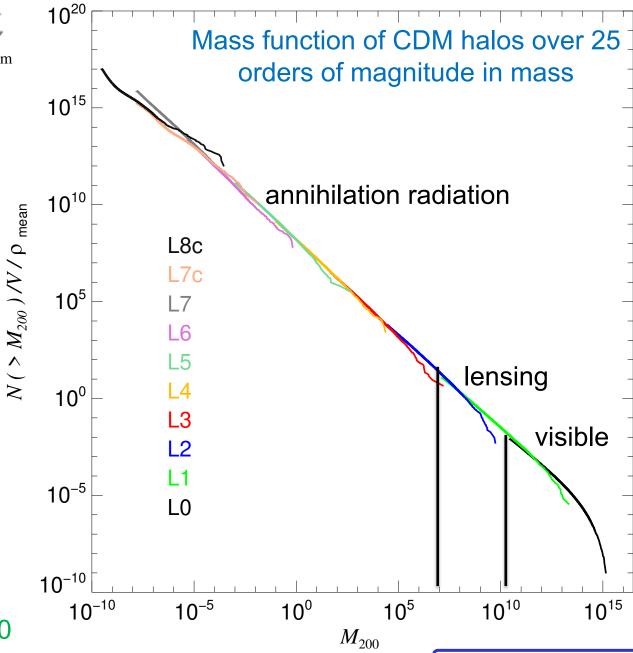
Posterior
distributions (mock
observations) for
power spectrum of
residuals

Constraints from forward modelling of 50 systems

He et al. '20









Indirect CDM detection through annihilation radiation

Supersymmetric particles are Majorana particles

annihilate into Standard Model particles (including γ-rays)

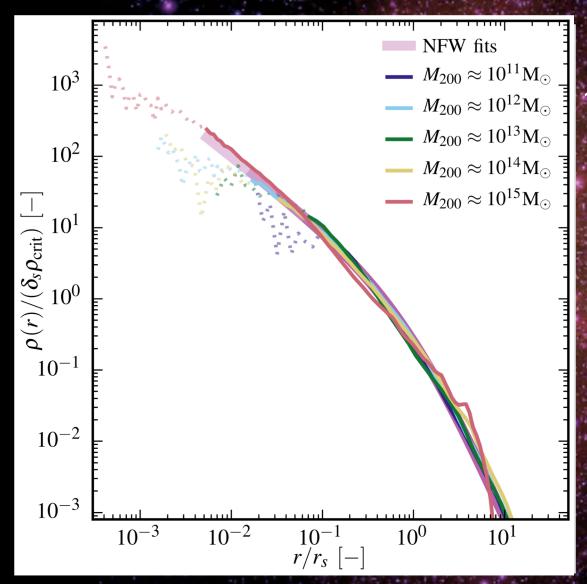
Intensity of annihilation radiation at x is:

$$I(x) = \frac{1}{8\pi} \sum_{f} \frac{dN_f}{dE} \langle \sigma_f v \rangle \int_{los} \left(\frac{\rho_{\chi}}{M_{\chi}} \right)^2 l dl$$
 cross-section (particle physics)

 $\langle \sigma v \rangle = 3 \times 10^{-26} cm^3 s^{-1}$ relic abundance in simple SUSY models

- \Rightarrow Theoretical expectation requires knowing $\rho(x)$
- → Accurate high resolution N-body simulations of halo formation from CDM initial conditions

The Density Profile of Cold Dark Matter Halos



Shape of halo profiles
~independent of halo mass &
cosmological parameters

Density profiles are "cuspy" - no `core' near the centre

Fitted by simple formula:

$$\frac{\rho(r)}{\rho_{crit}} = \frac{\delta_c}{(r/r_s)(1+r/r_s)^2}$$

(Navarro, Frenk & White '97)

More massive halos and halos that form earlier have higher densities (bigger δ)

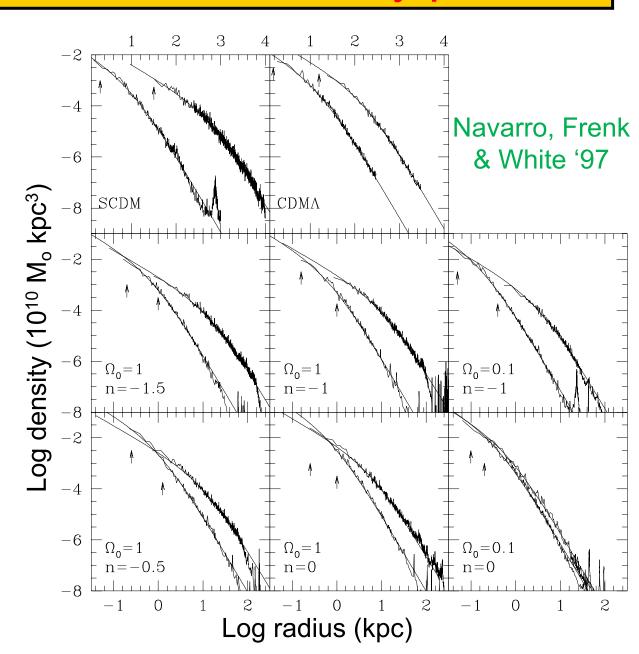


Universal halo density profiles

$$\frac{\rho(r)}{\rho_{crit}} = \frac{\delta_c}{(r/r_s)(1+r/r_s)^2}$$

Fits the spherically averaged density profiles of halos over a wide mass range.

2 parameters: Characteristic density δ_{C} radius: r_{s}



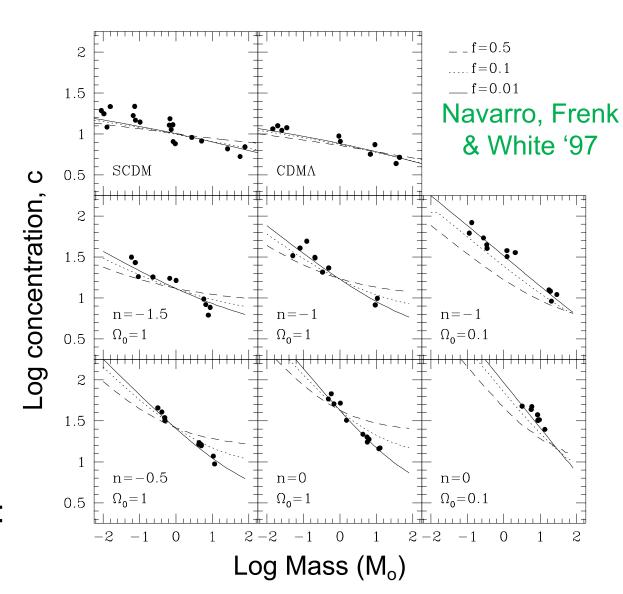


Universal halo density profiles

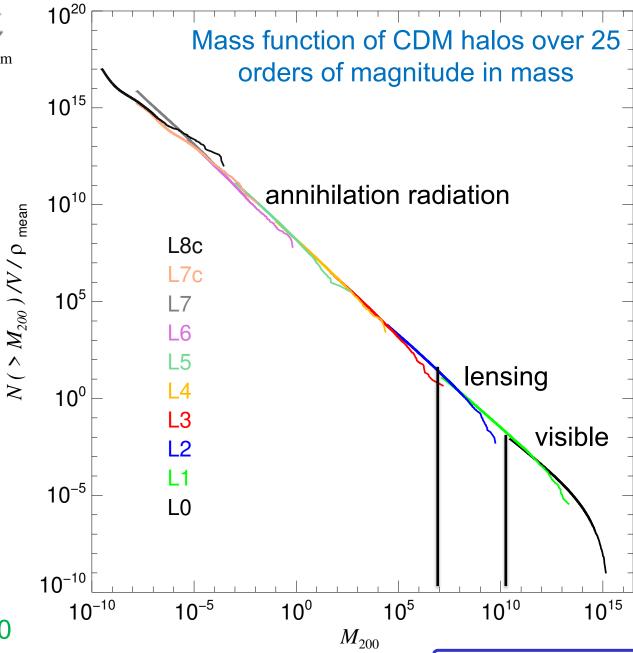
$$\frac{\rho(r)}{\rho_{crit}} = \frac{\delta_c}{(r/r_s)(1+r/r_s)^2}$$

2 parameters: Characteristic density, δ_{C} radius, r_{s}

The two parameters are related to halo mass in a way that is cosmology dependent: c \ as M \





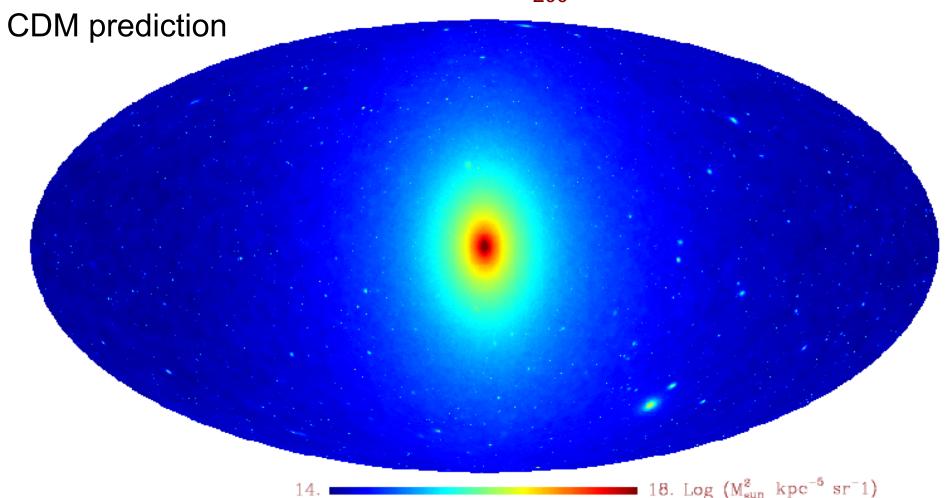




Springel et al. '08

The Milky Way seen in annihilation radiation

Aquarius simulation: $N_{200} = 1.1 \times 10^9$



Mass resolution − 10⁵ M_o



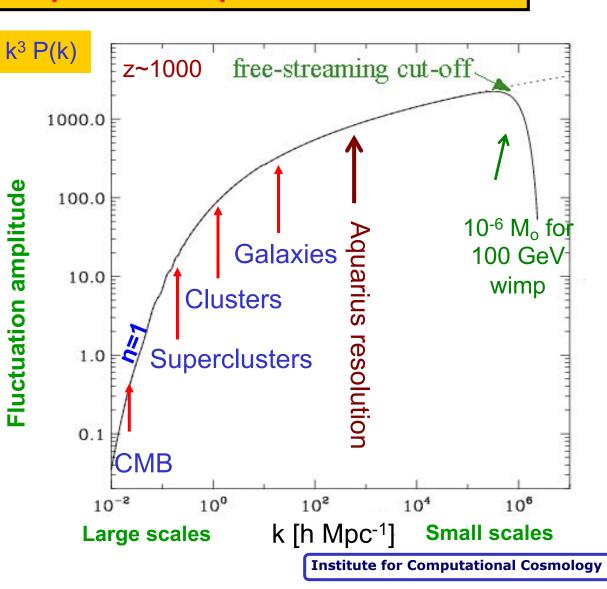
The cold dark matter linear power spectrum

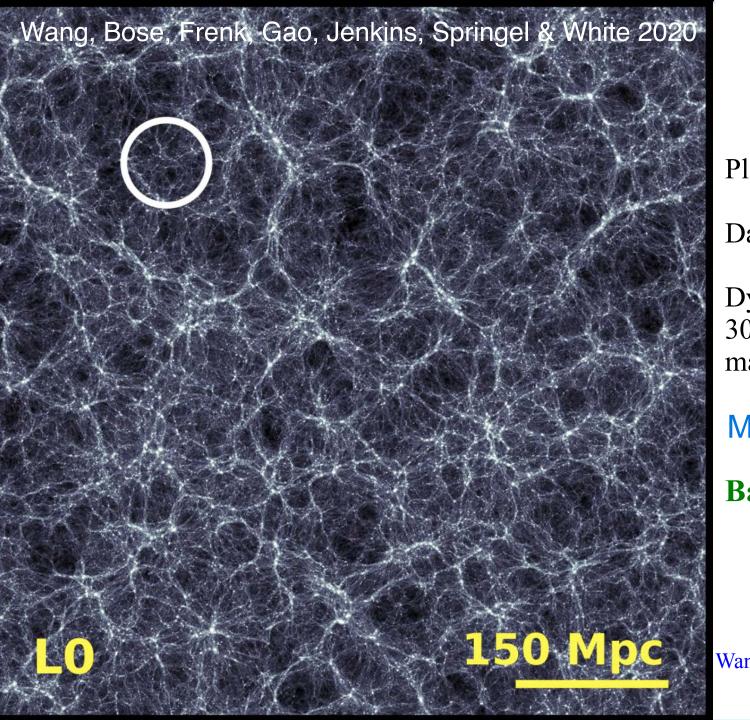
The linear power spectrum ("power per octave")

 $\lambda_{cut} \alpha m_x^{-1}$

Assumes a 100GeV wimp

Green et al '04





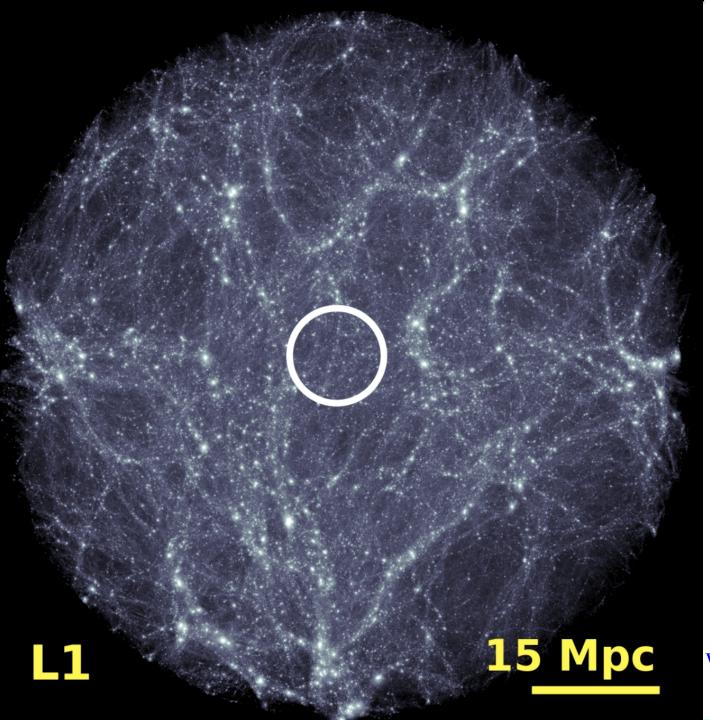
Planck cosmology

Dark matter only

Dynamic range of 30 orders of magnitude in mass

 $M_{char} = 10^{14} M_{\odot}$

Base Level



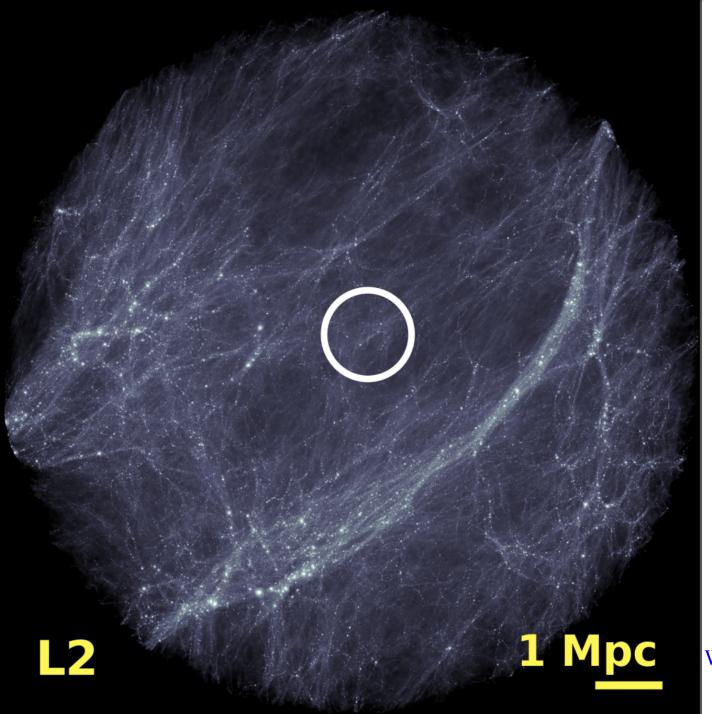
Planck cosmology

Dark matter only

Dynamic range of 30 orders of magnitude in mass

 $M_{char} = 10^{12} M_{\odot}$

Zoom Level 1



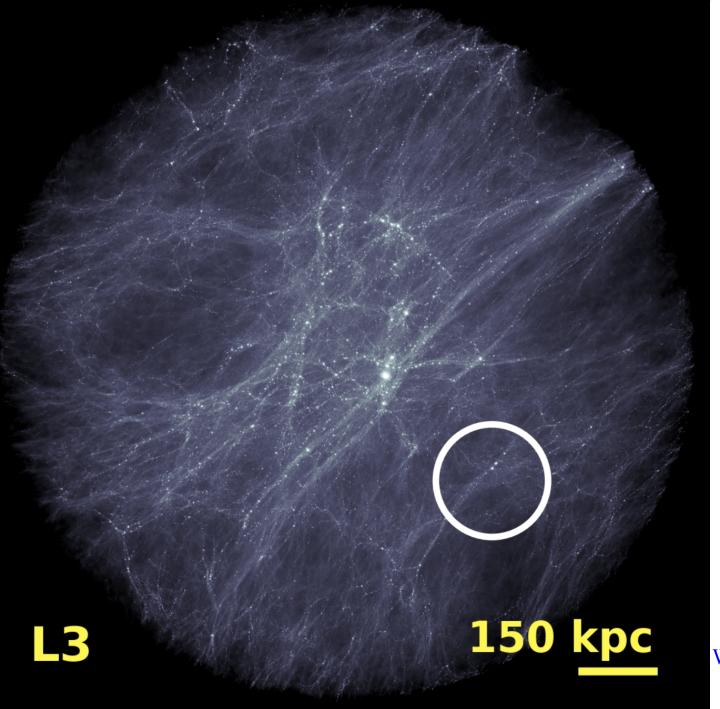
Planck cosmology

Dark matter only

Dynamic range of 30 orders of magnitude in mass

 $M_{char} = 10^9 M_{\odot}$

Zoom Level 2



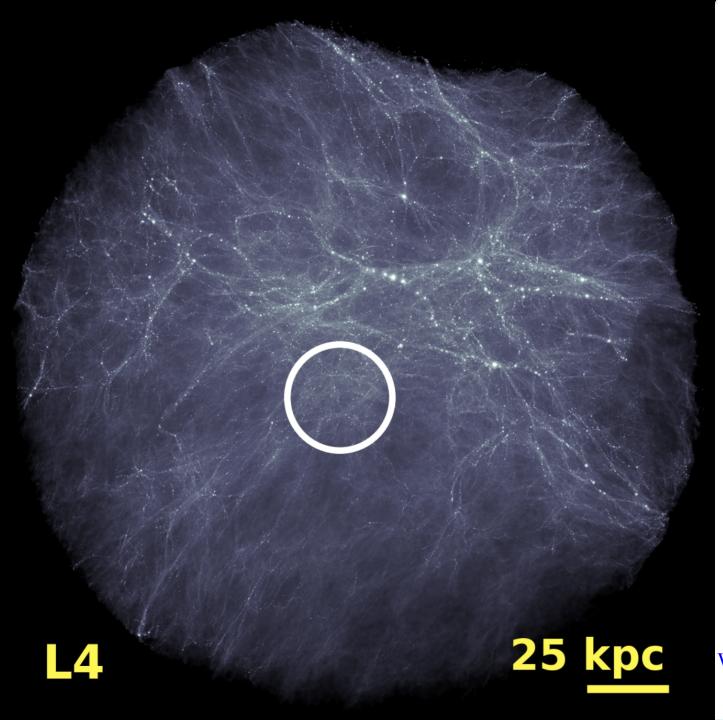
Planck cosmology

Dark matter only

Dynamic range of 30 orders of magnitude in mass

 $M_{char} = 10^6 M_{\odot}$

Zoom Level 3



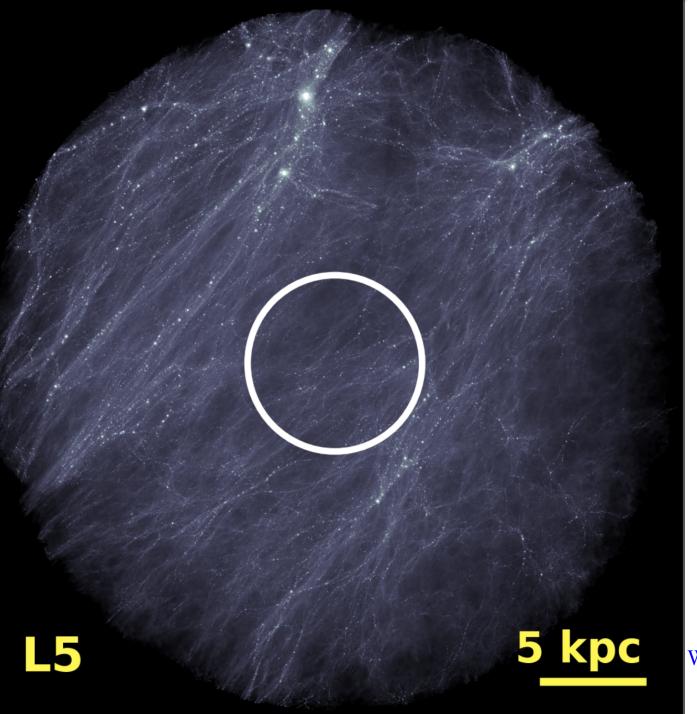
Planck cosmology

Dark matter only

Dynamic range of 30 orders of magnitude in mass

 $M_{char} = 10^3 M_{\odot}$

Zoom Level 4



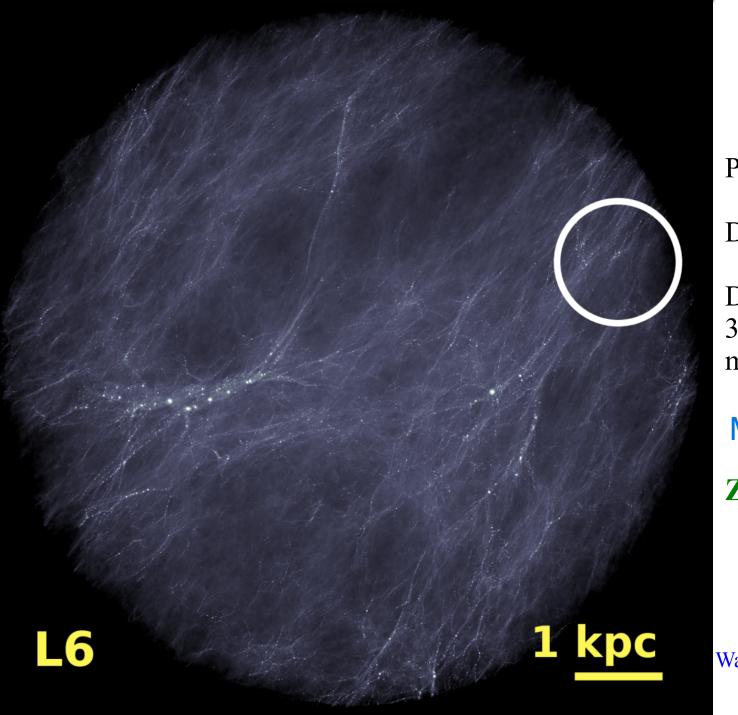
Planck cosmology

Dark matter only

Dynamic range of 30 orders of magnitude in mass

 $M_{char} = 10 M_{\odot}$

Zoom Level 5



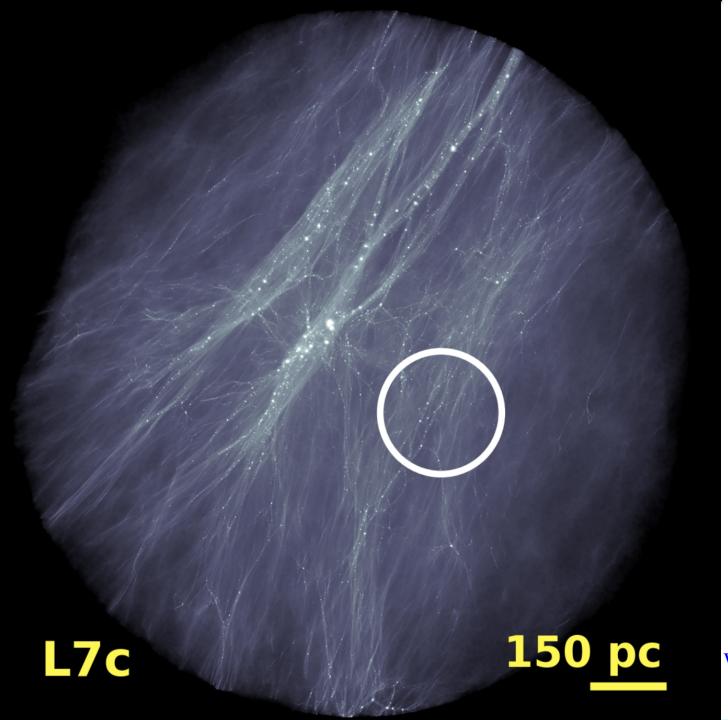
Planck cosmology

Dark matter only

Dynamic range of 30 orders of magnitude in mass

 $M_{char} = 10^{-1} M_{\odot}$

Zoom Level 6



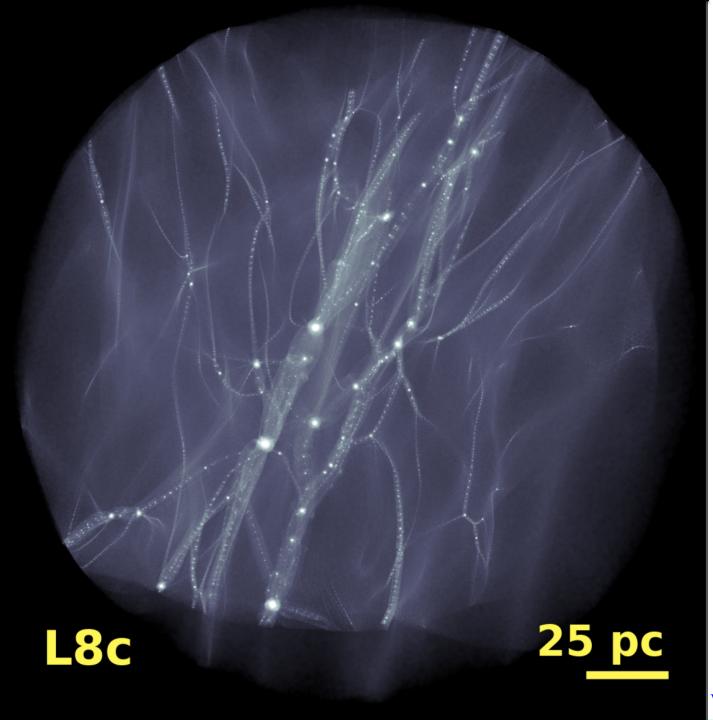
Planck cosmology

Dark matter only

Dynamic range of 30 orders of magnitude in mass

 $M_{char} = 10^{-4} M_{\odot}$

Zoom Level 7



Planck cosmology

Dark matter only

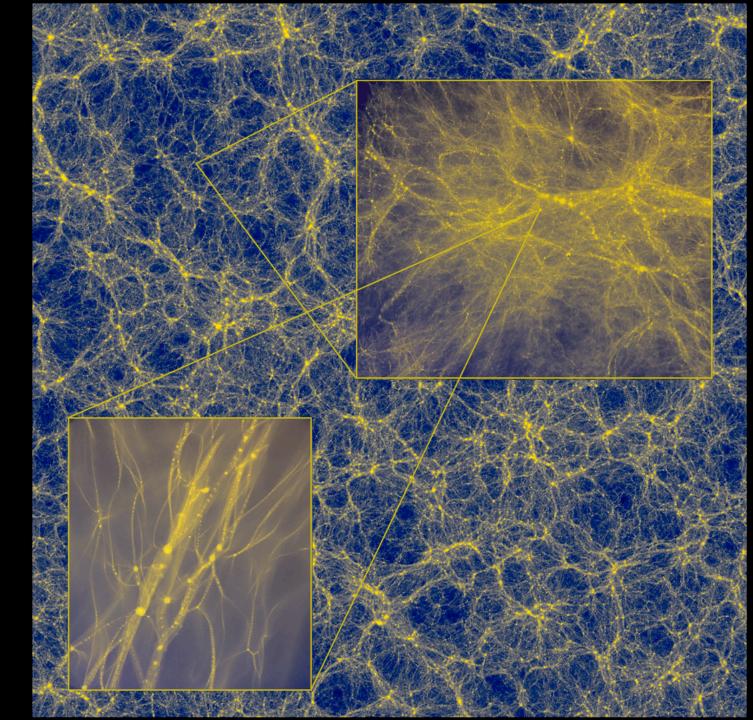
Dynamic range of 30 orders of magnitude in mass

 $M_{char} = 10^{-6} M_{\odot}$

Zoom Level 8

The density of this region is only ~3% of the cosmic mean



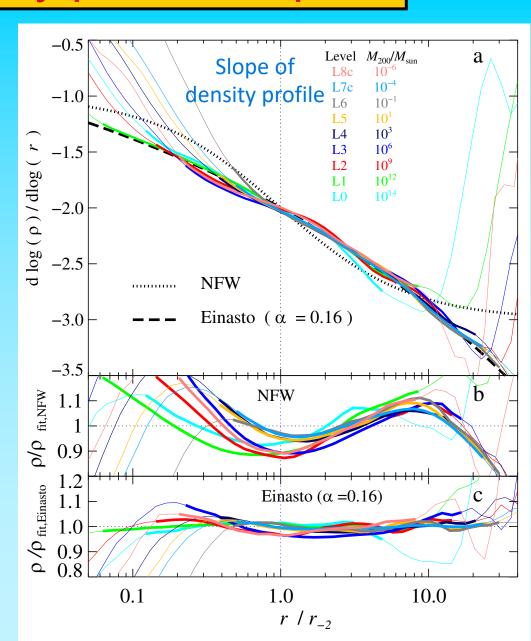


Wang et al '20



Density profile shapes

Over 19 orders of magnitude in halo mass and 4 orders of magnitude in density, the mean density profiles of halos are fit by NFW to within 20% and by Einasto $(\alpha = 0.16)$ to within 7%



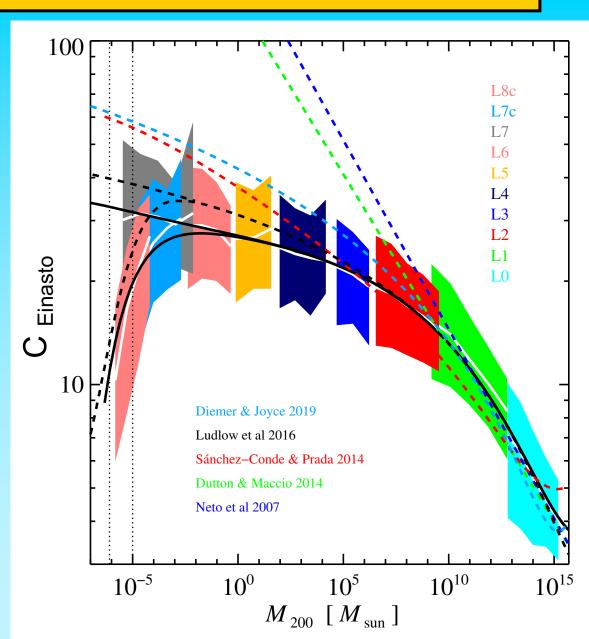


Concentration-mass relation

Concentrations at small mass are lower than all previous extrapolations by up to factors of tens.

A turndown at 10³ Earth masses is due to the freestreaming limit.

The scatter depends only weakly on halo mass

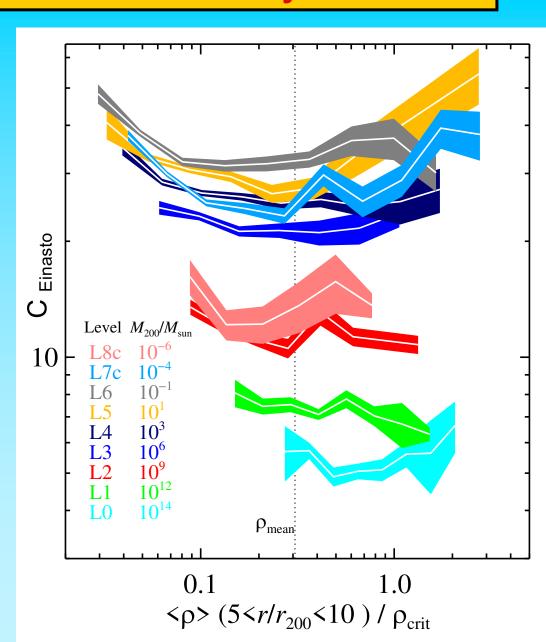




Concentration-density relation

At given halo mass, concentration does not depend on local environment density

The range of local environment density does not depend strongly on halo mass





Annihilation luminosity

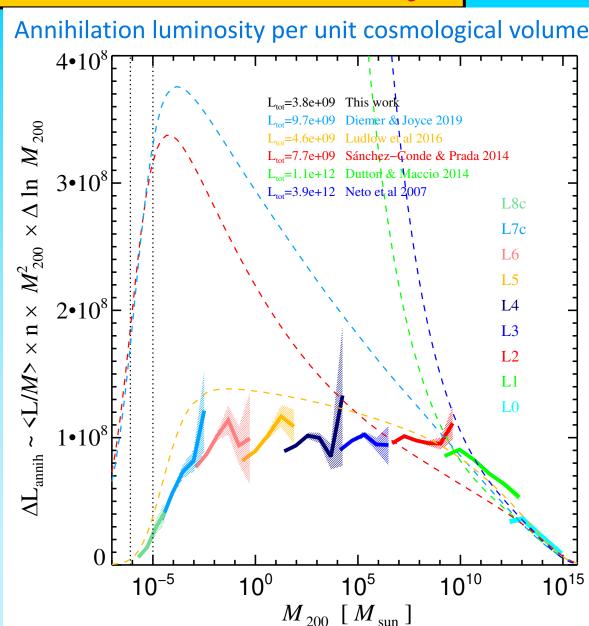
The contribution of halos
to the mean z = 0
luminosity density of the
Universe is almost
independent of their mass
over the mass range

$$10^{-4}~{\rm M}_{\odot} < {\rm M}_{\rm halo} < 10^{12}~{\rm M}_{\odot}$$

It is lower than previously estimated by factors between 3 and 1000

This still neglects the substructure contribution to halo luminosity

Wang, Bose, CSF + '20





Conclusions

- ΛCDM: great success on scales > 1Mpc: CMB, LSS, gal evolution
- But on these scales ACDM cannot be distinguished from WDM
- The identity of the DM makes a big difference on small scales
- CDM makes many small subhalos but most (~5.10⁸M₀)
 are dark → No satellite problem in CDM or WDM
- 2. Distortions of strong gravitational lenses offer a clean test of CDM vs WDM → and can potentially rule out CDM!
- 3. Halos of all masses (21 orders of magnitude) have NFW profiles -> small haloes may dominate annihiln radn